



PASCOS '04

Northeastern university, Boston, 16-22 August 2004

Dark Matter At LHC

Stefano Lacaprara

Stefano.Lacaprara@pd.infn.it

INFN and Padova University



Outline



- Dark Matter and SuperSymmetry (SUSY),
- SUSY Scenarios in mSUGRA,
- LHC reach for SUSY,
- Sparticles spectroscopy,
- SUSY Dark Matter with LHC,
- Conclusion

- Large fraction of universe mass is dark (WMAP)

in units of $\rho_c = 1.86 \cdot 10^{-29} \text{ g/cm}^3$

Barion density: $\Omega_b = 0.044 \pm 0.004$

Dark energy: $\Omega_\Lambda = 0.73 \pm 0.04$

Dark matter: $\Omega_m = 0.23 \pm 0.04$

- Dark Matter density: $0.0094 < \Omega_m h^2 < 0.129$ (2σ range)
- SuperSymmetry provide a good candidate for Dark Matter
 - If **R-Parity** is conserved, the lightest supersymmetric particle (LSP) is stable
 - Relic density of LSP can explain (fraction) of Dark Matter
- “*Natural*” candidate from a theory developed to solve different problems!

- Use constrained MSSM: *mSUGRA*
- Only 5 free parameters (105 + 19 SM for generic MSSM)

m_0	universal scalar mass
$m_{1/2}$	universal gaugino mass
$\tan \beta$	ratio of two Higgs doublets VEV's
$\text{sign}(\mu)$	Higgs mixing parameter sign
A_0	trilinear SUSY breaking parameter

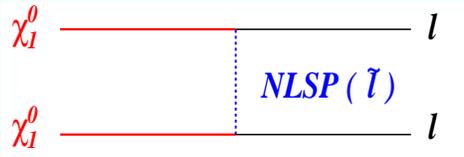
- LSP is the lightest neutralino: χ_1^0
- NLSP is usually a slepton: $\tilde{\tau}$
- Heaviest SP is usually the gluino: \tilde{g}

Dark Matter due to relic χ_1^0

Relic density: $\rho_{\chi_1^0} = n_{\chi_1^0} \times m_{\chi_1^0}$

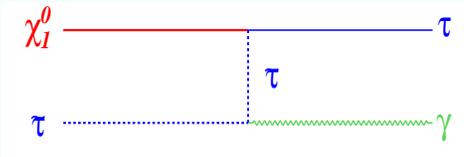
$n_{\chi_1^0}$ depends on LSP annihilation rate at early universe

Bulk:

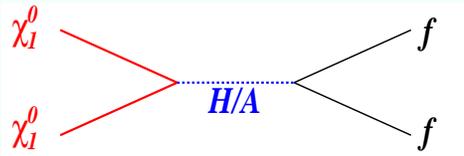


Co-Annihilation tail: $\tilde{\tau} \rightarrow \tau \chi_1^0$ forbidden.

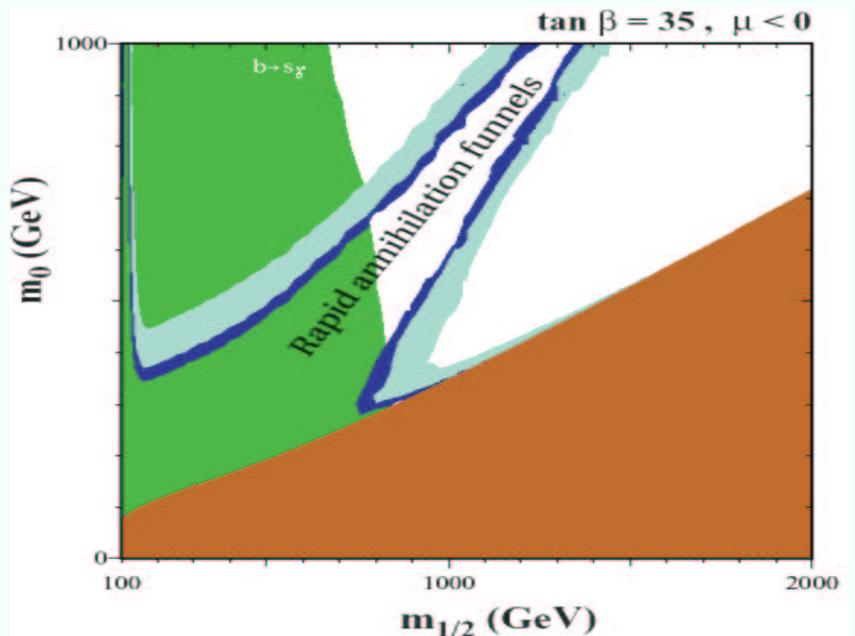
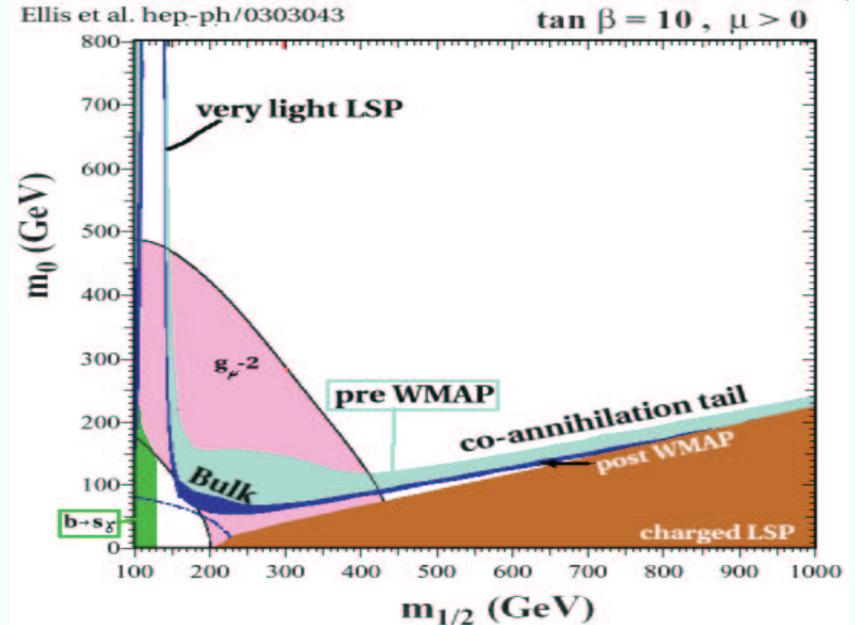
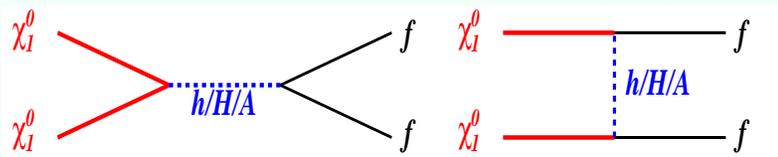
$\tilde{\tau}$ lives longer



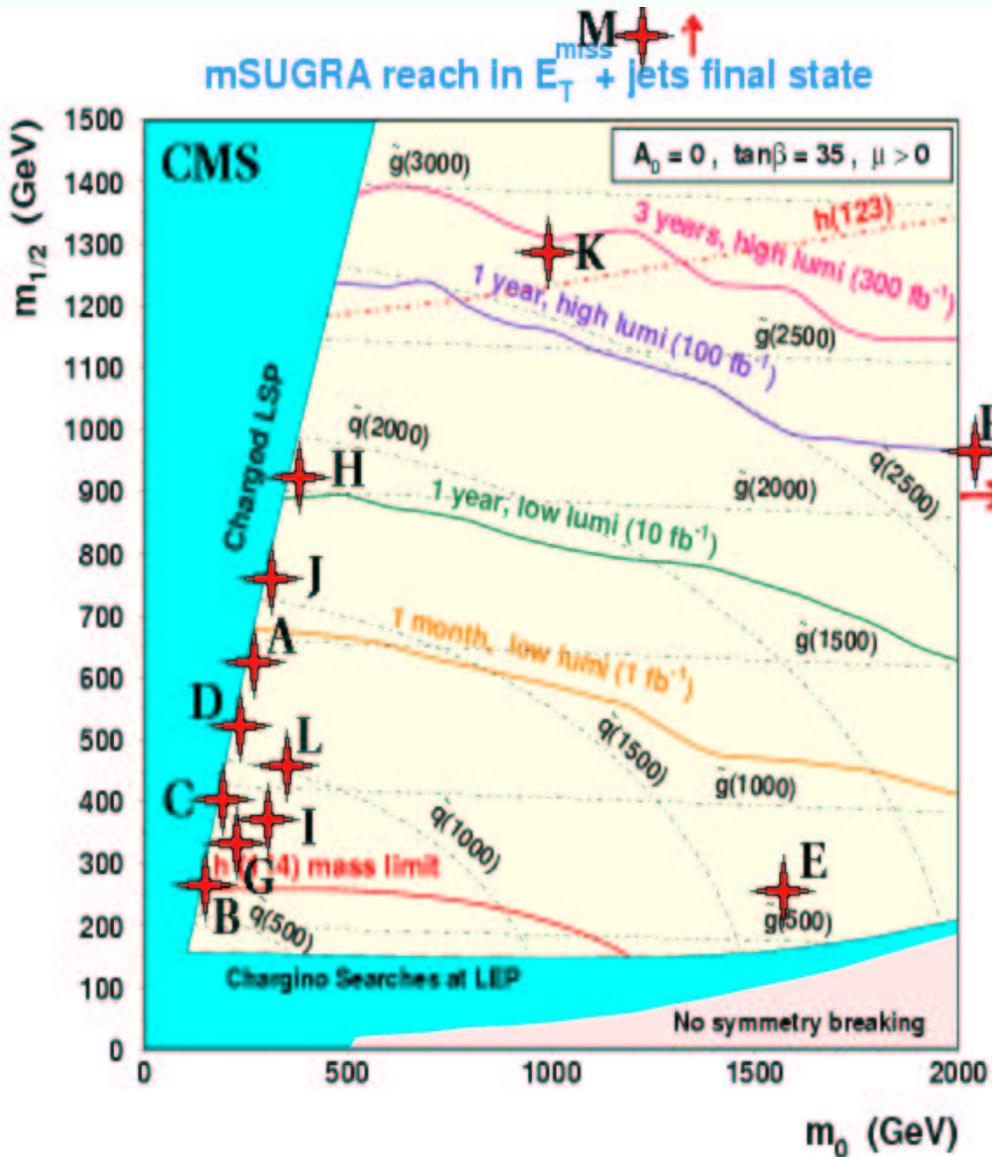
Funnels (large $\tan \beta$): $m_0 \sim m_{1/2}$,
 $m_{\chi_1^0} \sim 1/2 m_{A/H}$ resonant annihilation (if exactly on-peak too rapid)



Focus point: $\chi_1^0 \sim 100\%$ higgsino
 off-shell annihilation via h, H, A dominant

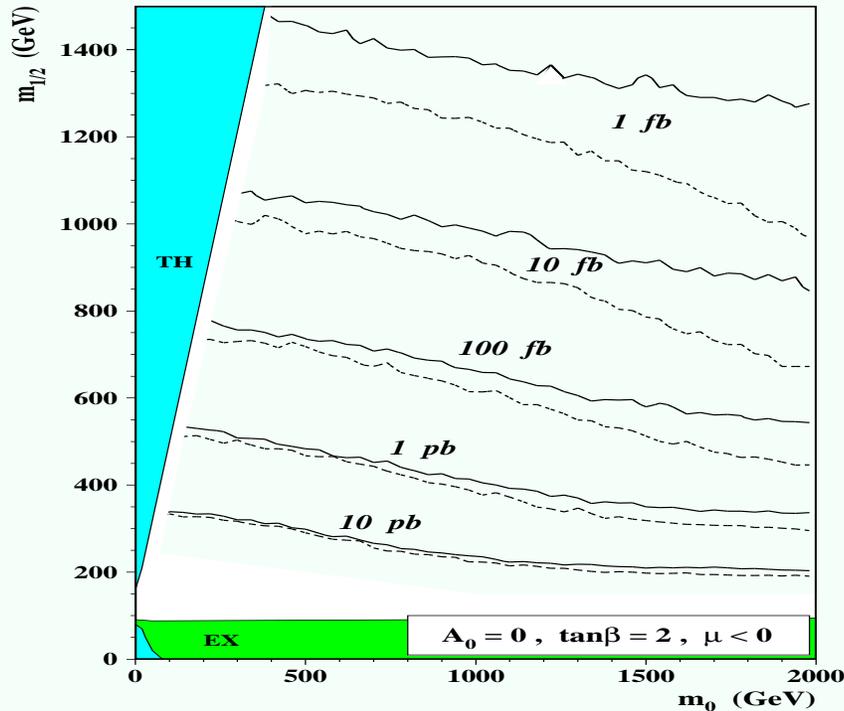


Benchmark points on LHC reach from \tilde{g}, \tilde{q} production (jets+missing E_T)



Point	$m_{1/2}$	m_0	$\tan\beta$	$\text{sign}(\mu)$
A	600	140	5	+
B	250	100	10	+
C	400	90	10	+
D	525	125	10	-
E	300	1500	10	+
F	1000	3450	10	+
G	375	120	20	+
H	1500	419	20	+
I	350	180	35	+
J	750	300	35	+
K	1150	1000	35	-
L	450	350	50	+
M	1900	1500	50	+

- At LHC dominant production for \tilde{g} and \tilde{q}
- Strong production \Rightarrow huge cross section!

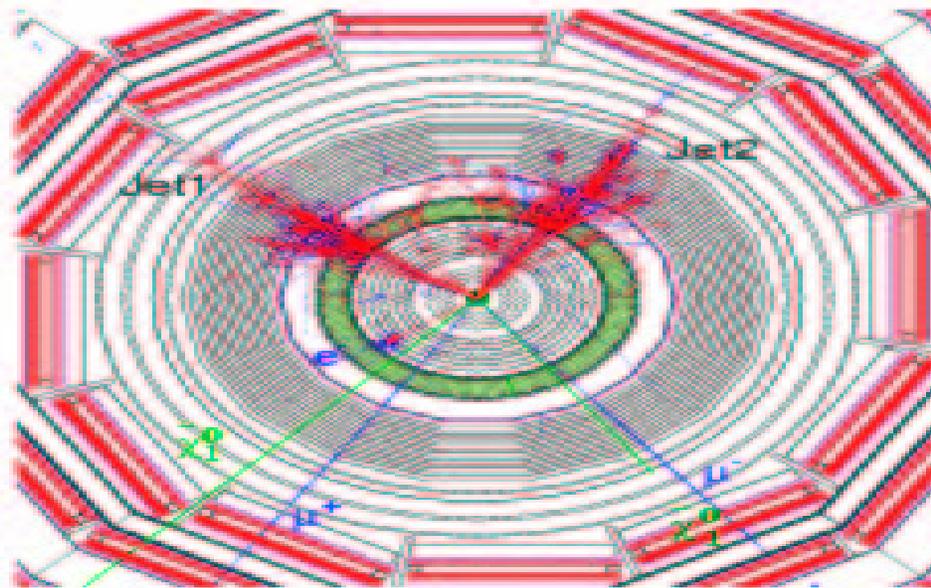
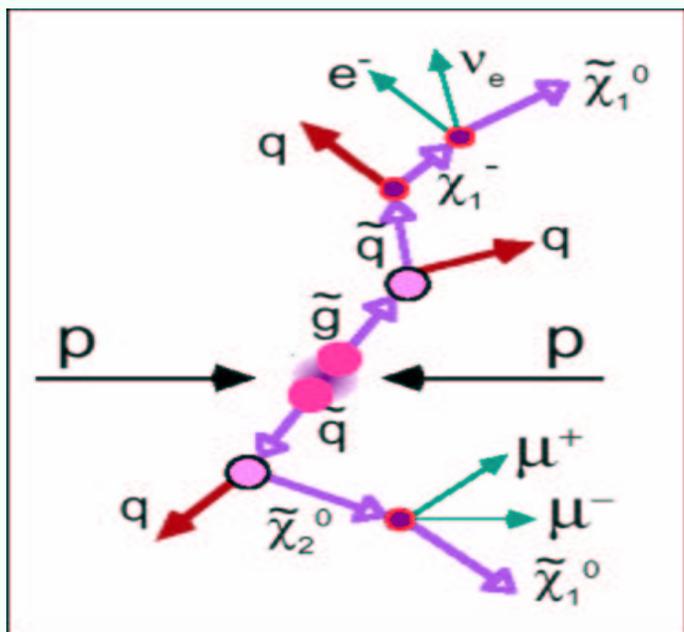
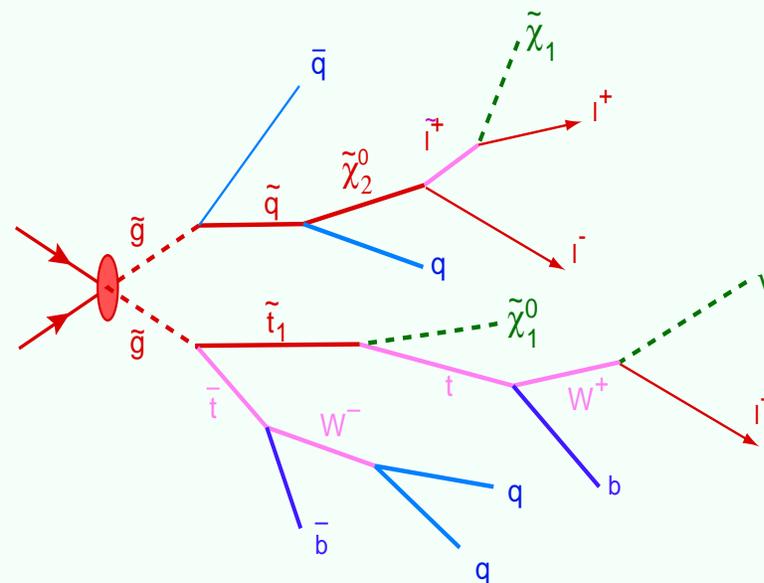


- Almost insensitive to $\tan \beta$ and $\text{sign}(\mu)$
- Some reference:
 - $\sigma(tt\bar{t}) \sim 1 \text{ nb}$
 - $\sigma(Z \rightarrow l\nu) \sim 3 \text{ nb}$
 - At High Level Trigger:
Rate(MET+jets) $\sim 5 \text{ Hz}$

- LHC low luminosity: ($\mathcal{L} = 2\text{nb}^{-1}\text{s}^{-1}$)
- for $m_0, m_{1/2}$ just above limits:
SUSY visible at trigger rate!

Typical SUSY event topology:

- isolated leptons (3)
- b -jets (2)
- jets (4)
- large missing energy

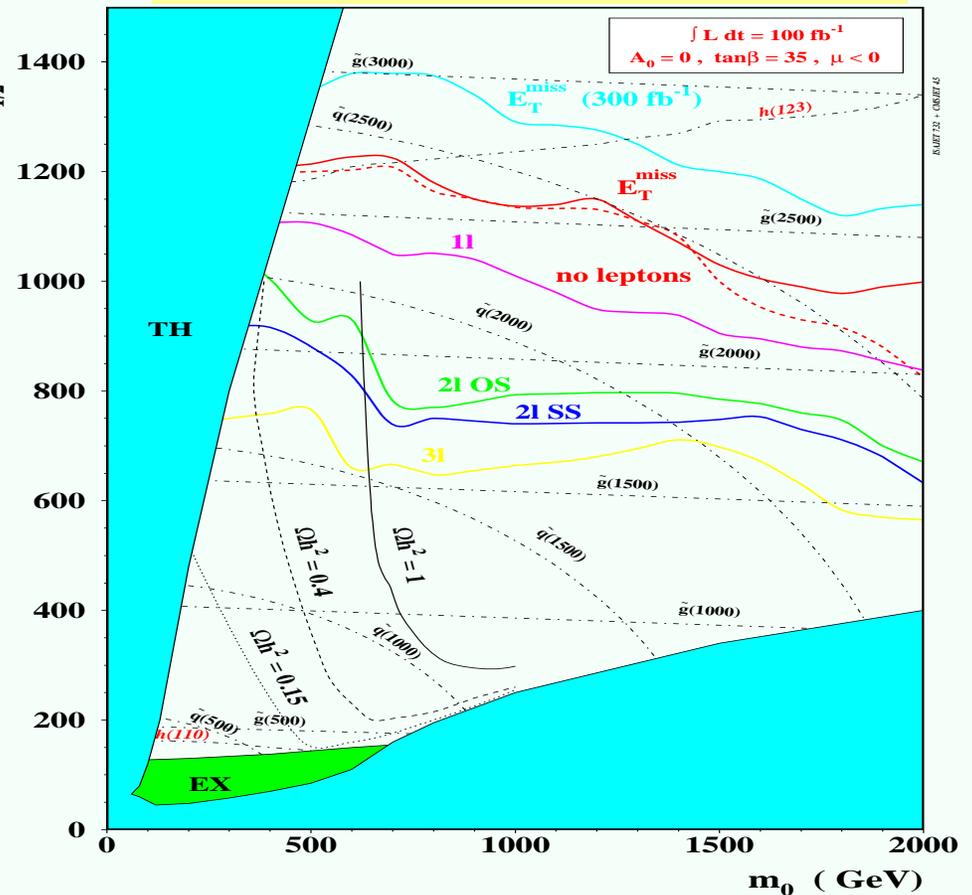
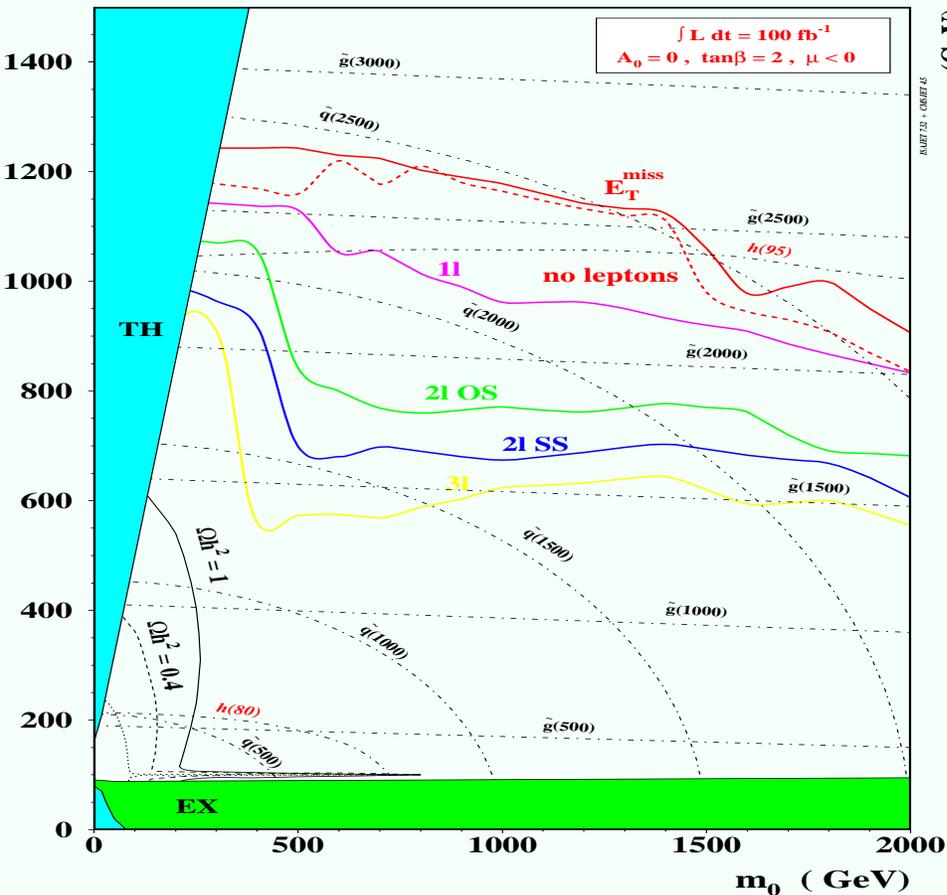


Trigger strategy is rather simple: many handles (CMS)

- Get maximum trigger efficiency and keep under control trigger rate
- Apply very simple selection to reduce model-dependent bias
- Low lumi:
 - 1 jets $E_T > 180$ GeV & missing $E_T > 123$ GeV
 - 4 jets $E_T > 113$ GeV
 - Efficiency $\epsilon \sim 60 \div 70\%$ for low $M_{\tilde{q}} \sim 400$ GeV
 - Background rate (QCD) ~ 12 Hz
- High lumi:
 - missing $E_T > 239$ GeV
 - 4 jets $E_T > 185$ GeV
 - $\epsilon(M_{\tilde{q}} \sim 2$ TeV) $\sim 75 \div 90\%$
- More complex exclusive trigger possible
- Trigger consider also R_p violating models

First step: discover SUSY!
 Several final states investigated:
 Final reach $m_{\tilde{g}, \tilde{q}} \sim 2.6 \div 3$ TeV

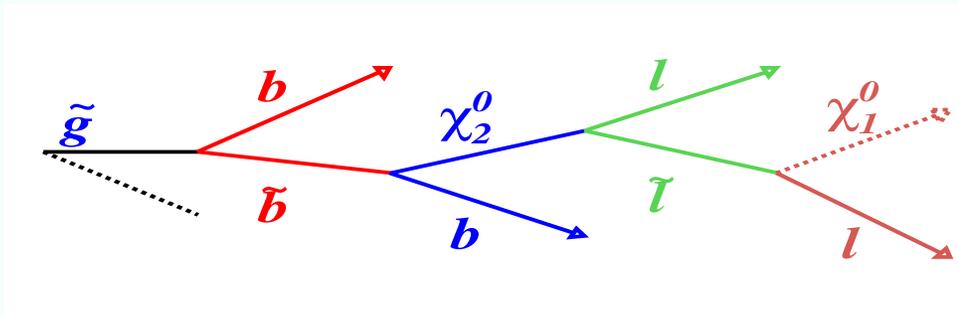
- $E_T^{miss} + \text{jets}$
- $E_T^{miss} + \text{jets} + \text{no leptons}$
- $E_T^{miss} + \text{jets} + 1 \text{ lepton}$
- $E_T^{miss} + \text{jets} + 2 \text{ OS leptons}$
- $E_T^{miss} + \text{jets} + 2 \text{ SS leptons}$
- $E_T^{miss} + \text{jets} + 3 \text{ leptons}$



Second step: measure sparticle spectrum!

- Use benchmark point **B**:
- $m_0 = 100 \text{ GeV}$, $m_{1/2} = 250 \text{ GeV}$, $\tan \beta = 10$, $\mu > 0$, $A_0 = 0$
- Most favorable scenario: very light sparticles
- $\sigma_{SUSY}^{tot} \approx 58 \text{ pb}$

Use long decay chains:



- ≥ 2 SFOS isolated leptons
- ≥ 2 b -jets
- E_T^{miss}
- Almost SM background free with simple cuts

Start with $\chi_2^0 \rightarrow \chi_1^0 l^+ l^-$ decay

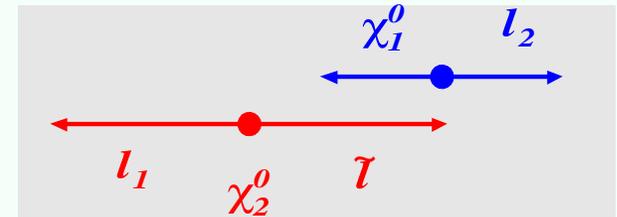
- $\text{BR}(\chi_2^0 \rightarrow \chi_1^0 l^+ l^-) \approx 0.04\%$
- $\text{BR}(\chi_2^0 \rightarrow \tilde{l}^\pm l^\mp \rightarrow \chi_1^0 l^+ l^-) \approx 16.4\%$
- $\text{BR}(\chi_2^0 \rightarrow \tilde{\tau}^\pm \tau^\mp \rightarrow \chi_1^0 \tau^+ \tau^-) \approx 83.2\%$

- Final states with e^\pm or μ^\pm sizeable
- At larger $\tan \beta$ (eg point **G**), $\text{BR}(\rightarrow \tau^\pm) \approx 100\%$ less favorable!

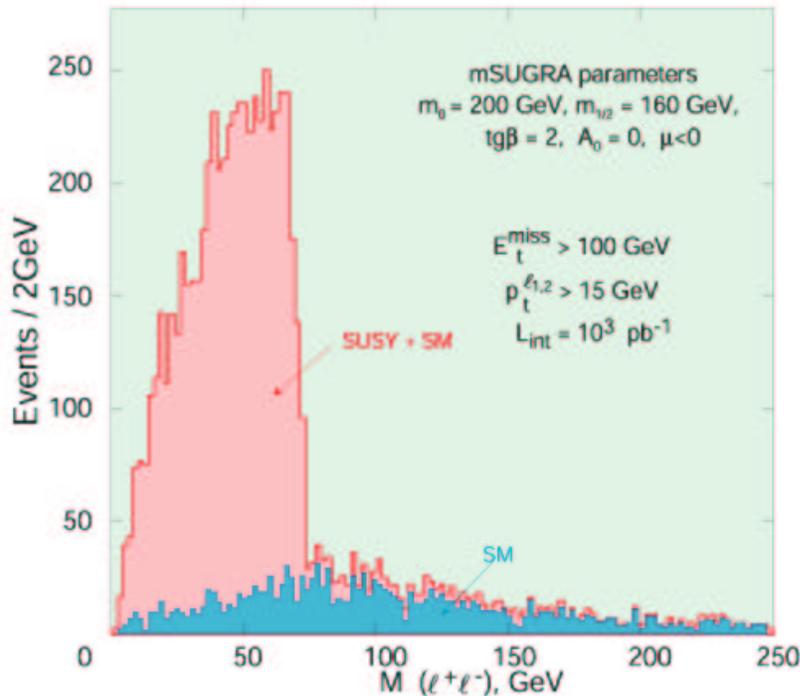
- Dilepton invariant mass has sharp upper edge
- due to 2 body decay of χ_2^0

$$M_{ll}^{max} = \frac{1}{m_{\tilde{l}}} \sqrt{\left(m_{\chi_2^0}^2 - m_{\tilde{l}}^2\right) \left(m_{\tilde{l}}^2 - m_{\chi_1^0}^2\right)}$$

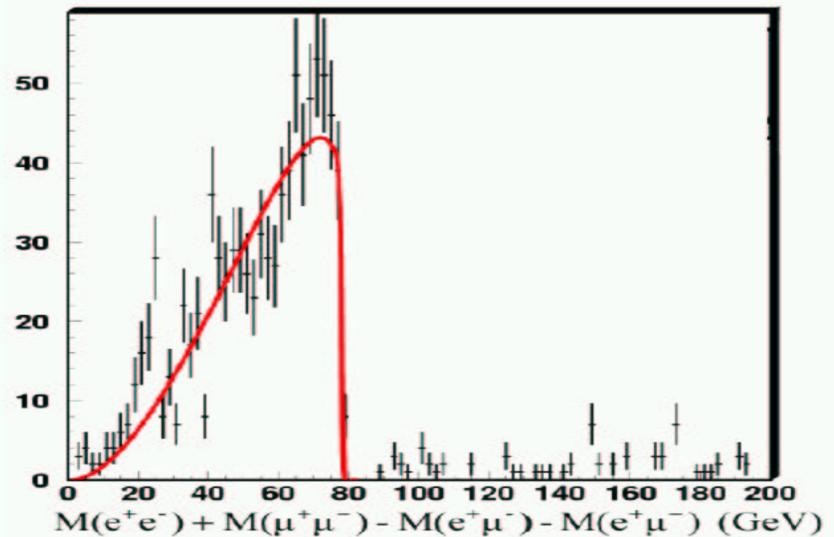
When l^\pm back to back in χ_2^0 rest frame



Inclusive $l^+l^- + E_t^{miss}$ final states



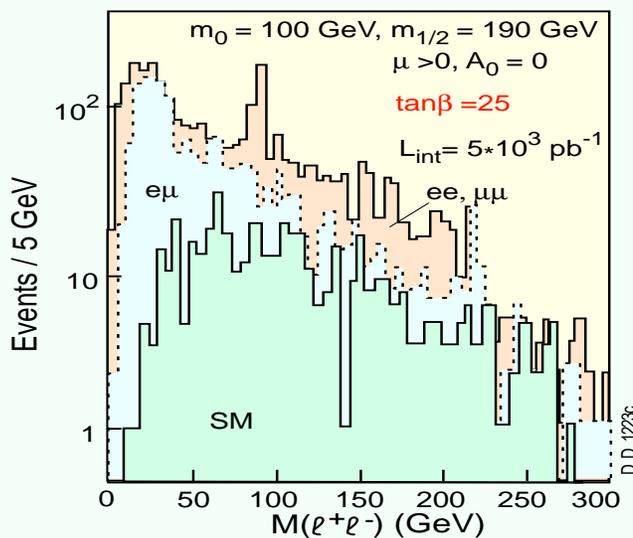
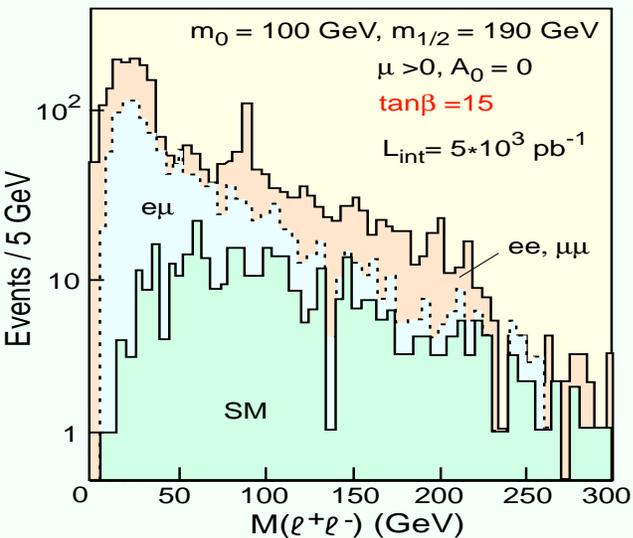
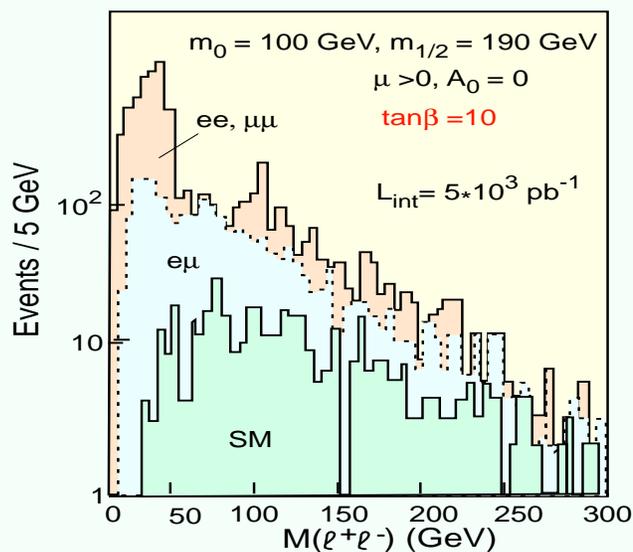
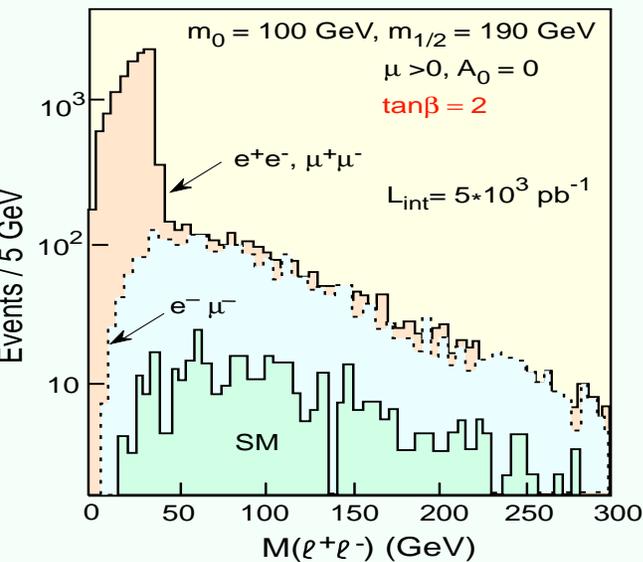
Subtract OFOS leptons reduce background (mostly from SUSY)



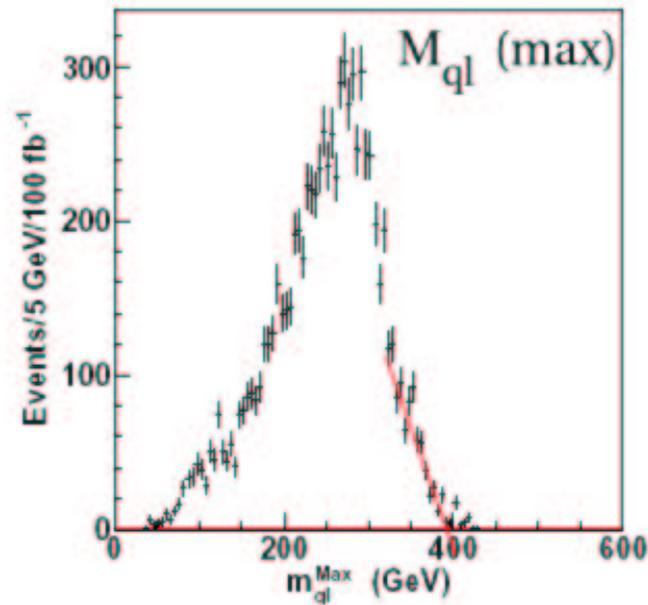
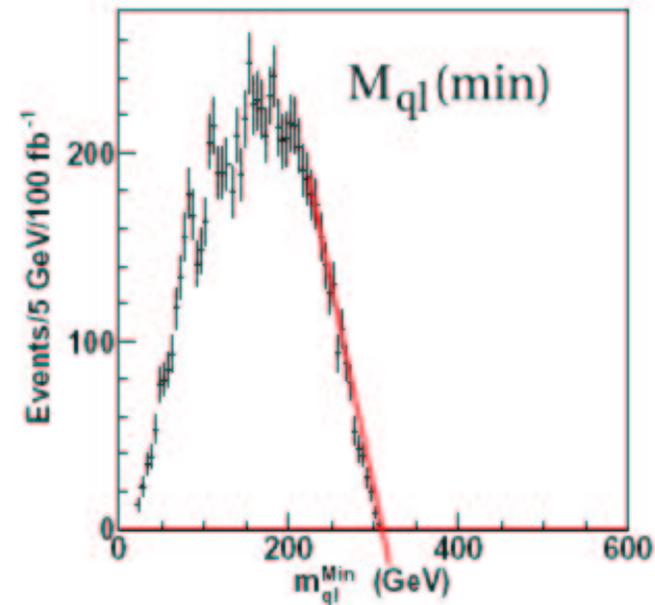
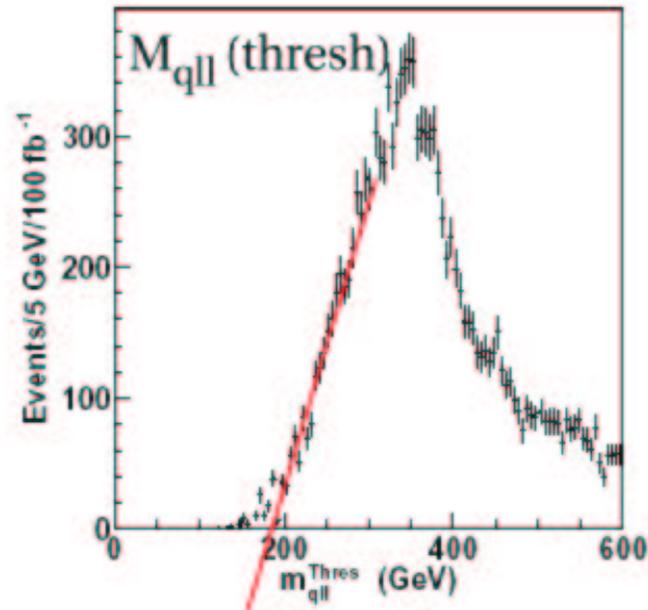
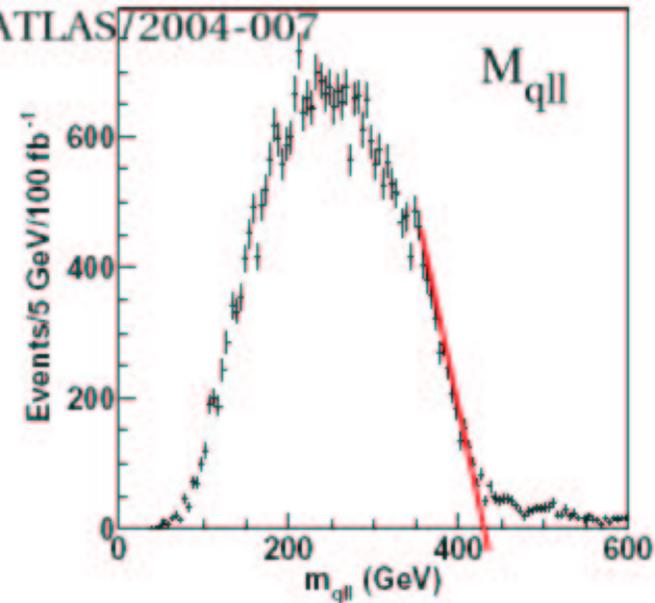
Dilepton mass edge vs $\tan \beta$

Dilepton invariant mass vs $\tan \beta$

at fixed $m_0 = 100$ GeV, $m_{1/2} = 190$ GeV, $A_0 = 0$, $\mu < 0$



- Evidence for a M_{II} edge \Rightarrow SUSY!
- Structure less evident with increasing $\tan \beta$: $\chi_2^0 \rightarrow \chi_1^0 \tau^\pm$ dominates



Edges present for other invariant mass distribution

-  M_{qll}
-  M_{ql}^{min}
-  M_{ql}^{max}
-  M_{qll}^{thres}

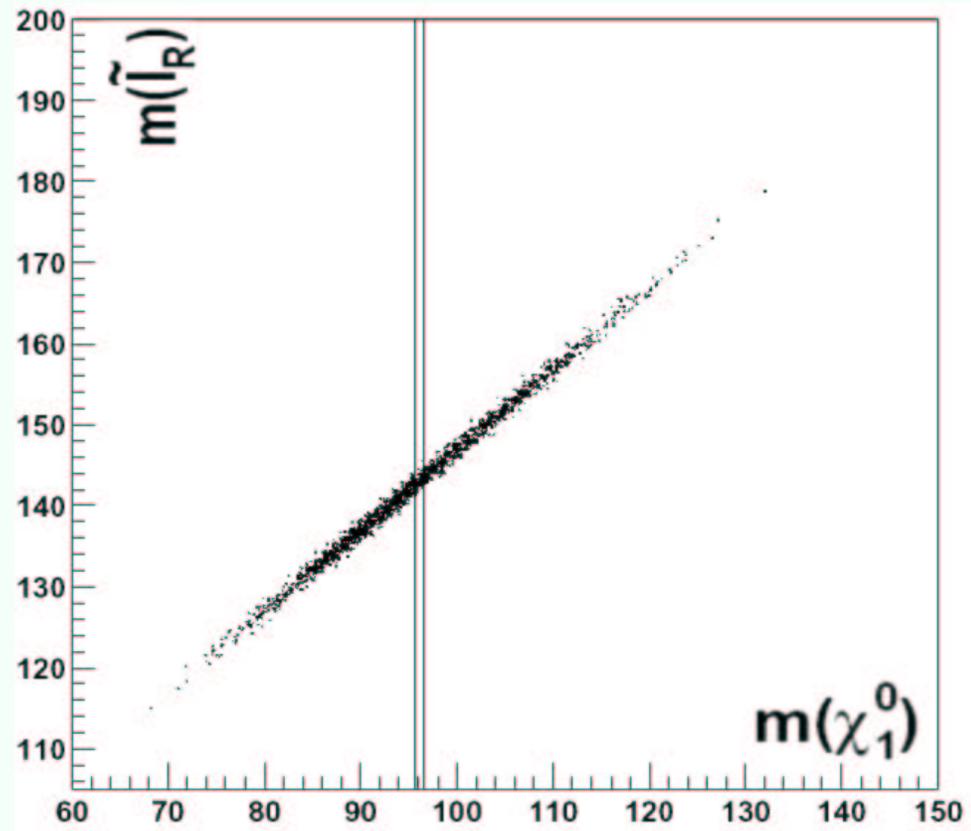
All functions of m_p^{2art} differences

Error on endpoints: statistical, energy scale, fit model

Fitted values agree well with MC input

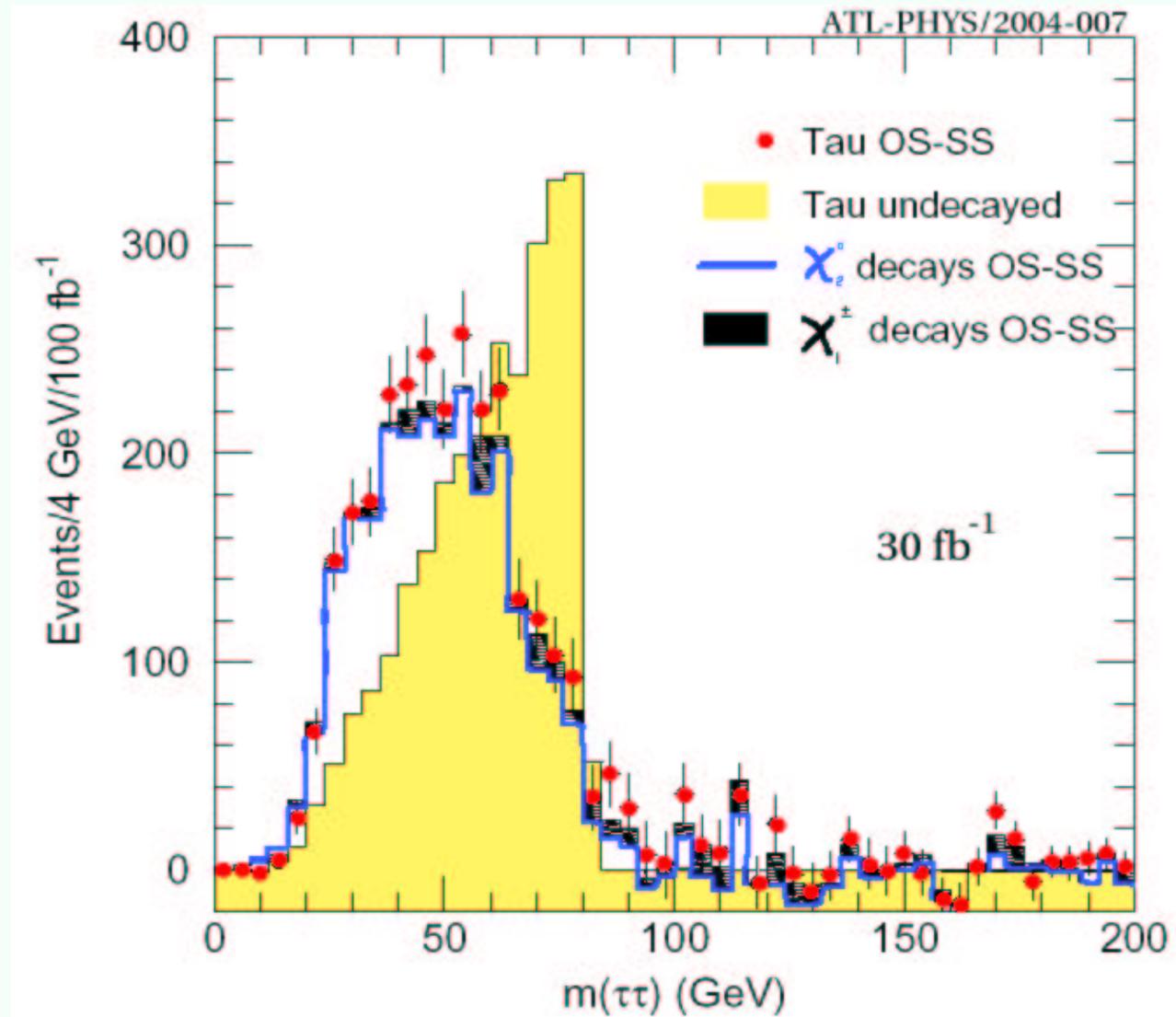
- Use 6 edge measurements: $M_{\tilde{q}_L}, M_{\tilde{b}}, M_{\chi_1^0}, M_{\chi_2^0}, M_{\tilde{l}_R}$
- Edges depends on mass difference \Rightarrow strong dependencies!
- Measure mass relation as a function of $M_{\chi_1^0}$ to $\sim 1\%$
- $M_{\chi_1^0}$ determined to $\sim 10\%$

- If $M_{\chi_1^0}$ from extra LHC (eg LC) \Rightarrow fix mass scale for all other sparticles!
- Use to measure heavier sparticle masses



- At point **B** dominant χ_2^0 decay mode into $\tilde{\tau}\tau$ ($\sim 80\%$)
- At higher $\tan\beta$ (eg point **G**, point **I**) BR higher

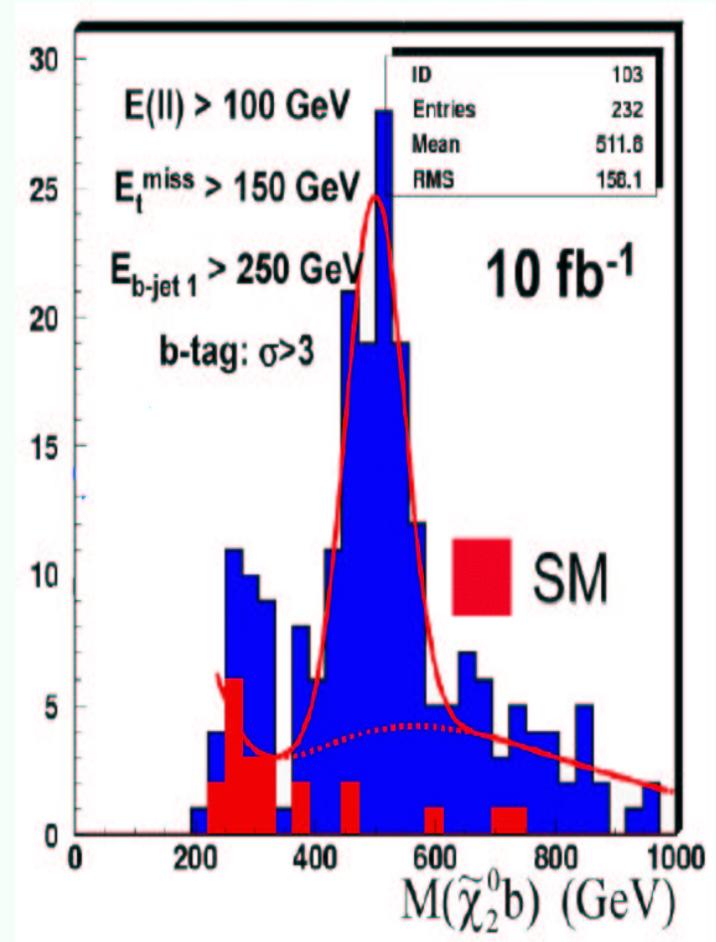
- Study final state with τ s
- Use τ hadronic decay:
 τ -tagged jets
- Subtract same sign τ 's
to reduce SUSY
background
- SM background small
- Edge on $M_{\tau\tau}$ invariant
mass partially spoiled by
 ν_{τ} 's
- Distribution sensitive to
edge position



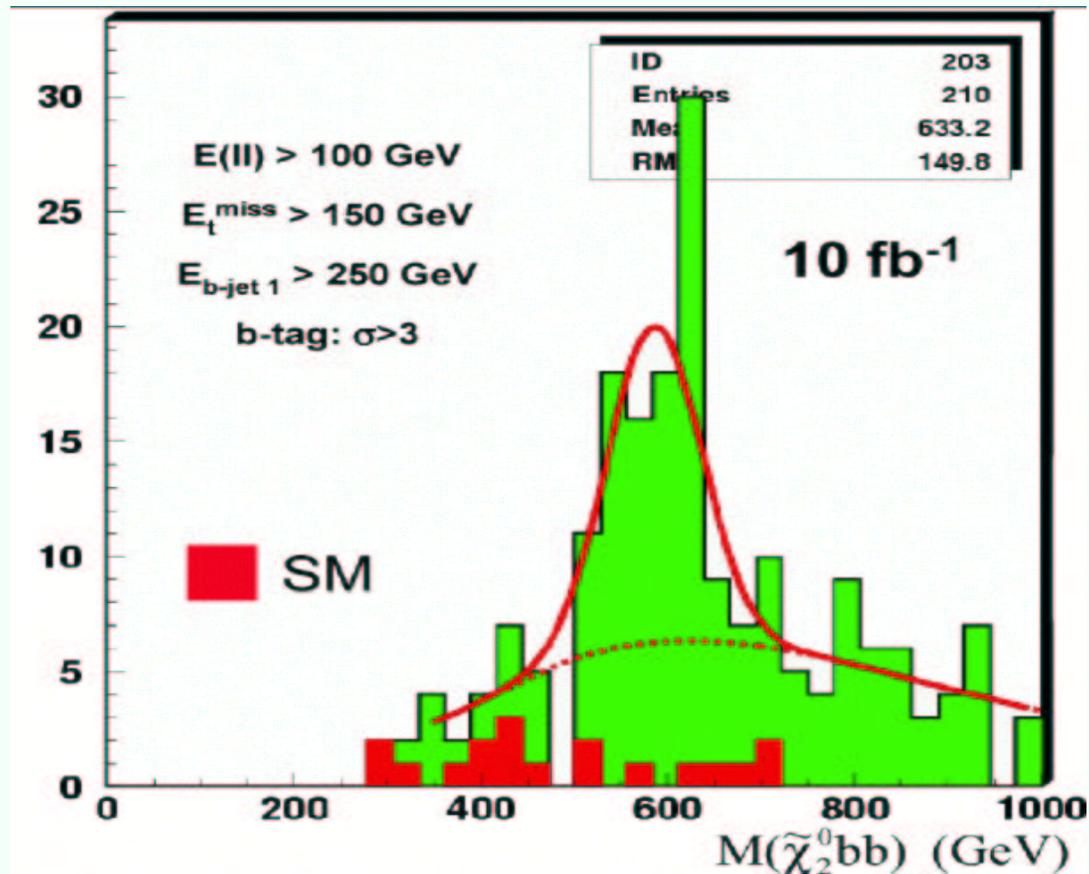
- Near the M_{l^l} edge:

$$\vec{p}_{\chi_2^0} = \left(1 + \frac{M_{\chi_1^0}}{M_{ll}} \right) \vec{p}_{ll}$$

- Need to know $M_{\chi_1^0}$!
- Add most energetic b -jet to reconstruct sbottom
- Fit result: $M_{\chi_2^0 b} = 500 \pm 7$ GeV
- Generated: $M_{\tilde{b}_1} = 496$ GeV,
 $M_{\tilde{b}_2} = 524$ GeV
- With full statistics (300 fb^{-1}) possible to resolve \tilde{b}_1 and \tilde{b}_2



- Combine \tilde{b} with b -jet
- Add other b -jets (closest in ϕ)
- Get gluino mass peak!
- Fit result: $M_{\chi_2^0 bb} = 594 \pm 7 \text{ GeV}$
- Generated: $M_{\tilde{g}} = 596 \text{ GeV}$



- Same procedure also for non b squark ($\tilde{u}, \tilde{d}, \tilde{c}, \tilde{s}$)
- Use non b -jets
- Higer statistics!
- \tilde{q} mass peak well visible with just 1 fb^{-1}
- Higer combinatorial background

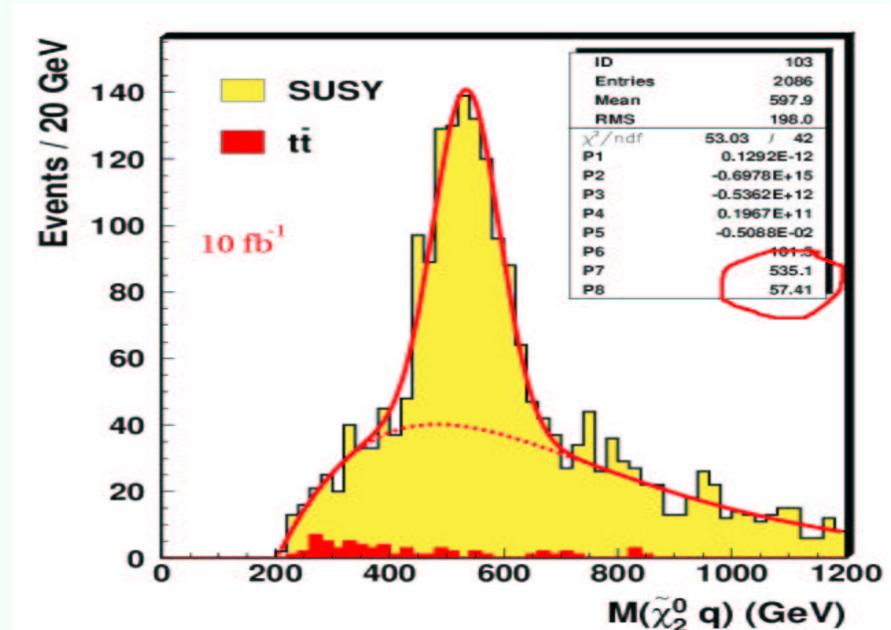
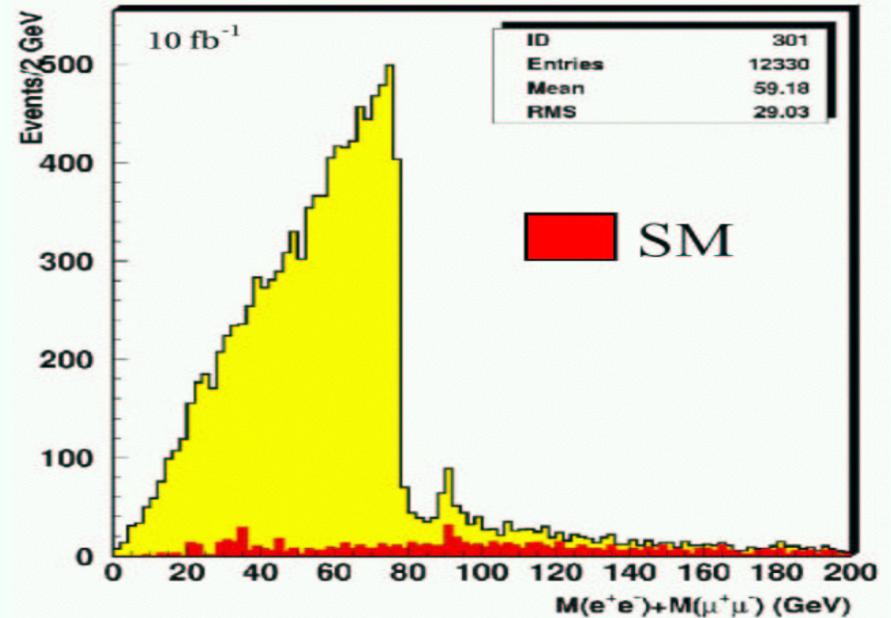
Fit result:

$$M_{\chi_2^0 q} = 535 \pm 3 \text{ GeV}$$

Generated:

$$M_{\tilde{u}_L} = M_{\tilde{c}_L} = 537 \text{ GeV}$$

$$M_{\tilde{d}_L} = M_{\tilde{s}_L} = 542 \text{ GeV}$$



Alternative approach (ATL-PHYS/2003-039)

Use same decay chain

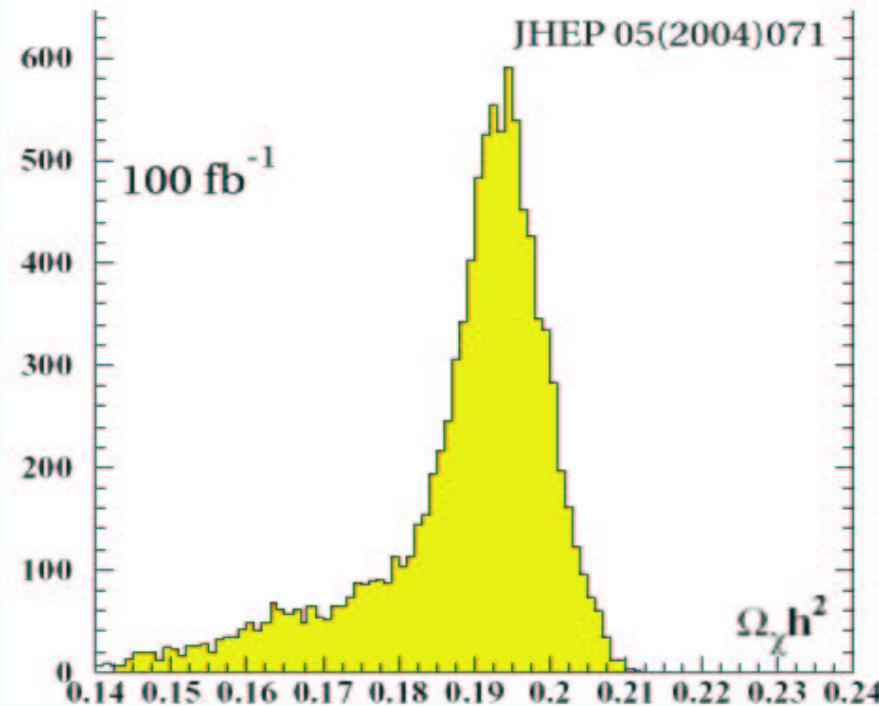
$$\tilde{g} \rightarrow \tilde{b}b \rightarrow \chi_2^0 bb \rightarrow \tilde{l}bbl \rightarrow \chi_1^0 bbl$$

For each event set of kinematical equations:

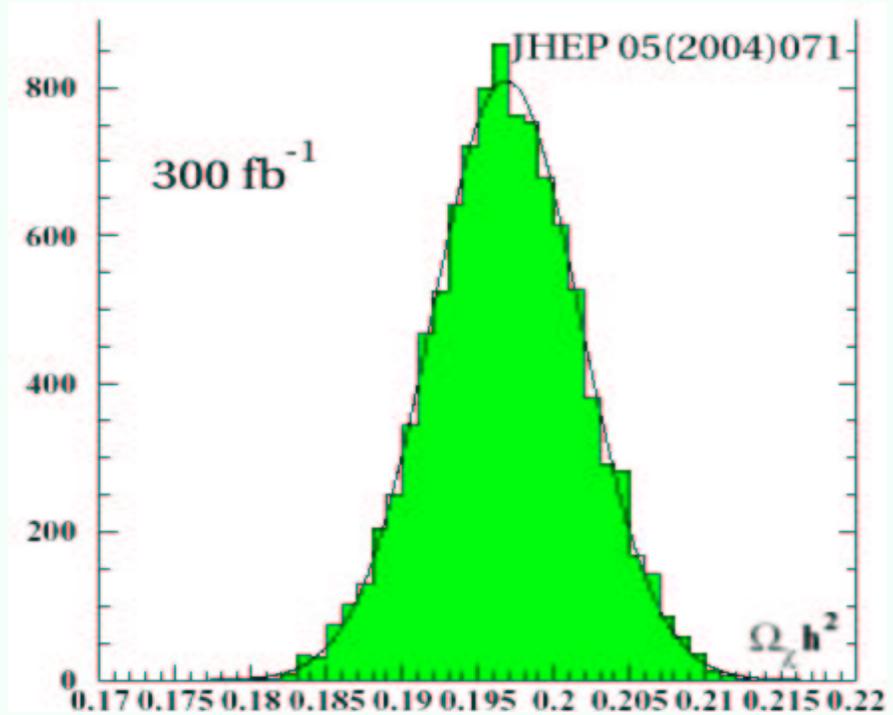
$m_{\chi_1^0}^2 = p_{\chi_1^0}^2$	● 5 equations
$m_{\tilde{l}}^2 = (p_{\chi_1^0} + p_{l_1})^2$	● 9 unknowns:
$m_{\chi_2^0}^2 = (p_{\chi_1^0} + p_{l_1} + p_{l_2})^2$	● 5 masses
$m_{\tilde{b}}^2 = (p_{\chi_1^0} + p_{l_1} + p_{l_2} + p_{b_1})^2$	● χ_1^0 4-momentum
$m_{\tilde{g}}^2 = (p_{\chi_1^0} + p_{l_1} + p_{l_2} + p_{b_1} + p_{b_2})^2$	

- The masses are the same for all the events!
- With 5 events 25 equation and 25 unknowns ($5(m) + 5 \times 4 (p)$)
- Use all events (not only near edge): increase statistics
- Don't need to know $m_{\chi_1^0}$

- Model-dependent mSUGRA fit to all measurements (like SM at LEP)
- Scan on a grid of SUSY models, get best model by χ^2 comparison
- Deduce LSP mass and relic density $\Omega_\chi h^2$
- Generate many simulated experiments for all measurements and expected uncertainties



- Tail at 100 fb^{-1} \tilde{b}_1, \tilde{b}_2 not resolved



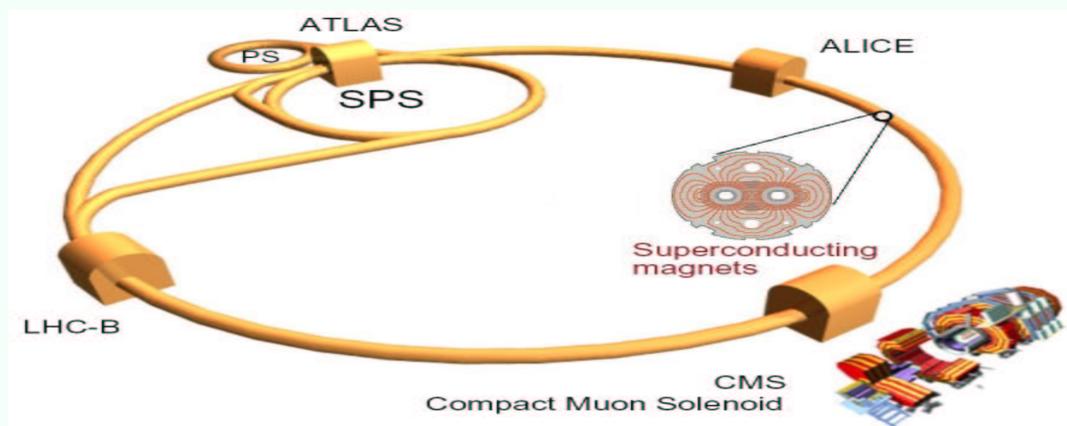
- Uncertainty $\sigma(\Omega_\chi h^2)/\Omega_\chi h^2 \sim 3\%$
- $\Omega_\chi h^2 = 0.197 \pm 0.003$

Result presented for **B**: most favourably!

- Bulk region:
 - **G** like **B** but $\text{BR}(\chi_2^0 \rightarrow \chi_1^0 ll) \approx 2.3\%$ (was 16.4%) and σ smaller (~ 8 pb vs ~ 58 pb)
 - Borderline even with final LHC statistics
 - **I** $\text{BR}(\chi_2^0 \rightarrow \chi_1^0 \tau\tau) \sim 100\%$
 - Need work to fully exploit τ 's final state
- Coannihilation tail
 - **D,J,L**: $\text{BR}(\chi_2^0 \rightarrow \chi_1^0 \tau\tau) \sim 100\%$
 - **A,C,H**: similar to **B** but $10 \div 100 \times \int \mathcal{L} dt$
- Focus points and Rapid annihilation funnels
 - $M_{\tilde{g}, \tilde{q}}$ very high
 - (Focus Points) Even heavier with with higher M_{top} !
($M_{top} = 180$ GeV D0)
 - Only lighter scalar Higgs h visible at LHC

- LHC will discover SUSY (if exist)
- Most favorable mSUGRA in Bulk region: with $\sim 100 \text{ fb}^{-1}$ sound estimate for Dark Matter
- Less favorable scenarios: need full statistics ($\sim 300 \text{ fb}^{-1}$) and more work for $\chi_1^0 \rightarrow \tau' s$ decay
- Focus points, Rapid annihilation funnels: SUSY spectroscopy \sim hopeless! sparticles too heavy
- Caveat: mSUGRA model, R-parity conserved, ...

- LHC will discover SUSY (if exist)
- Most favorable mSUGRA in Bulk region: with $\sim 100 \text{ fb}^{-1}$ sound estimate for Dark Matter
- Less favorable scenarios: need full statistics ($\sim 300 \text{ fb}^{-1}$) and more work for $\chi_1^0 \rightarrow \tau's$ decay
- Focus points, Rapid annihilation funnels: SUSY spectroscopy \sim hopeless! sparticles too heavy
- **Caveat: mSUGRA model, R-parity conserved, ...**
- With a bit of luck, it could be that looking inside the largest microscope (LHC)



- LHC will discover SUSY (if exist)
- Most favorable mSUGRA in Bulk region: with $\sim 100 \text{ fb}^{-1}$ sound estimate for Dark Matter
- Less favorable scenarios: need full statistics ($\sim 300 \text{ fb}^{-1}$) and more work for $\chi_1^0 \rightarrow \tau' s$ decay
- Focus points, Rapid annihilation funnels: SUSY spectroscopy \sim hopeless! sparticles too heavy
- Caveat: mSUGRA model, R-parity conserved, ...
- With a bit of luck, it could be that looking inside the largest microscope (LHC) we will see the Universe!

