

# PONTECORVO'S IDEA (III)



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IFIC-CSIC-U.VALENCIA  
CERN SUMMER STUDENT LECTURES  
2008





Бруно Понтекорво

Bruno Pontecorvo was one of the “Ragazzi de via Panisperma” together with another genius, Ettore Majorana.

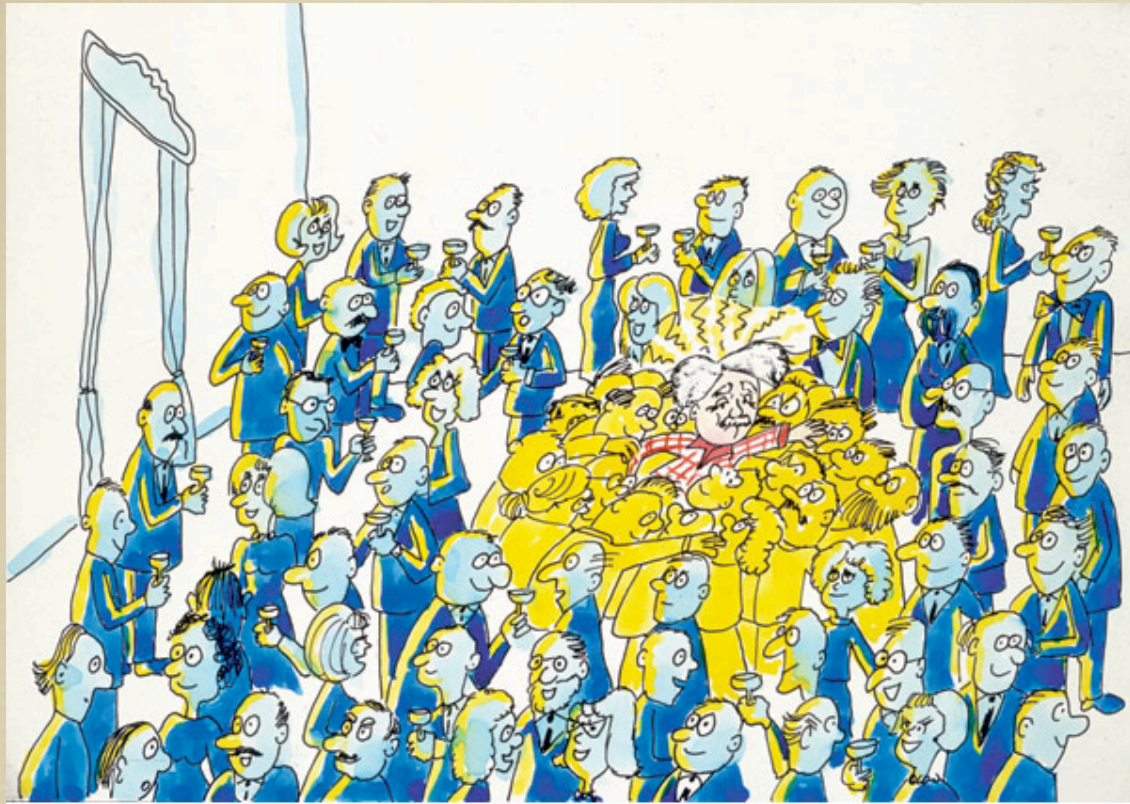
Assistant of Fermi, he made numerous contributions to physics. Convinced socialist he is also famous for his move in 1950 to the USSR.

Among his scientific ideas:

- 1) The delayed coincidence that made possible Reines & Cowan experiment (discovery of  $\bar{\nu}_e$ )
- 2) The prediction that electron neutrinos are not the same than muon neutrinos
- 3) The mechanism of neutrino oscillations



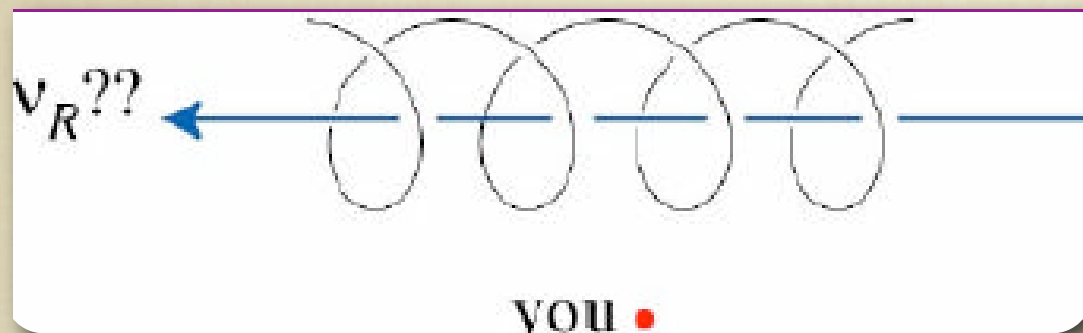
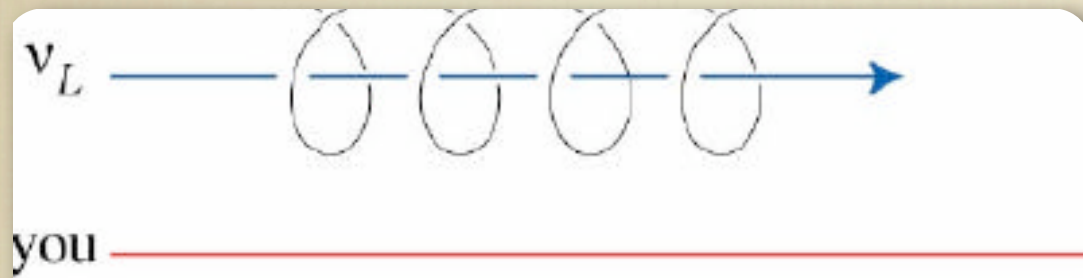
# FERMION MASSES & HIGGS BOSON



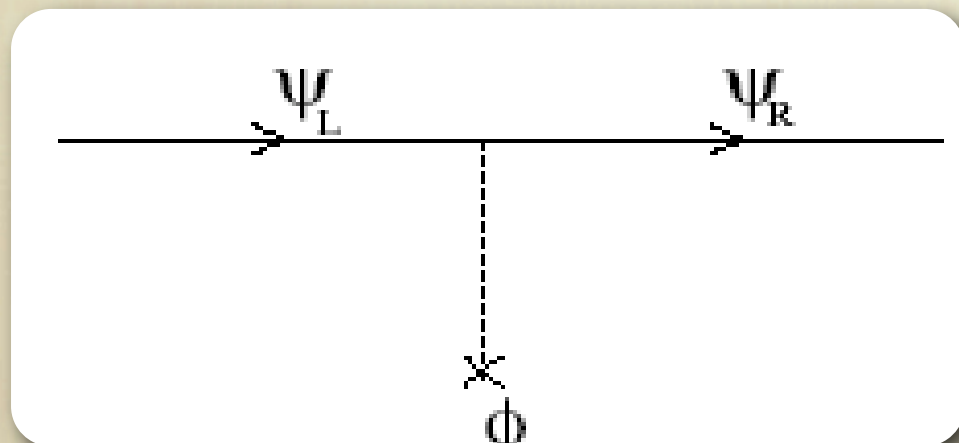
The behaviour of physicists in a crowded social event at a conference is an analogy for the Higgs mechanism, as proposed by David Miller (University College London). The physicists represent a non-trivial medium permeating space. In the upper panel, the physicists cluster around a famous scientist who enters the room, slowing the scientist's progress. In much the same way, a particle passing through the Higgs–Brout–Englert field slows down and acquires a mass. In the lower panel, a rumor propagates. This is an excitation of the medium — the group of physicists — itself, forming a body with a large mass; this is analogous to the formation of a Higgs boson.



# LEFT AND RIGHT HANDED PARTICLES



IF A PARTICLE IS MASSIVE LEFT AND RIGHT STATES MUST EXIST

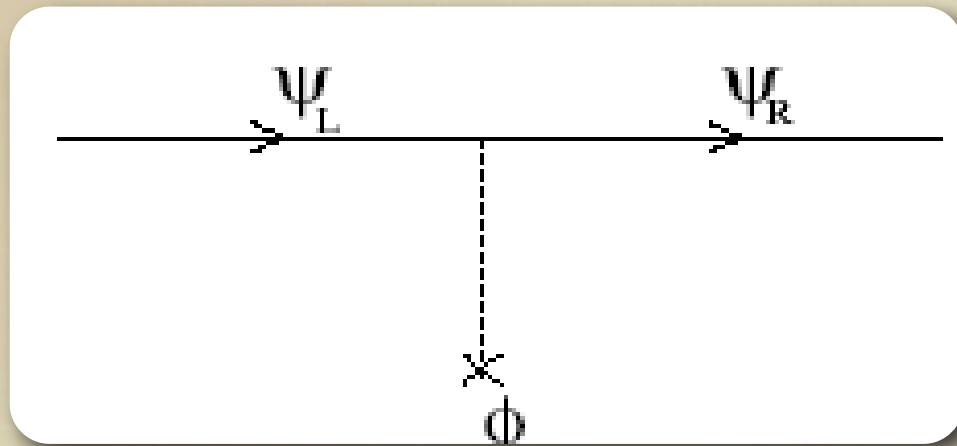


A (ELECTRICALLY) CHARGED FERMION SUCH AS THE ELECTRON HAS LEFT AND RIGHT STATES FOR PARTICLE AND ANTIPARTICLE (WHICH ARE DISTINCT BY ELECTRIC CHARGE)

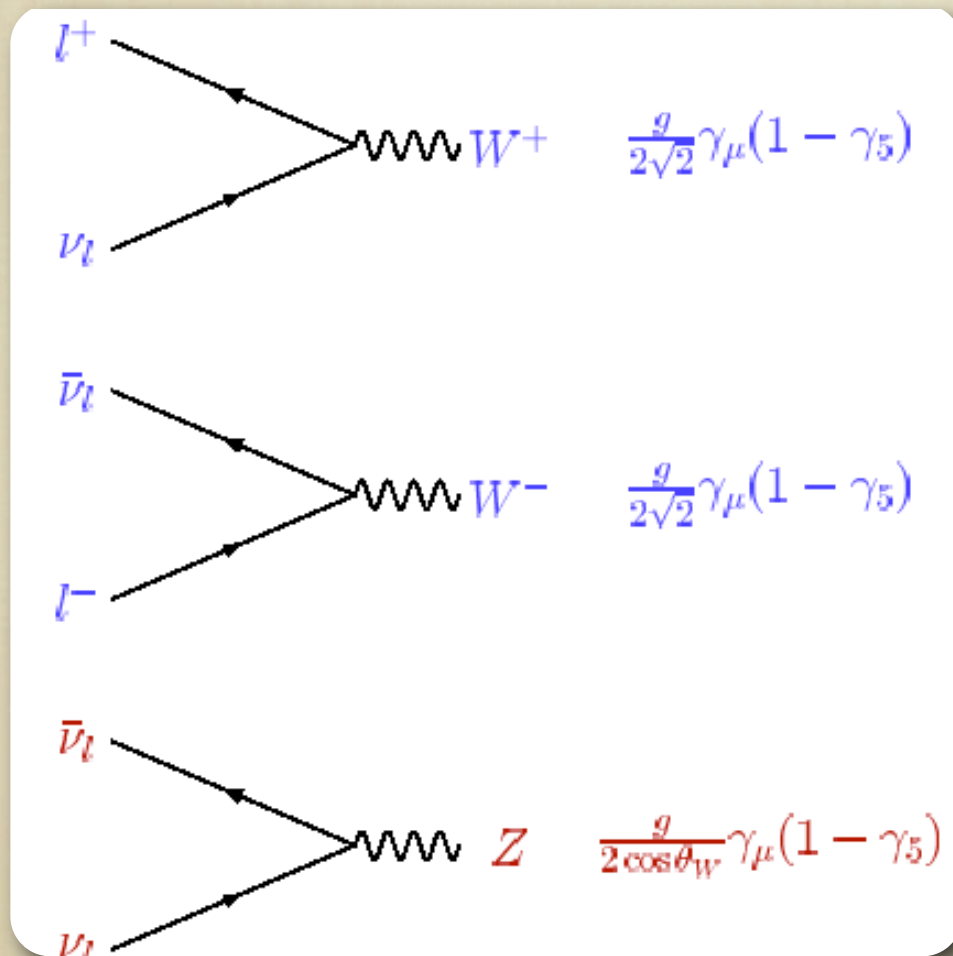
CHARGED FERMIONS COUPLE LEFT-RIGHT STATES TO A SCALAR (THE HIGGS) TO GENERATE MASSES



# MASS AND FLAVOR EIGENSTATES



**MASS STATES: OBJECTS THAT COUPLE TO HIGGS.**

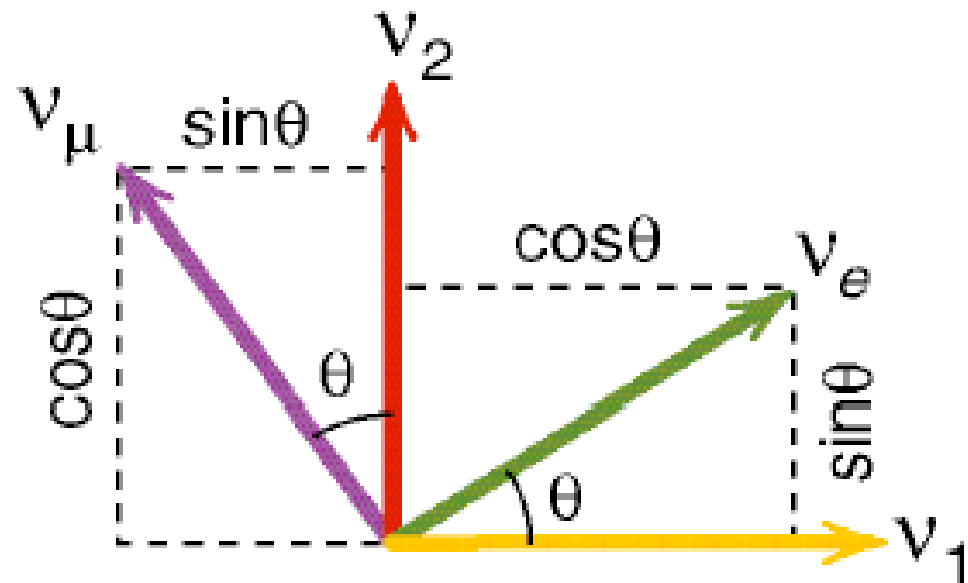


**WEAK STATES: OBJECTS THAT COUPLE TO WEAK BOSONS.**

**ARE THOSE TWO TYPES OF OBJECTS IDENTICAL?**

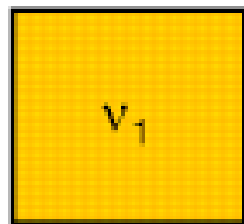


# MIXING



Mass states

First



Second

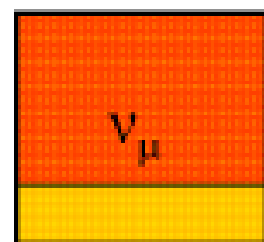


Weak states

First



Second

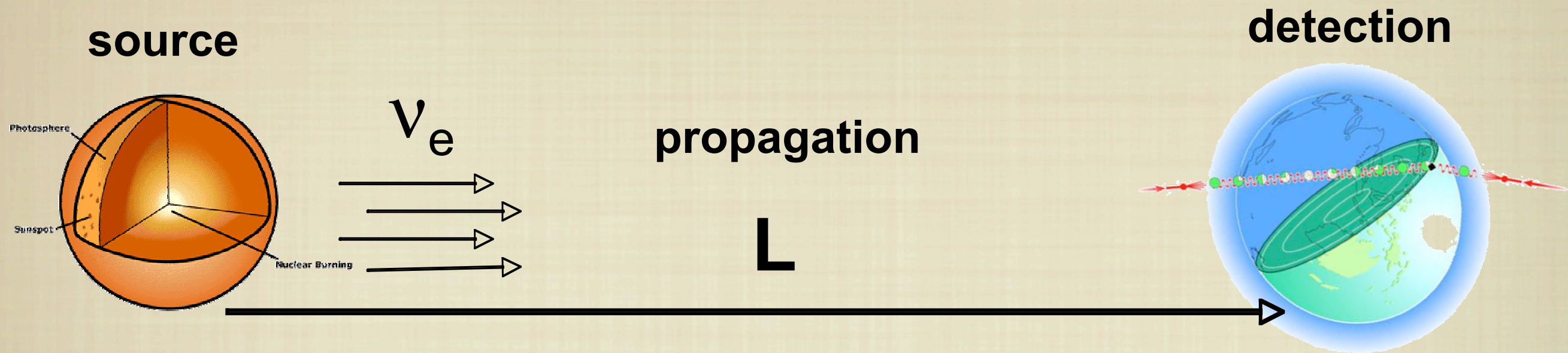


$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} = U \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

Mass and weak states are distinct and connected by means of an unitary transformation, the PMNS mixing matrix, which depends of a single parameter, the mixing angle  $\theta$



# OSCILLATIONS



The weak interaction produces neutrinos of a given flavor

$$\begin{aligned}
 |\nu(x_0)\rangle &= |\nu_e\rangle \\
 &= c|\nu_1\rangle + s|\nu_2\rangle
 \end{aligned}$$

The mass eigenstates Propagate at different velocities

$$\begin{aligned}
 |\nu(x)\rangle &= c|\nu_1\rangle e^{i(Et - \bar{k}_1 \bar{x})} \\
 &\quad + s|\nu_2\rangle e^{i(Et - \bar{k}_2 \bar{x})}
 \end{aligned}$$

Detection again via weak interaction

$$\begin{aligned}
 \nu_\mu N &\rightarrow \mu^- X \\
 \nu_e N &\rightarrow e^- X
 \end{aligned}$$

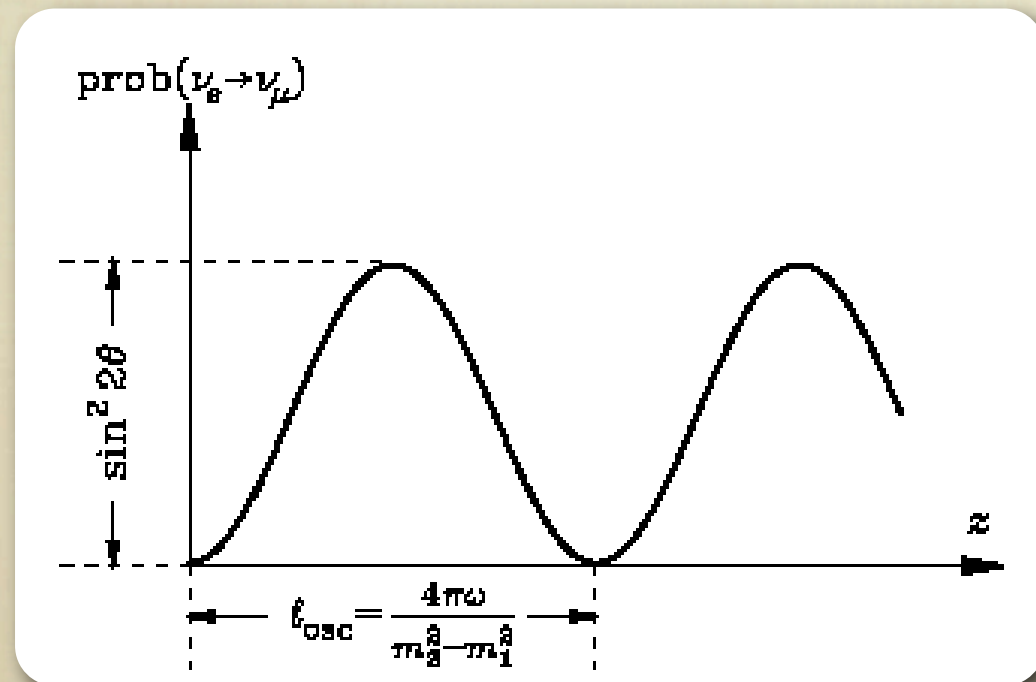
$$P(\nu_e \rightarrow \nu_\mu) = |\langle \nu_\mu | \nu(t) \rangle|^2$$



# OSCILLATION PROBABILITY

$$P_{\nu_e \rightarrow \nu_\mu}(L) = \sin^2(2\theta) \sin^2\left(1.27 \frac{\Delta m^2 (\text{eV}^2)}{E(\text{GeV})} L(\text{km})\right)$$

$$P_{\nu_e \rightarrow \nu_e}(L) = 1 - P_{\nu_e \rightarrow \nu_\mu}(L)$$

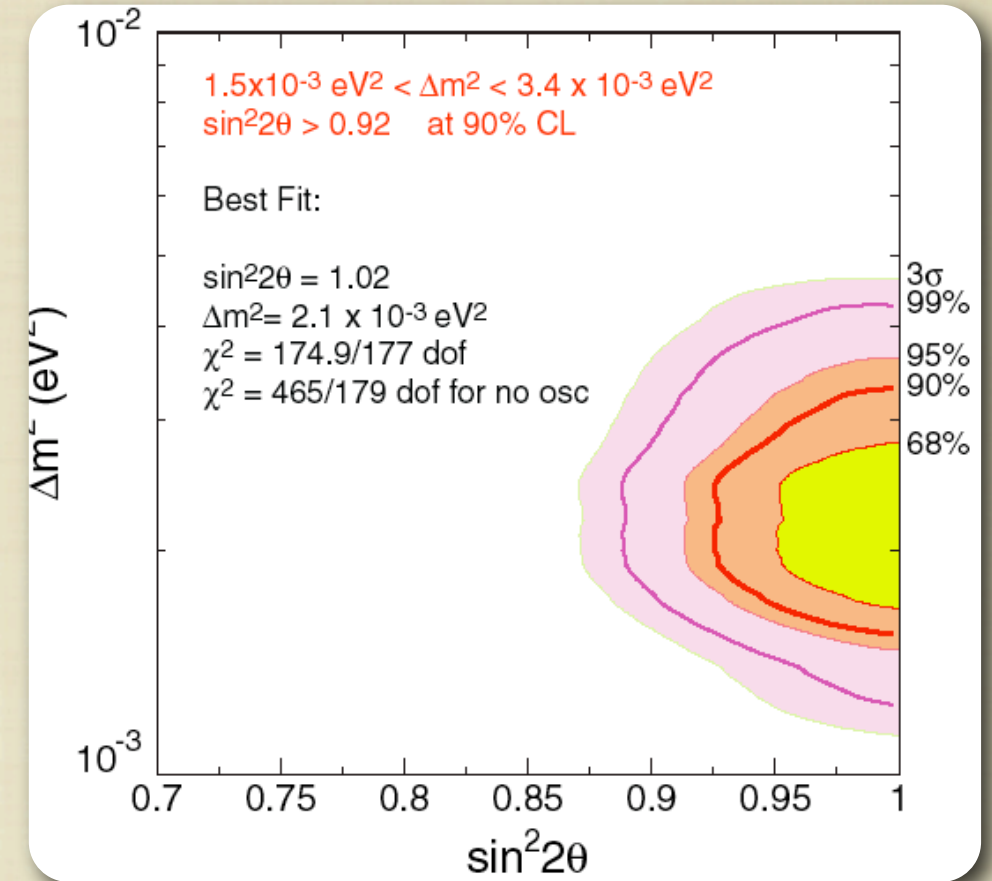
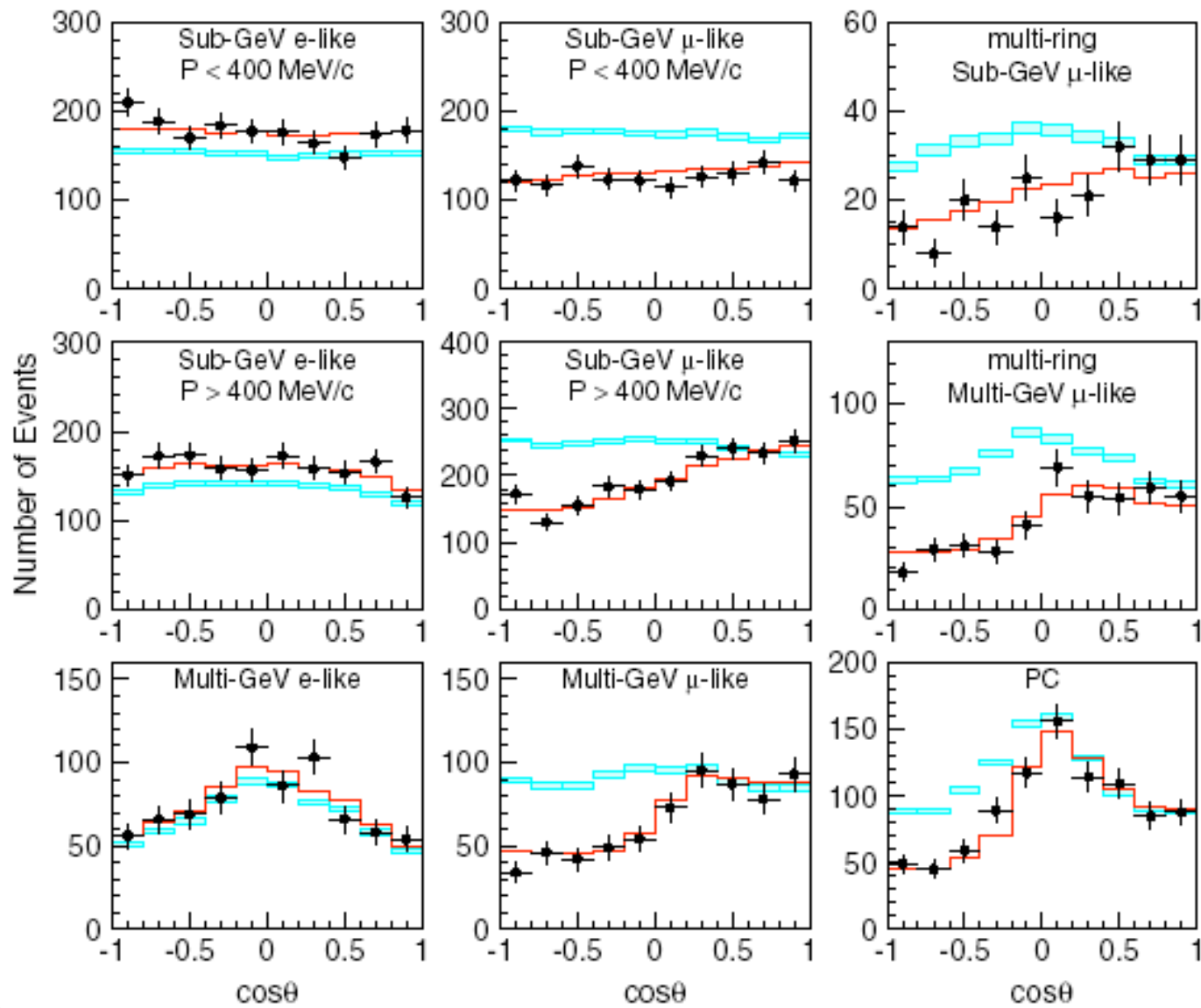


$$\Delta m^2 = m_2^2 - m_1^2$$

$$L_{\text{osc}}(\text{Km}) \approx \frac{E(\text{GeV})}{1.27 \Delta m^2 (\text{eV}^2)}$$



# ATMOSPHERIC NEUTRINO PROBLEM



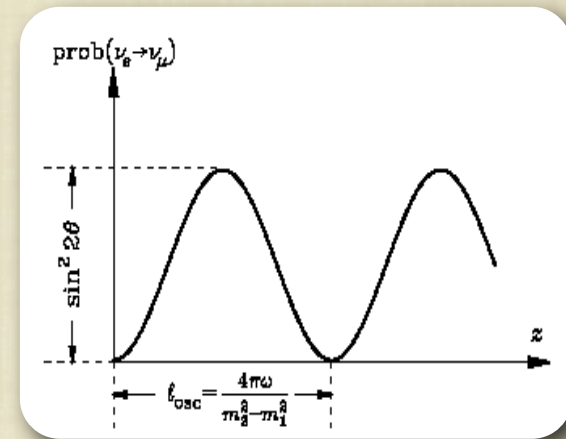
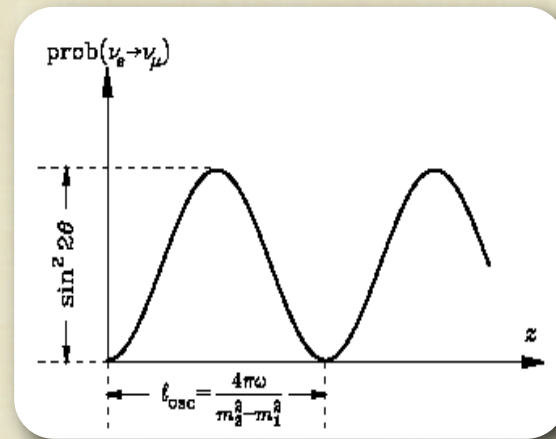
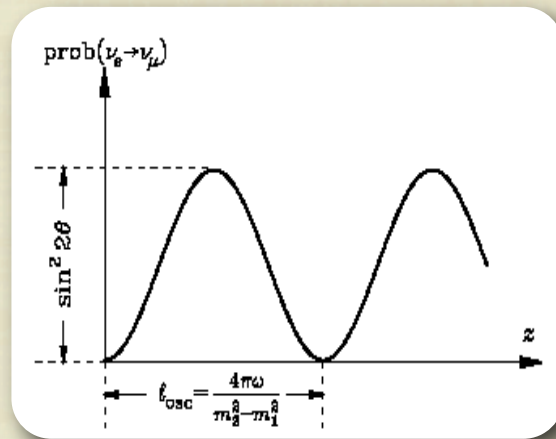
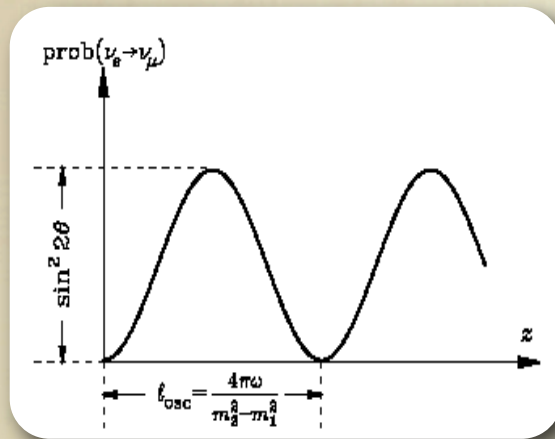
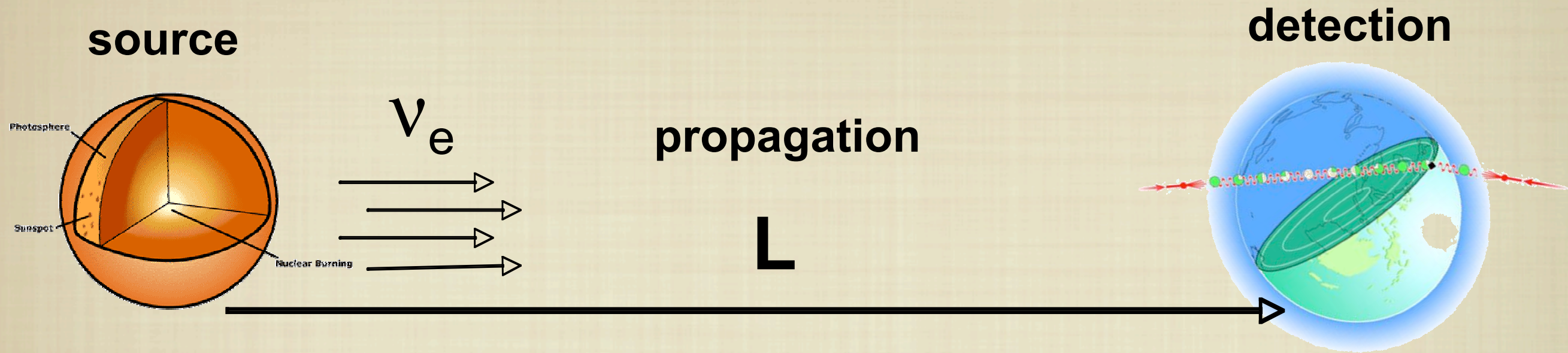
$\nu_\mu \rightarrow \nu_\tau$  oscillations

$$\Delta m^2 = 2.1 \cdot 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta \approx 1$$



# SOLAR NEUTRINOS & FINE TUNING

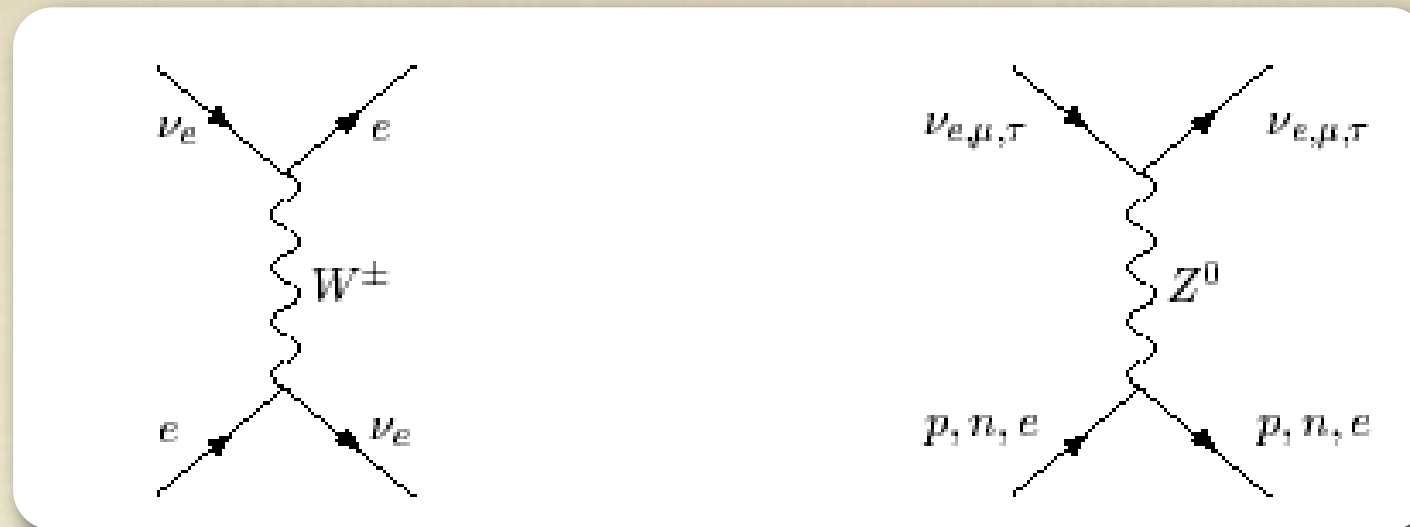


$$L_{osc} (Km) \approx \frac{E(GeV)}{1.27 \Delta m^2 (eV^2)}$$

$\Delta m^2$  and solar neutrino energy must be chosen to provide a  $L_{osc}$  that matches exactly distance from sun to earth.



# NEUTRINO OSCILLATION IN MATTER



Of all the three neutrinos,  $\nu_\mu$  and  $\nu_\tau$  interact via neutral currents with  $p$ ,  $n$  and  $e$ . But  $\nu_e$  is the only neutrino that can interact via CC and NC with the electrons of the medium.

This fact changes the oscillation probability with neutrinos propagate in dense matter. There can be a resonant enhancement of the oscillation probability. The Mikheyev-Smirnov-Wolfenstein (MSW) effect.

$P_{\text{osc}}^{\text{matter}}$  can be large (1) even if mixing angle in vacuum is small.

In practice this implies that (if MSW is at work)  $\nu_e$  can oscillate to  $\nu_\mu, \nu_\tau$  BEFORE exiting the sun

# OSCILLATION PROBABILITY IN MATTER

The probability of oscillation in matter has the same form as in  
vacuum

$$P(\nu_e \rightarrow \nu_\mu) = \sin^2 2\tilde{\theta} \sin^2 \left( 2\pi \frac{L}{\tilde{L}_{osc}} \right), \quad \tilde{L}_{osc} = \frac{2\pi E(\text{GeV})}{1.27 \Delta\tilde{m}^2 (\text{eV}^2)}$$



# MSW RESONANCE

For constant matter density there is an energy such that mixing in matter is maximal independently from the vacuum value.

$$\text{if } \Delta m^2 \cos 2\theta = A = \pm \sqrt{2} E G N_e$$

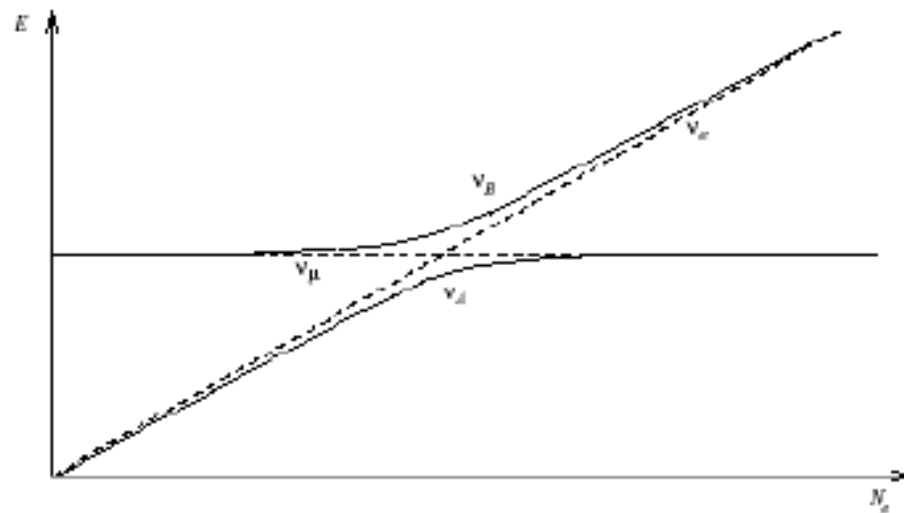
$$\sin^2 2\tilde{\theta} = \frac{(\Delta m^2)^2 \sin^2 2\theta}{(\Delta m^2 \cos 2\theta - A)^2 + (\Delta m^2)^2 \sin^2 2\theta} = 1$$

Thus the probability of neutrino transition in matter can be large even if the mixing angle is small

# ADIABATIC APPROXIMATION

In the sun  $N_e$  is not constant. However if the variation is sufficiently slow the eigenstates of  $H$  change slowly with the density and one can assume that the neutrino remains an eigenstate along the trajectory:

adiabatic approximation



$$\tilde{\nu}_1 = \nu_e \cos \tilde{\theta} + \nu_\mu \sin \tilde{\theta}$$

$$\tilde{\nu}_2 = -\nu_e \sin \tilde{\theta} + \nu_\mu \cos \tilde{\theta}$$

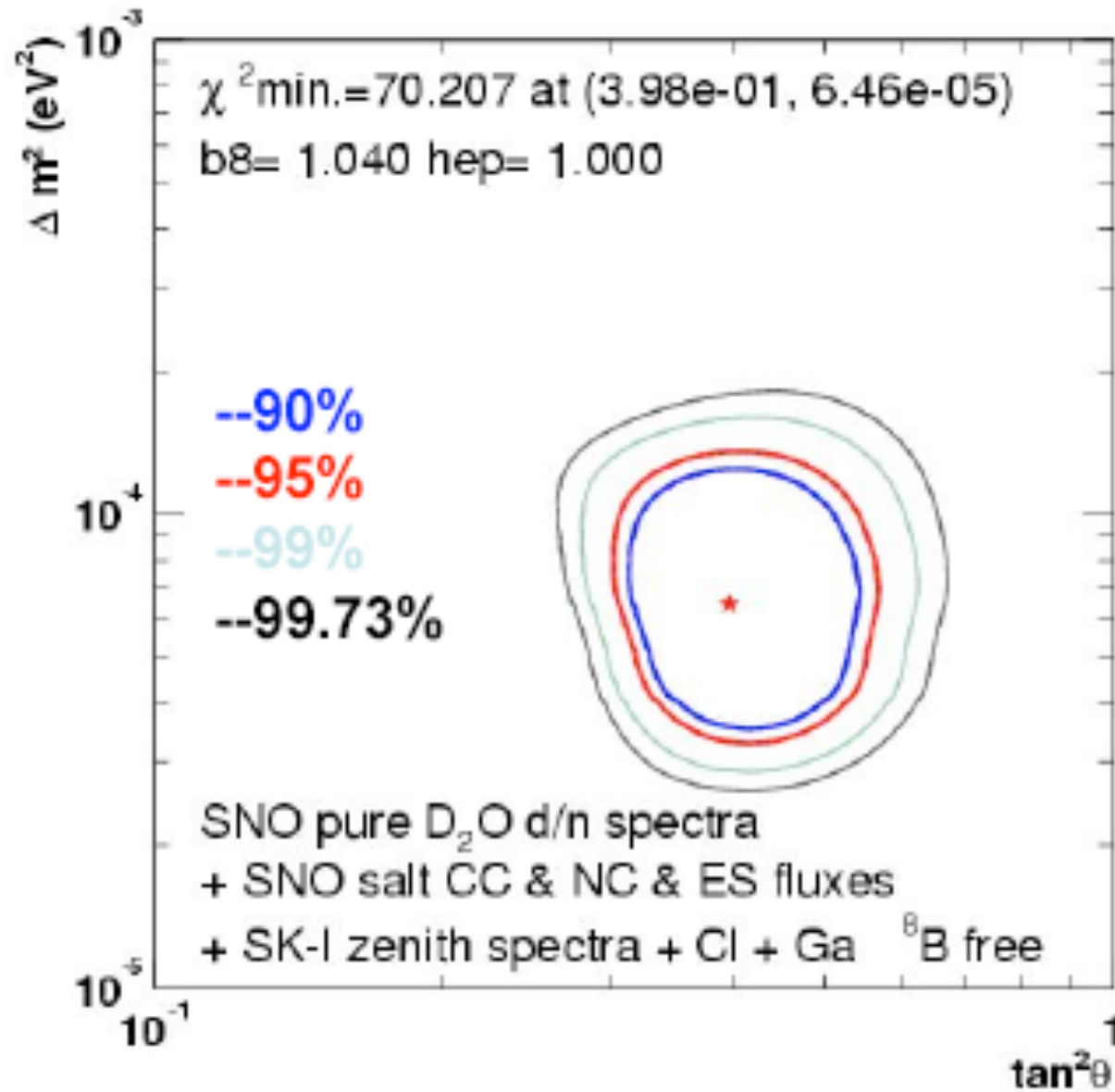
$$x = 0 \quad \text{if } A \gg \Delta m^2 \cos 2\theta \rightarrow \tilde{\theta} \approx \frac{\pi}{2} \Rightarrow \nu_e \approx \tilde{\nu}_2$$

$$x = R_{sun} \quad N_e = 0 \rightarrow \tilde{\theta} \approx \theta \Rightarrow \nu_\mu \approx \tilde{\nu}_2$$

A  $\nu_e$  produced at the sun core is the eigenstate  $\nu_2$  but this eigenstate outside the sun is mostly  $\nu_\mu$ . There is maximum  $\nu_e \rightarrow \nu_\mu$  conversion



# SOLAR OSCILLATIONS

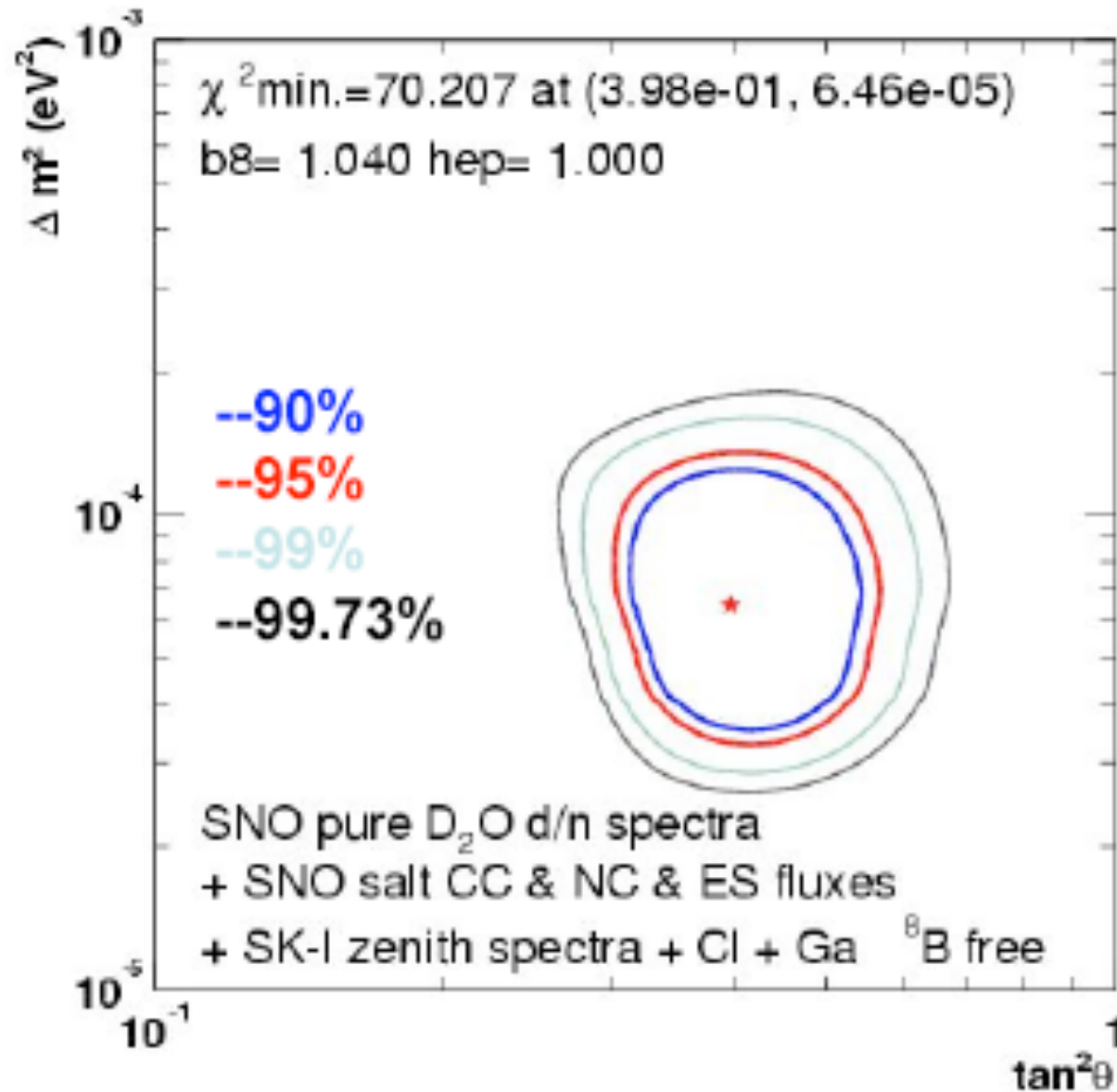


Neutrinos produced at the sun ( $\nu_e$ ) oscillate to other neutrinos via matter-enhanced MSW.

$$\Delta m^2 = 8 \times 10^{-5} \text{ eV}^2$$

$$\theta \approx 30^\circ$$

# SOLAR OSCILLATIONS



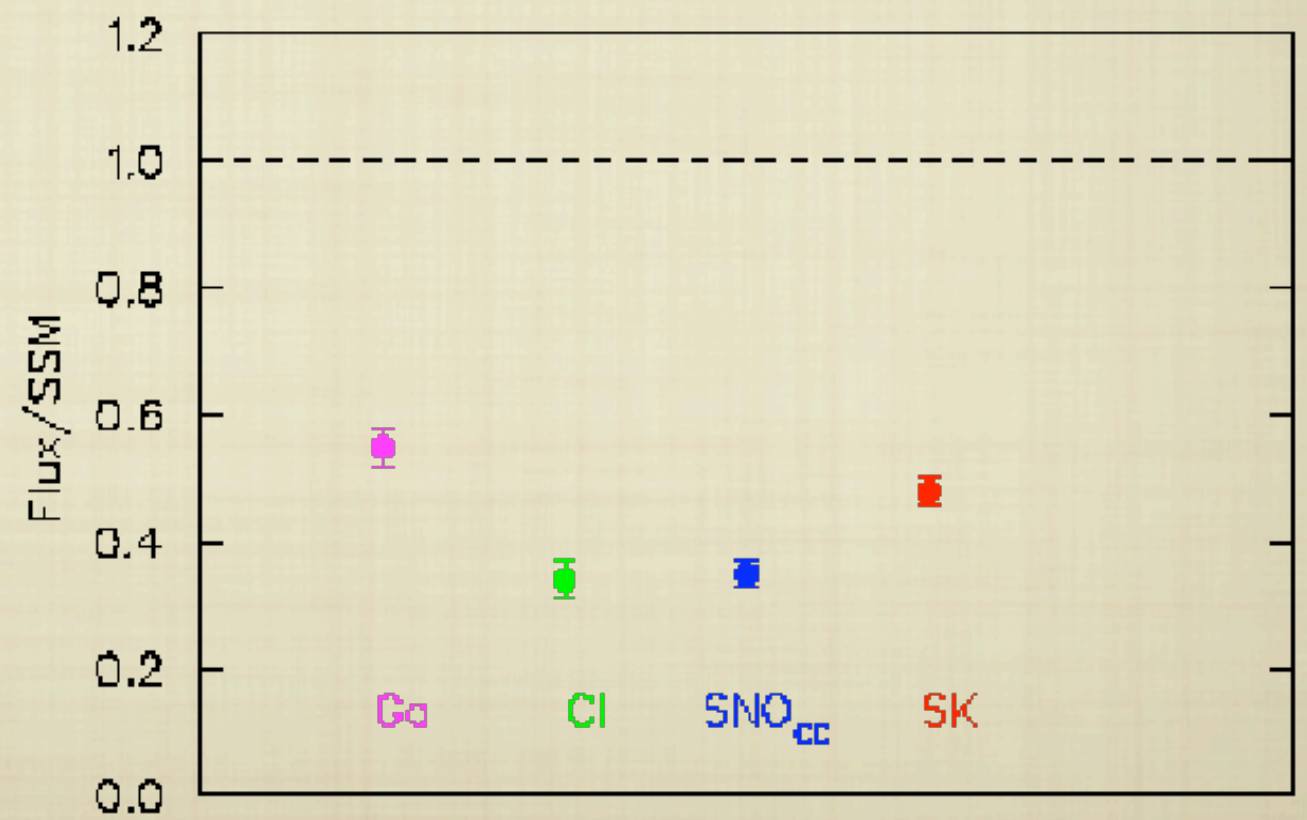
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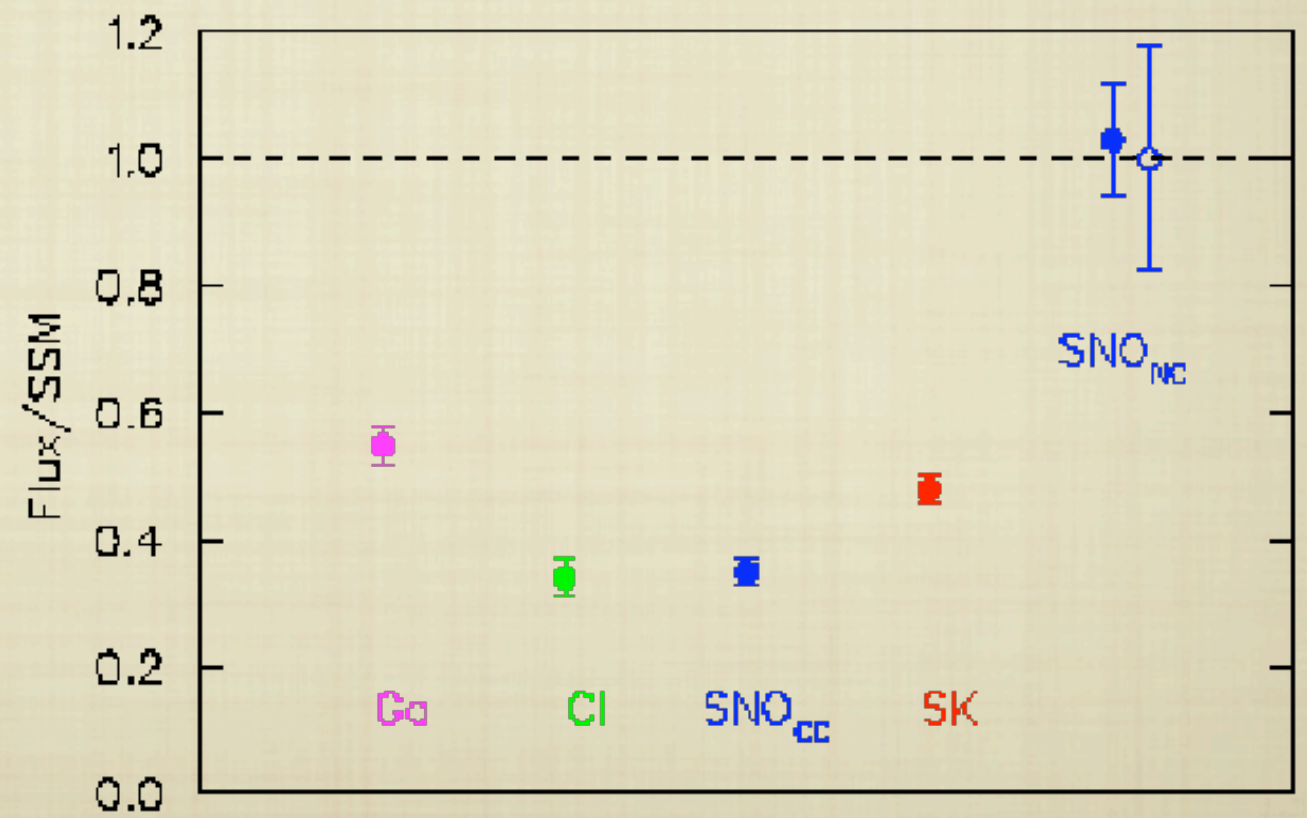
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# SOLAR NEUTRINO OSCILLATIONS



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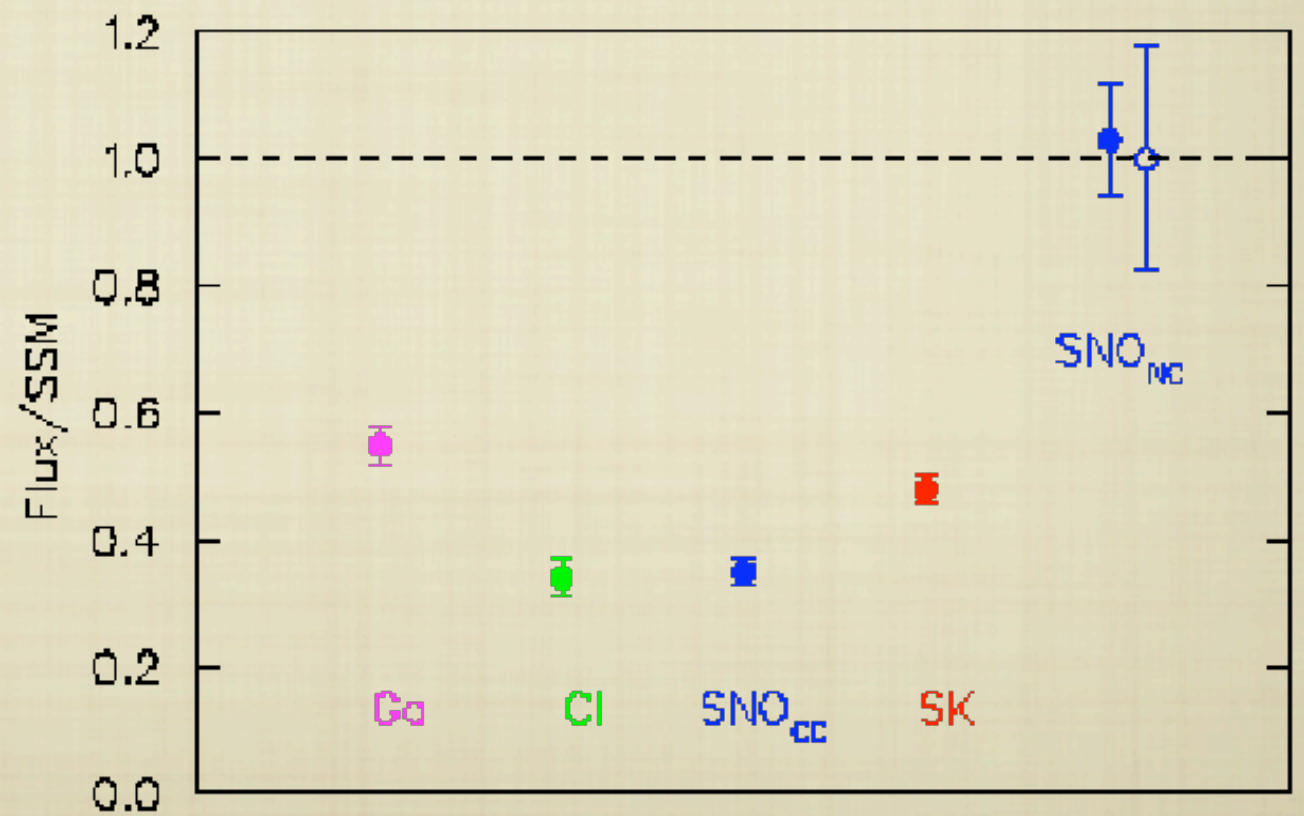
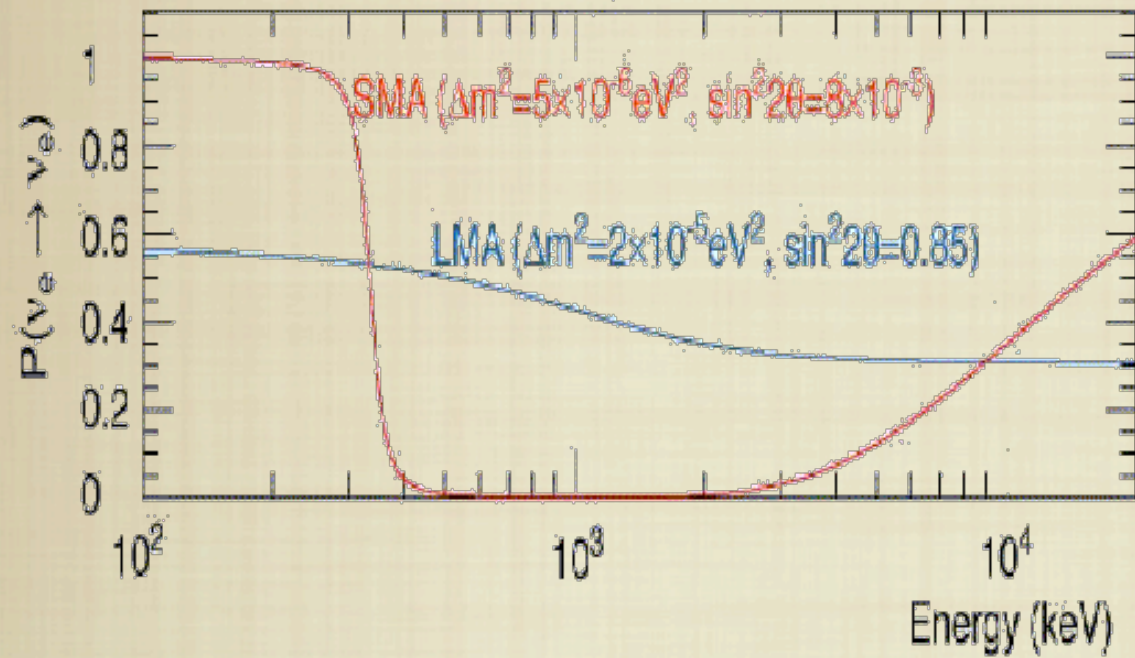




# SOLAR NEUTRINO OSCILLATIONS

## *Matter effect on $\nu_e$ from Sun to Earth*

Survival probability

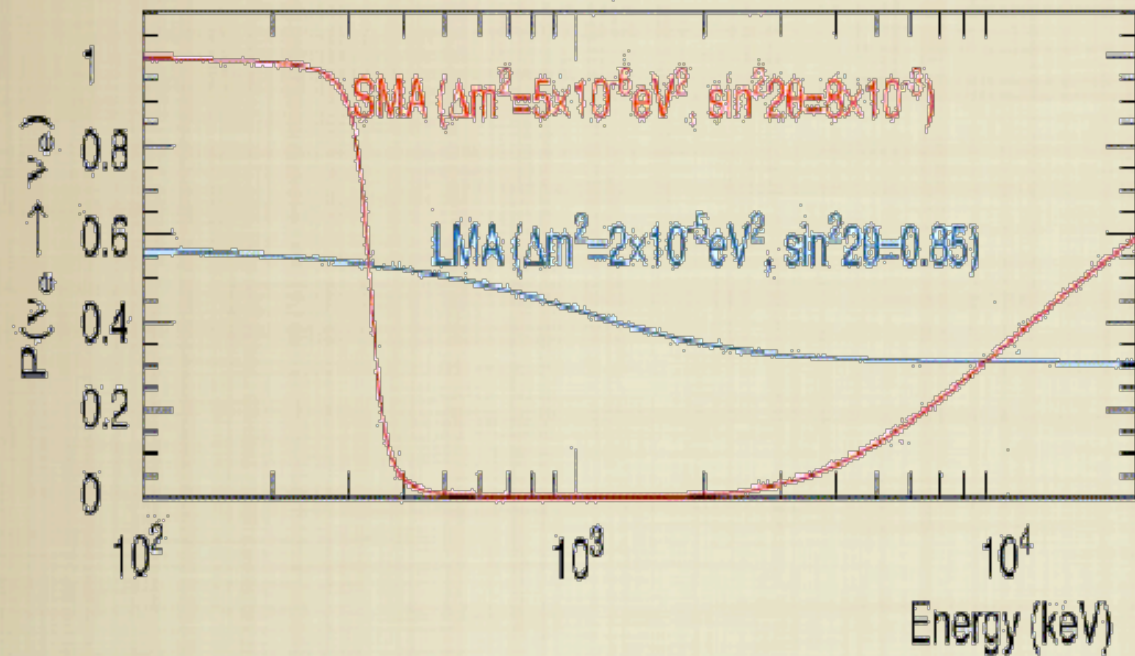




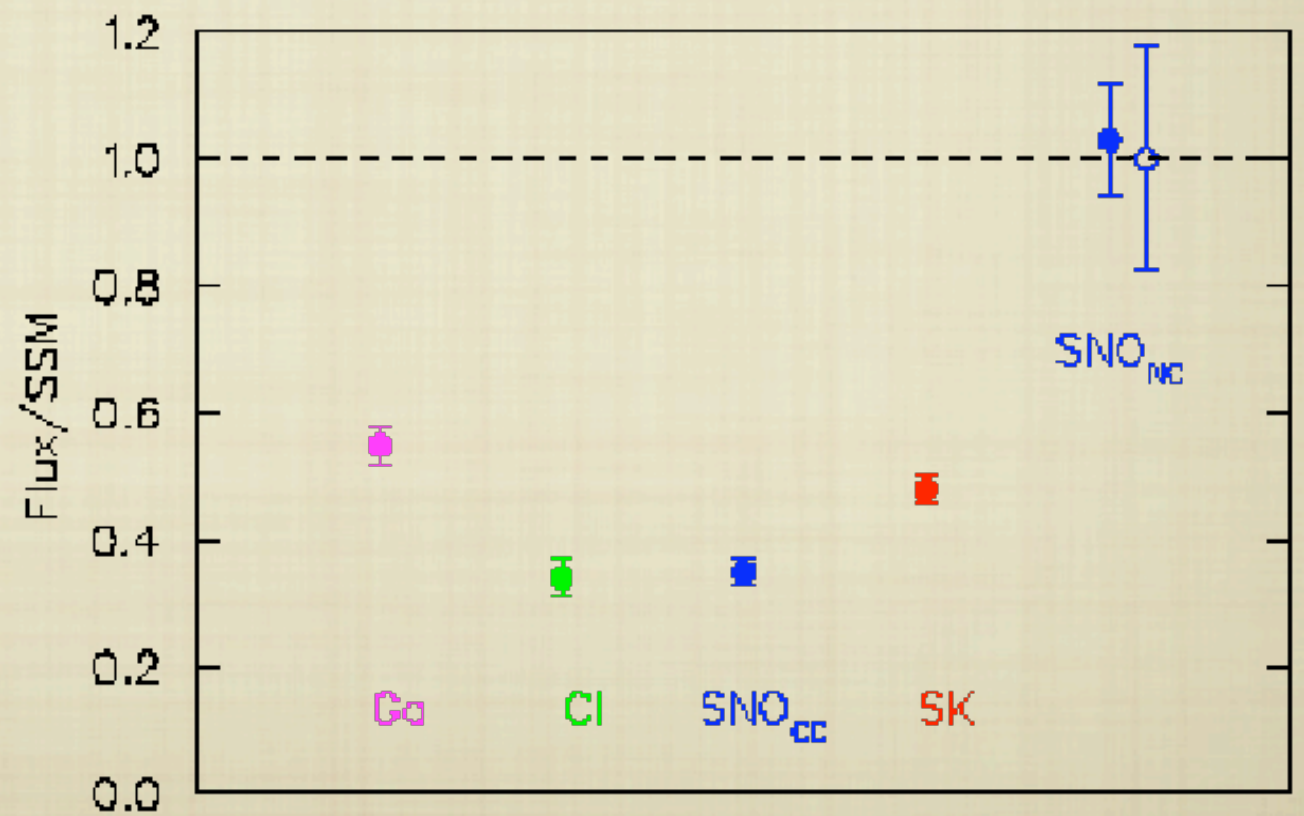
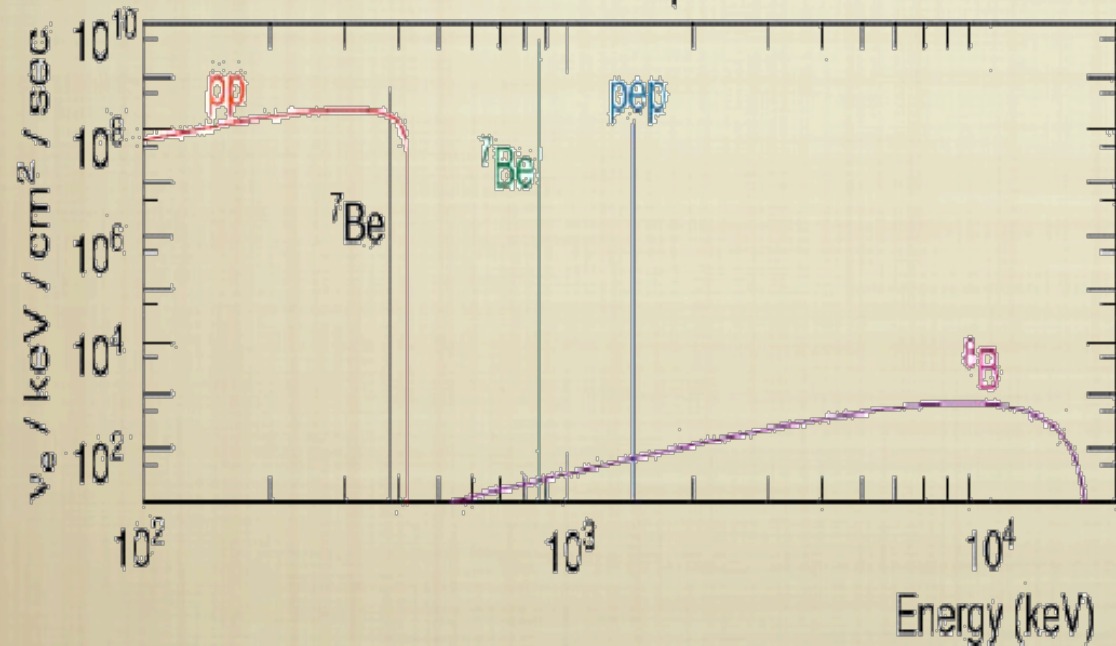
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Solar neutrino spectrum

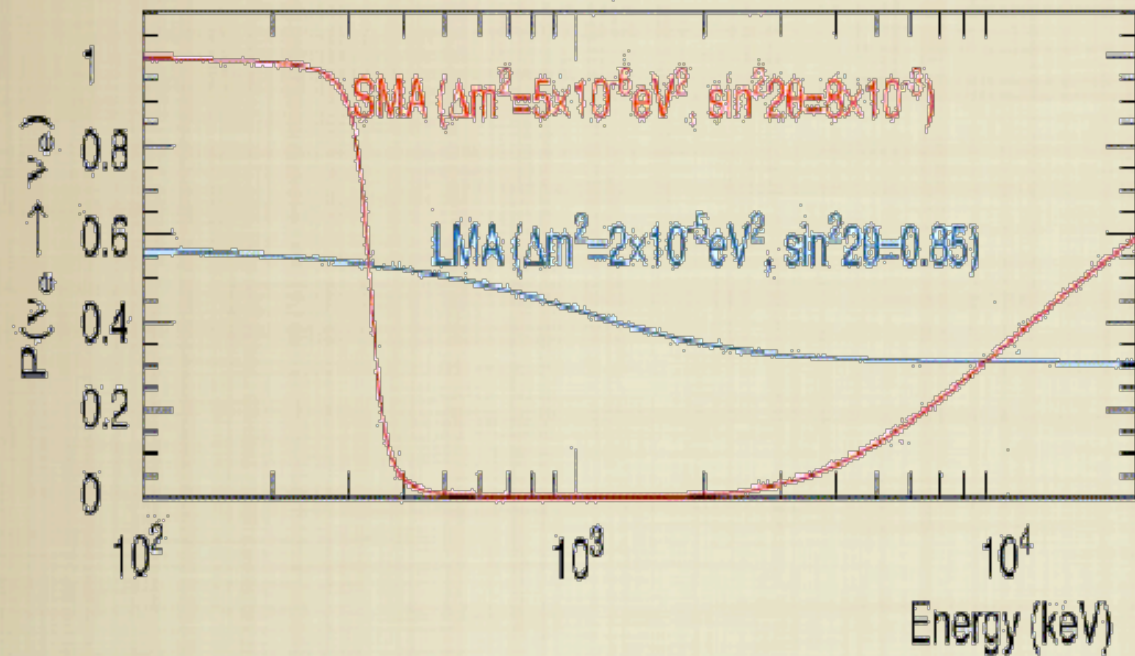




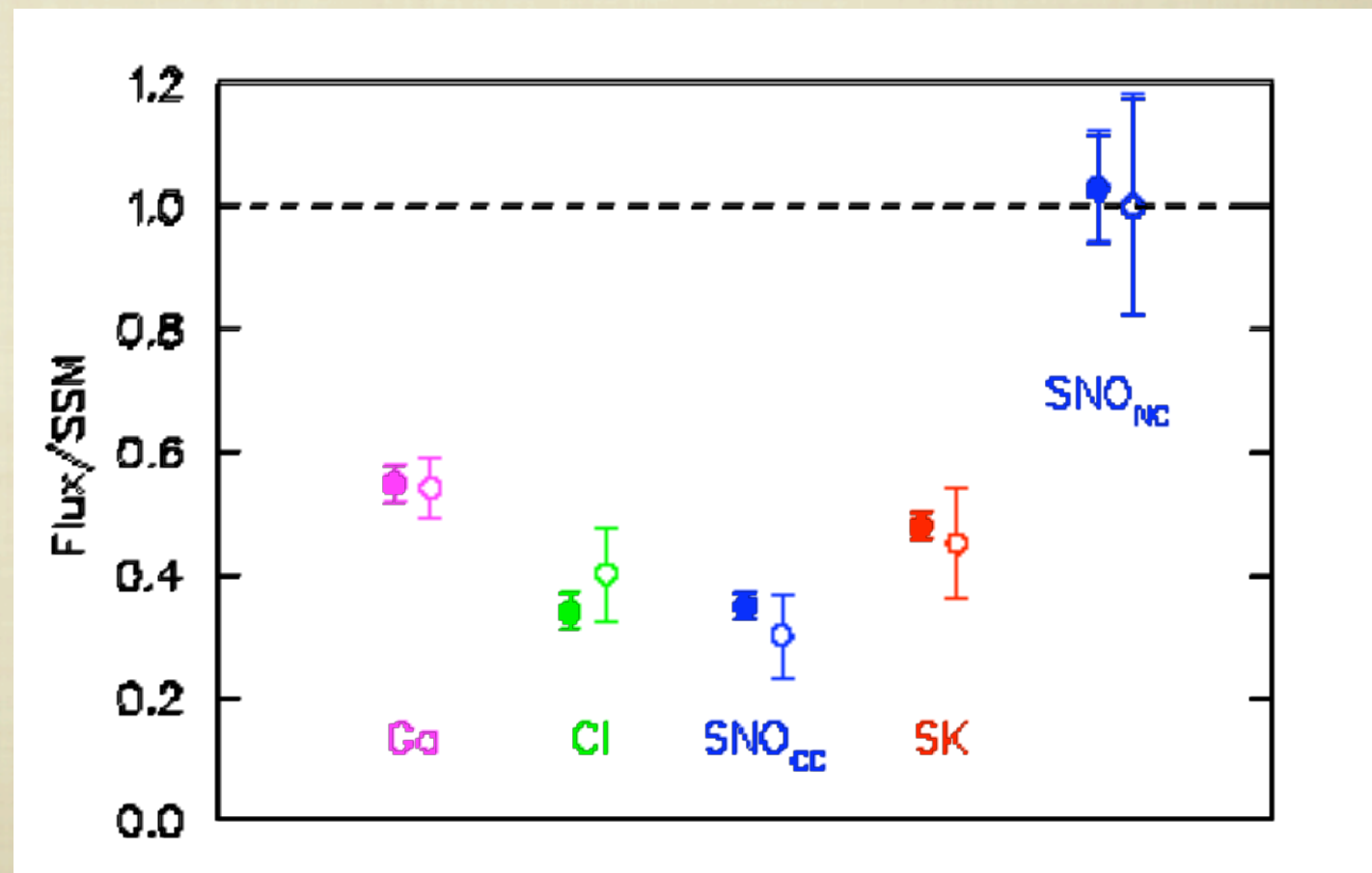
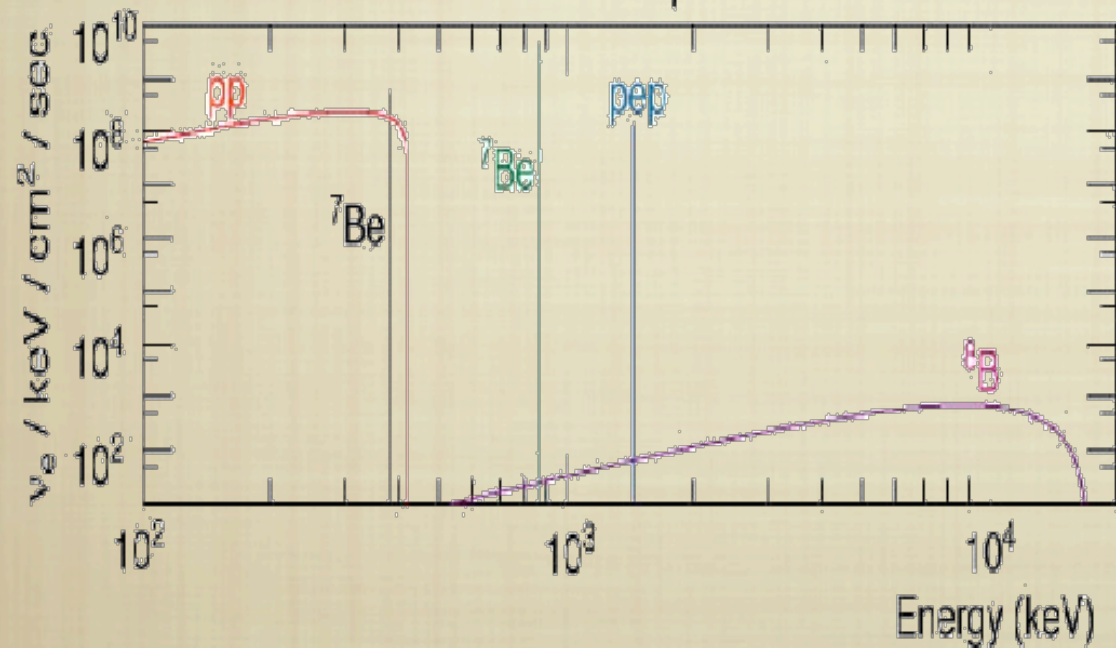
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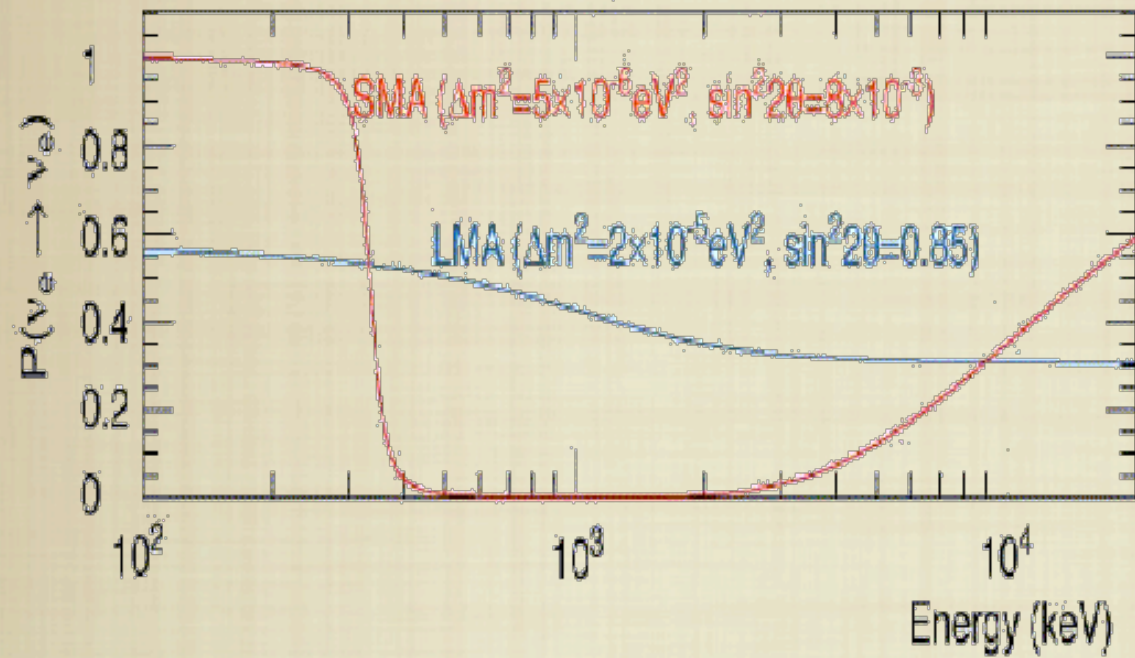
Solar neutrino spectrum



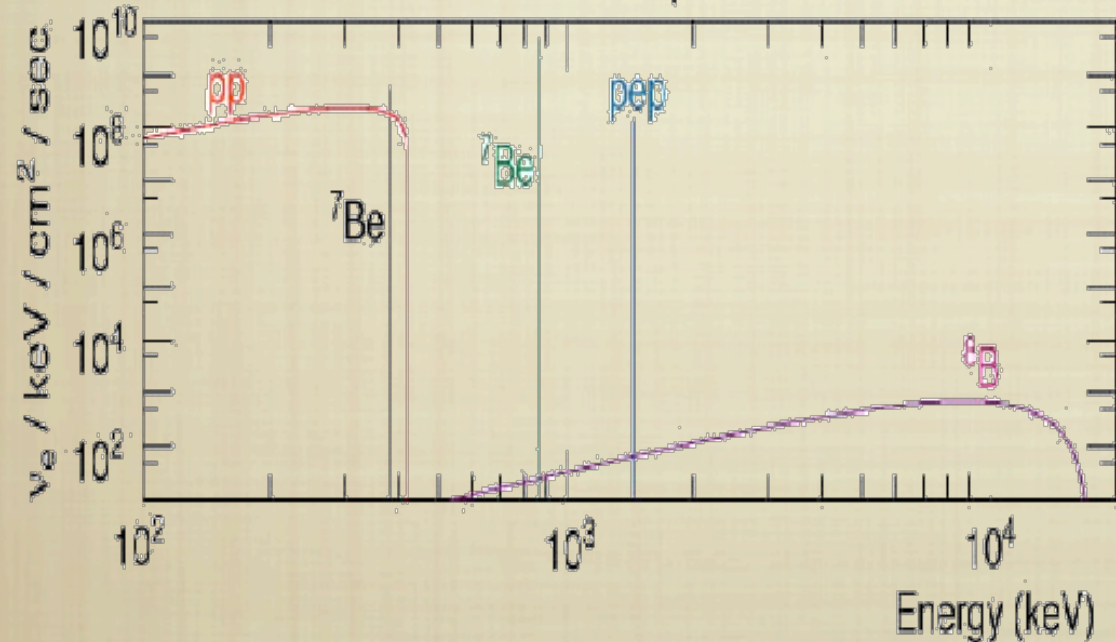
# SOLAR NEUTRINO OSCILLATIONS

## *Matter effect on $\nu_e$ from Sun to Earth*

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Solar neutrino spectrum

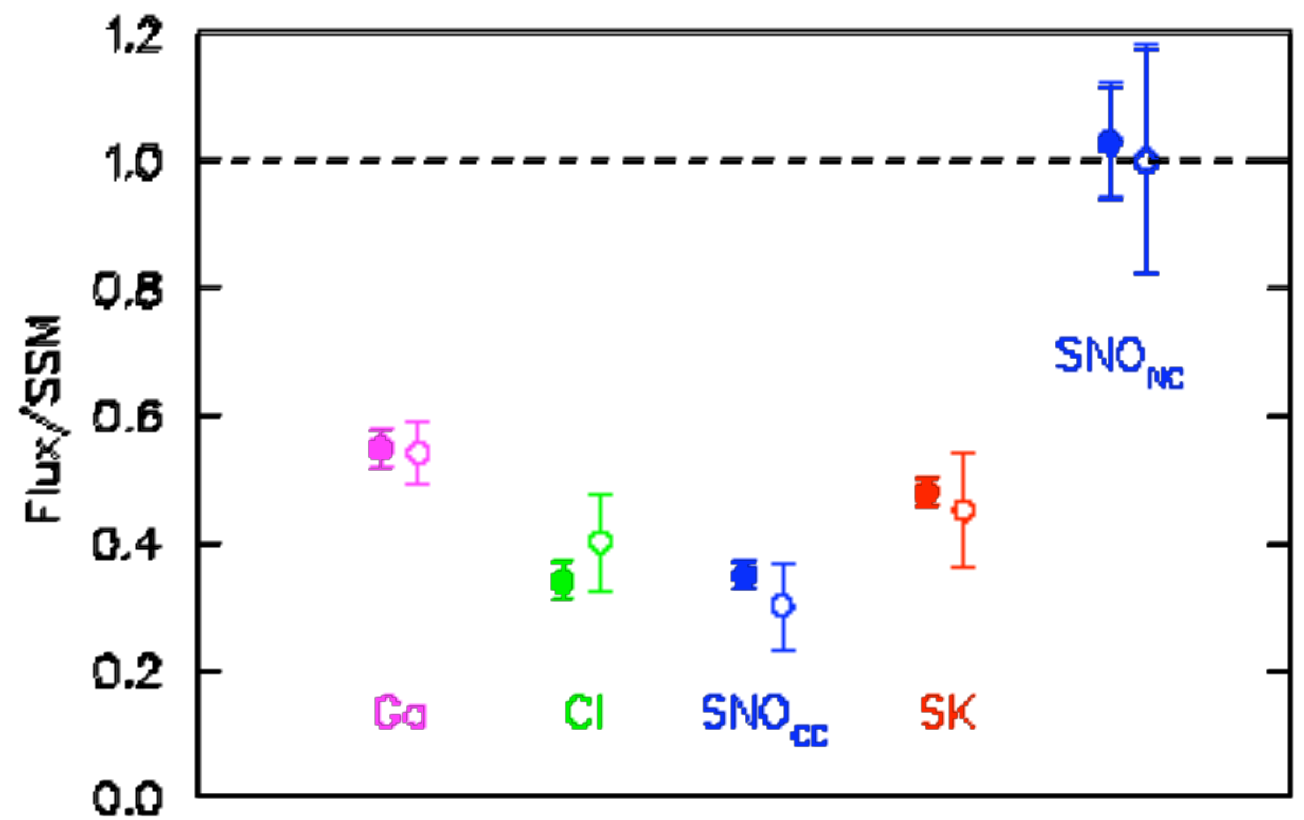


The LMA solar solution

+

matter effects

explain beautifully  
all solar neutrino experiments

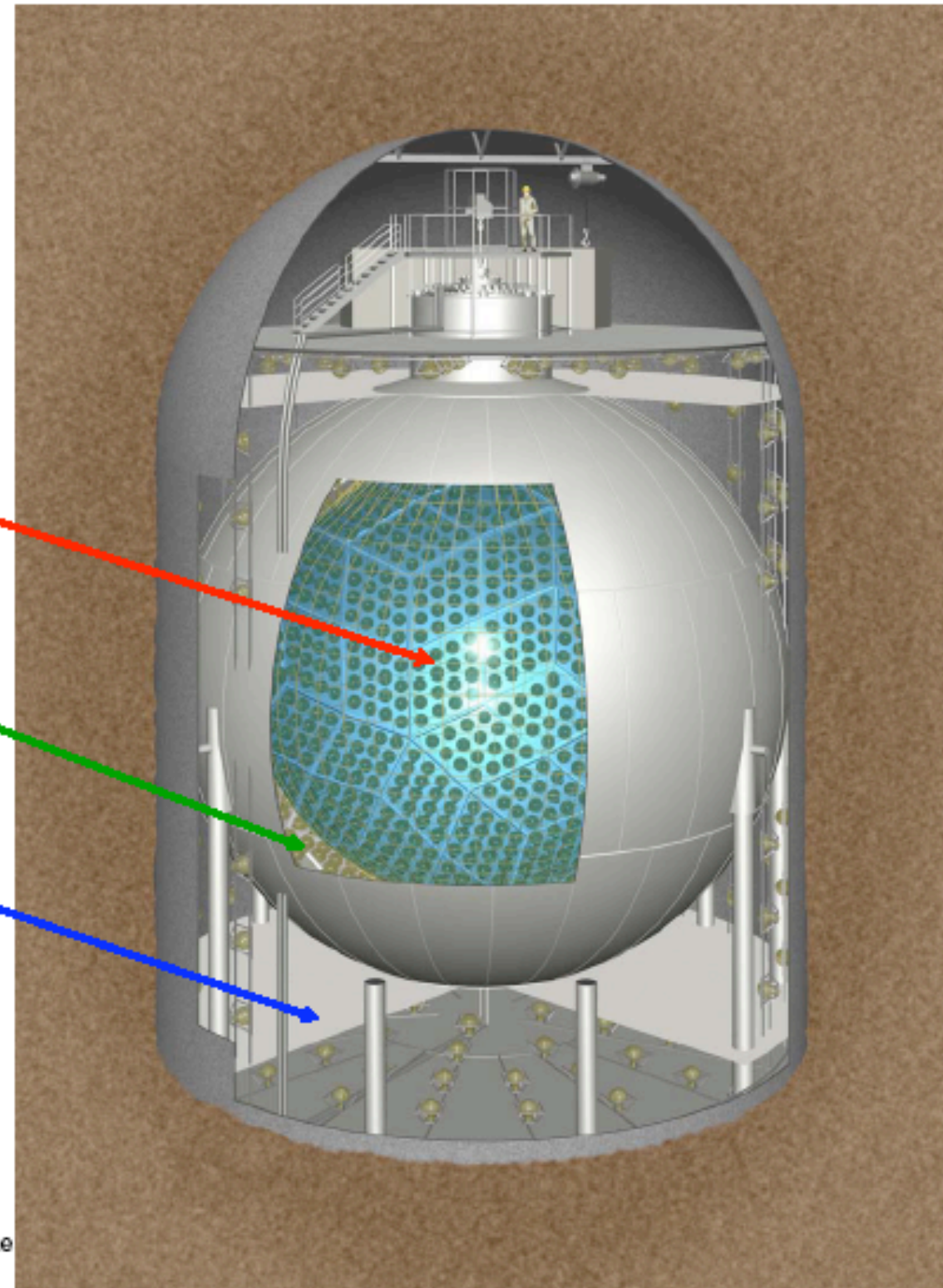




# KAMLAND

## KamLAND: Kamioka Liquid scintillator AntiNeutrino Detector

- 1 kton liq. Scint. Detector  
in the Kamiokande cavern
- 1325 17" fast PMTs
- 554 20" large area PMTs
- 34% photocathode coverage
- H<sub>2</sub>O Cerenkov veto counter





# KAMLAND LOCATION & FLUX



Many reactors contribute to the antineutrino flux at KamLAND

Site	Dist (km)	Cores (#)	$P_{\text{therm}}$ (GW)	Flux ( $\text{cm}^{-2} \text{s}^{-1}$ )	Rate noosc* ( $\text{yr}^{-1} \text{kt}^{-1}$ )
Kashiwazaki	160	7	24.3	$4.1 \cdot 10^5$	254.0
Ohi	179	4	13.7	$1.9 \cdot 10^5$	114.3
Takahama	191	4	10.2	$1.2 \cdot 10^5$	74.3
Tsuruga	138	2	4.5	$1.0 \cdot 10^5$	62.5
Hamaoka	214	4	10.6	$1.0 \cdot 10^5$	62.0
Mihama	146	3	4.9	$1.0 \cdot 10^5$	62.0
Sika	88	1	1.6	$9.0 \cdot 10^4$	55.2
Fukushima1	349	6	14.2	$5.1 \cdot 10^4$	31.1
Fukushima2	345	4	13.2	$4.8 \cdot 10^4$	29.5
Tokai2	295	1	3.3	$1.6 \cdot 10^4$	10.1
Onagawa	431	3	6.5	$1.5 \cdot 10^4$	9.3
Simane	401	2	3.8	$1.0 \cdot 10^4$	6.3
Ikata	561	3	6.0	$8.3 \cdot 10^3$	5.1
Genkai	755	4	10.1	$7.8 \cdot 10^3$	4.8
Sendai	830	2	5.3	$3.4 \cdot 10^3$	2.1
Tomari	783	2	3.3	$2.3 \cdot 10^3$	1.4
Ulchin	712	4	11.5	$9.9 \cdot 10^3$	6.1
Yonggwang	986	6	17.4	$7.8 \cdot 10^3$	4.8
Kori	735	4	9.2	$7.5 \cdot 10^3$	4.6
Wolsong	709	4	8.2	$7.1 \cdot 10^3$	4.3
<b>Total Nominal</b>	-	<b>70</b>	<b>181.7</b>	<b><math>1.3 \cdot 10^6</math></b>	<b>803.8</b>

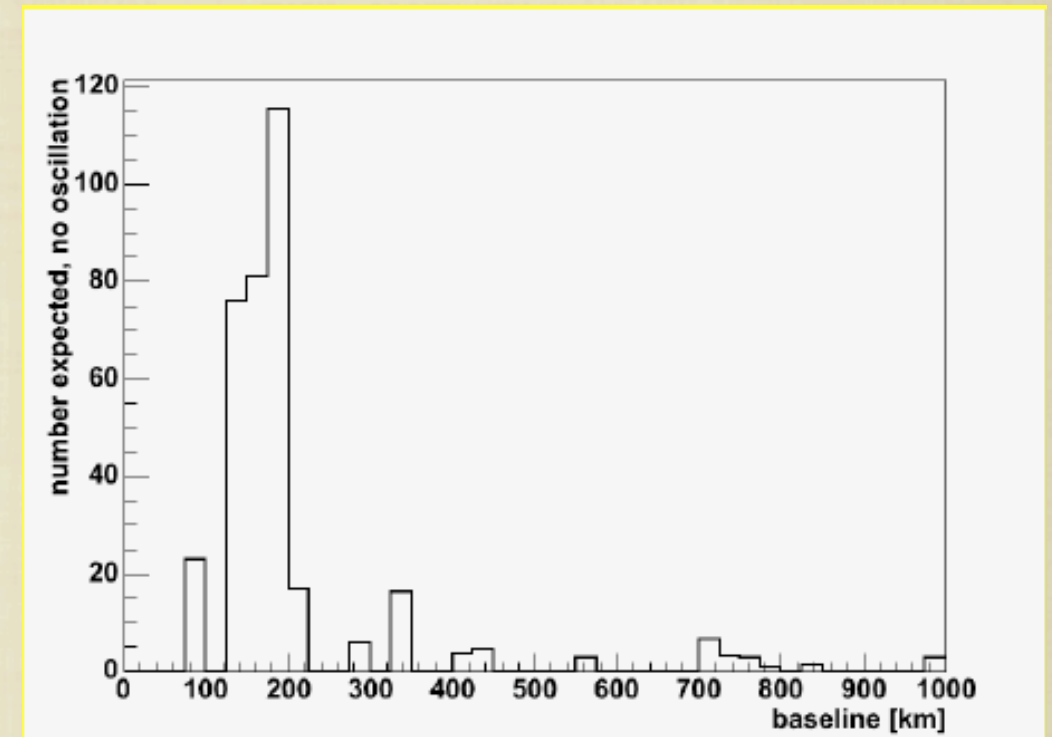
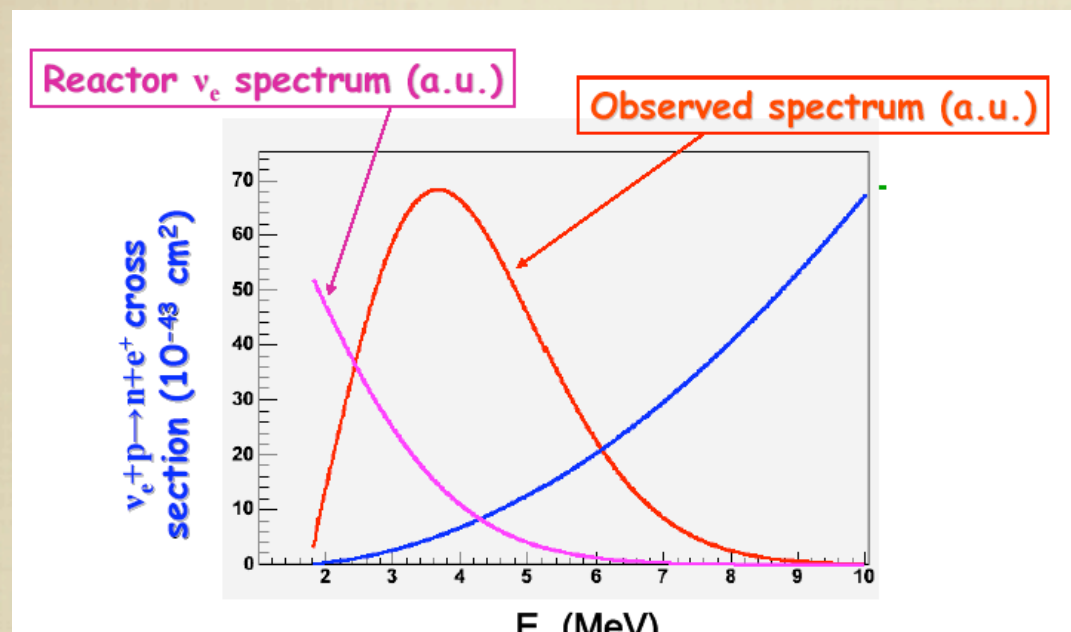
\* $E_{\nu} > 3.4 \text{ MeV}$   
( $E_{\text{prompt}} > 2.6 \text{ MeV}$ )

Detailed power and fuel  
Composition calculation used

From electrical power  
Japanese average  
fuel used



# SPECTRUM AND BASELINES



$$L_{osc} (Km) \approx \frac{E (GeV)}{1.27 \Delta m^2 (eV^2)}$$

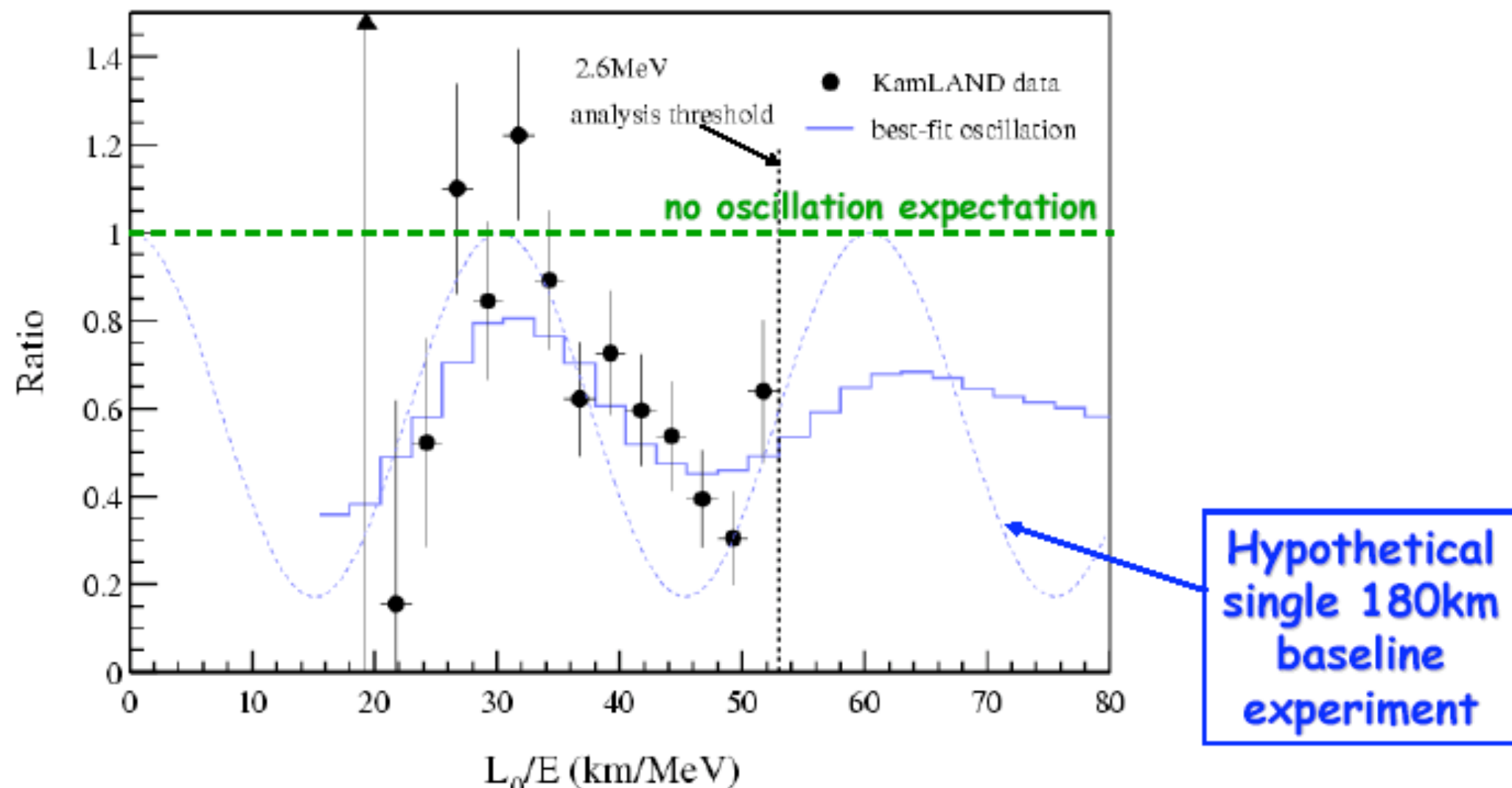
$$E \approx 3 \text{ MeV}, \quad \Delta m^2 \approx 10^{-5} \text{ eV}^2$$

$$L_{osc} \approx 200 \text{ Km}$$

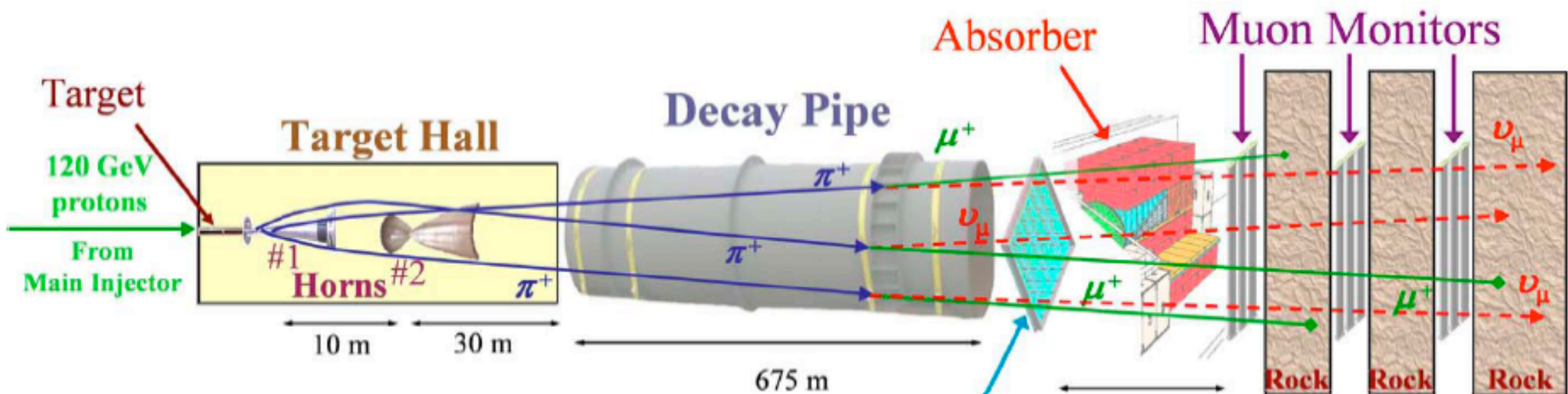
# L/E EFFECT IN KAMLAND

KamLAND uses a range of L and  
it cannot assign a specific L to each event

Nevertheless the ratio of detected/expected  
for  $L_0/E$  (or  $1/E$ ) is an interesting quantity, as it decouples  
the **oscillation pattern** from the reactor energy spectrum

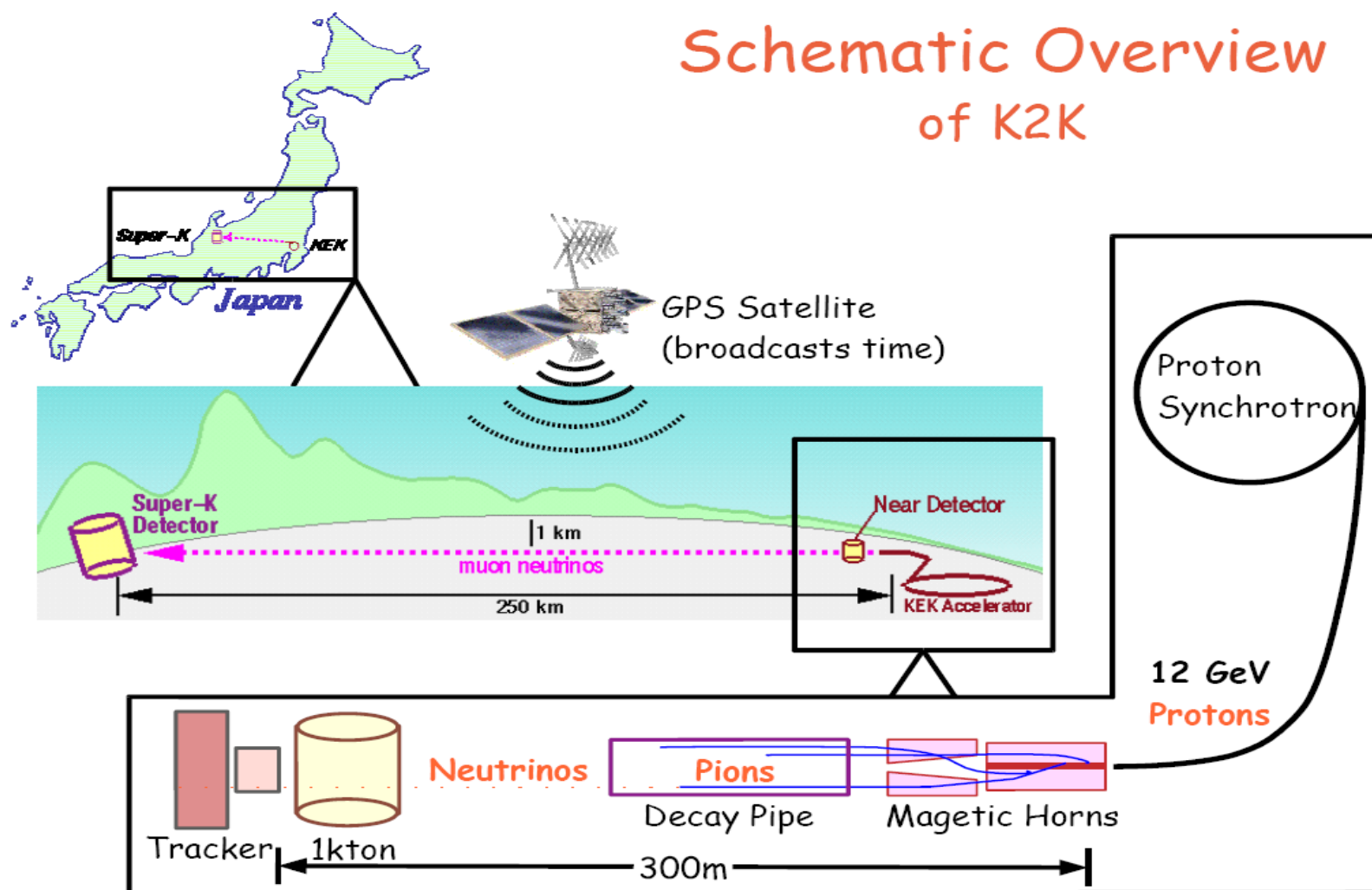






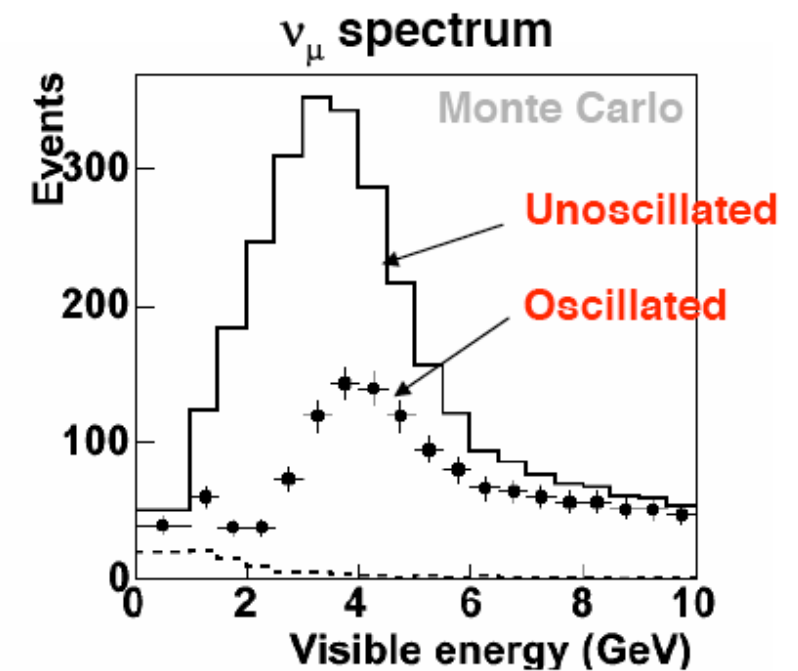
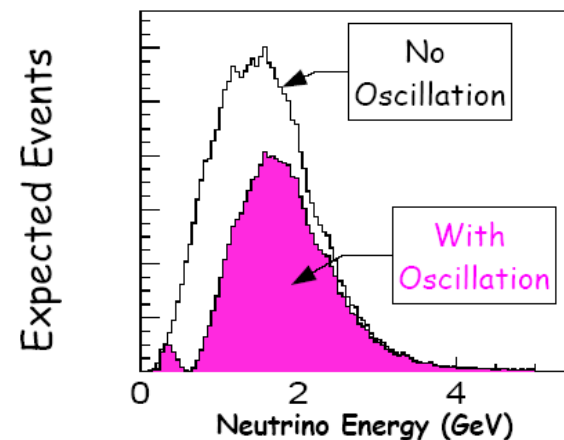
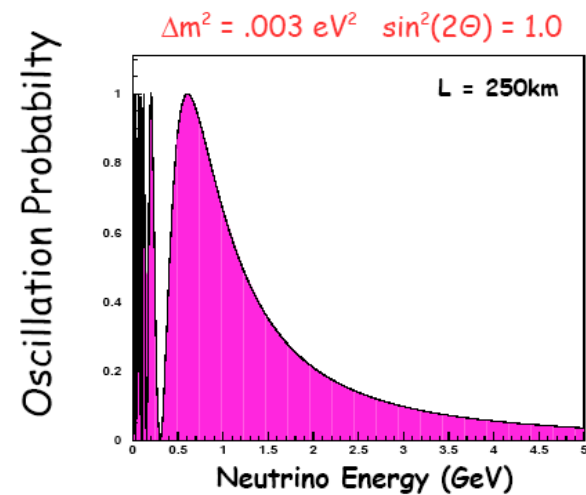
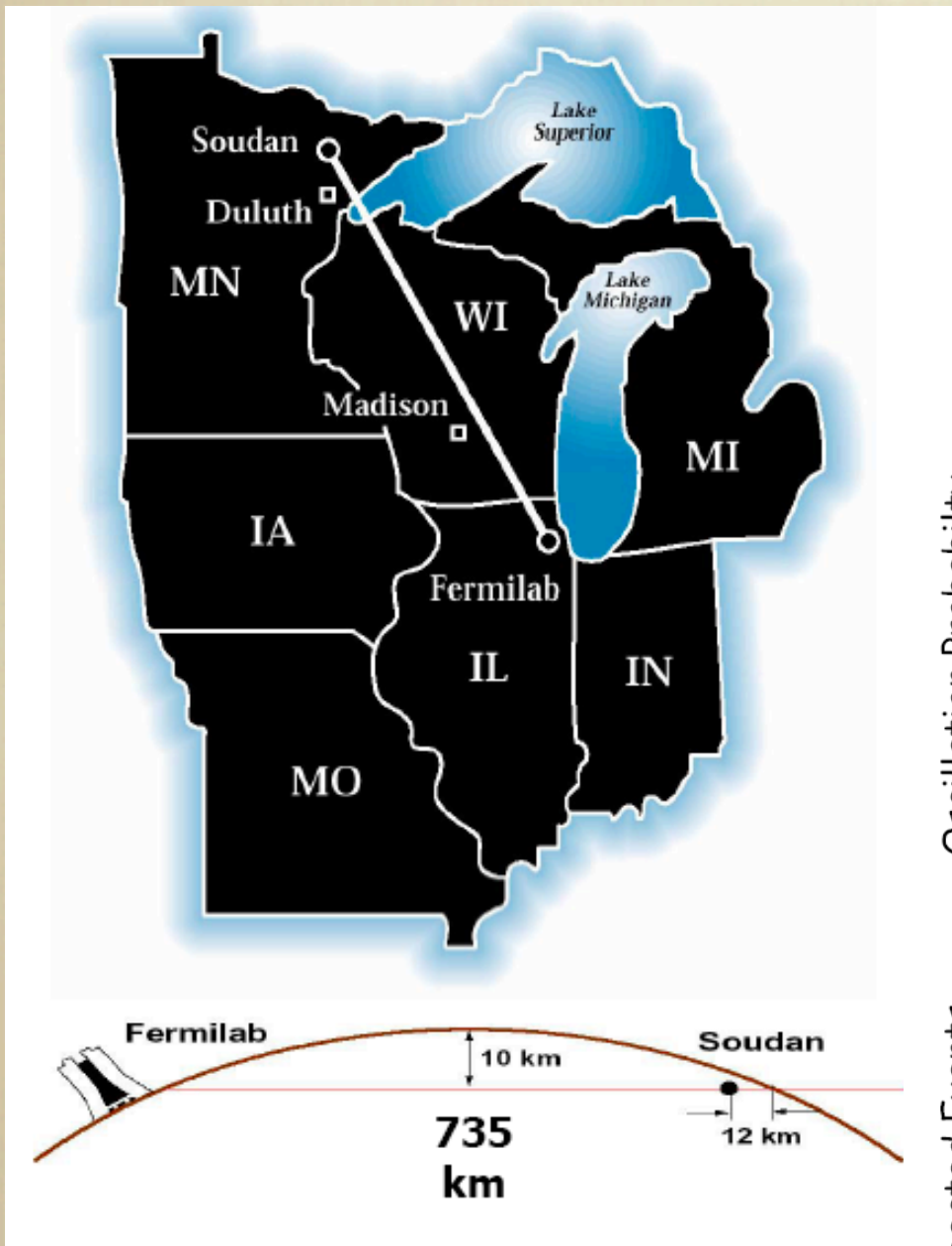
Schematic Overview of K2K

Minos & K2K:  
new skin for  
the old  
ceremony





# K2K/MINOS: CONFIRM ATMOSPHERIC OSCILLATION WITH A CONTROLLED BEAM



$E_{K2K} \sim 1\text{GeV} \Rightarrow L \sim 250 \text{ Km}$

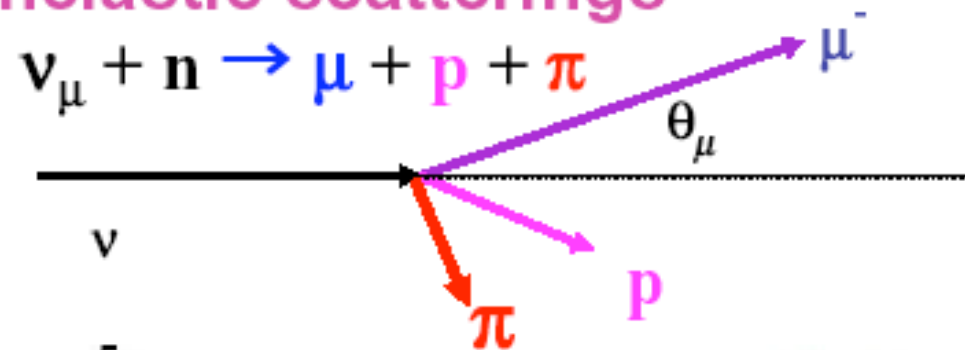
$E_{Numi} \sim 3\text{GeV} \Rightarrow L \sim 750 \text{ Km}$



# CROSS SECTIONS AND ENERGY RECONSTRUCTION

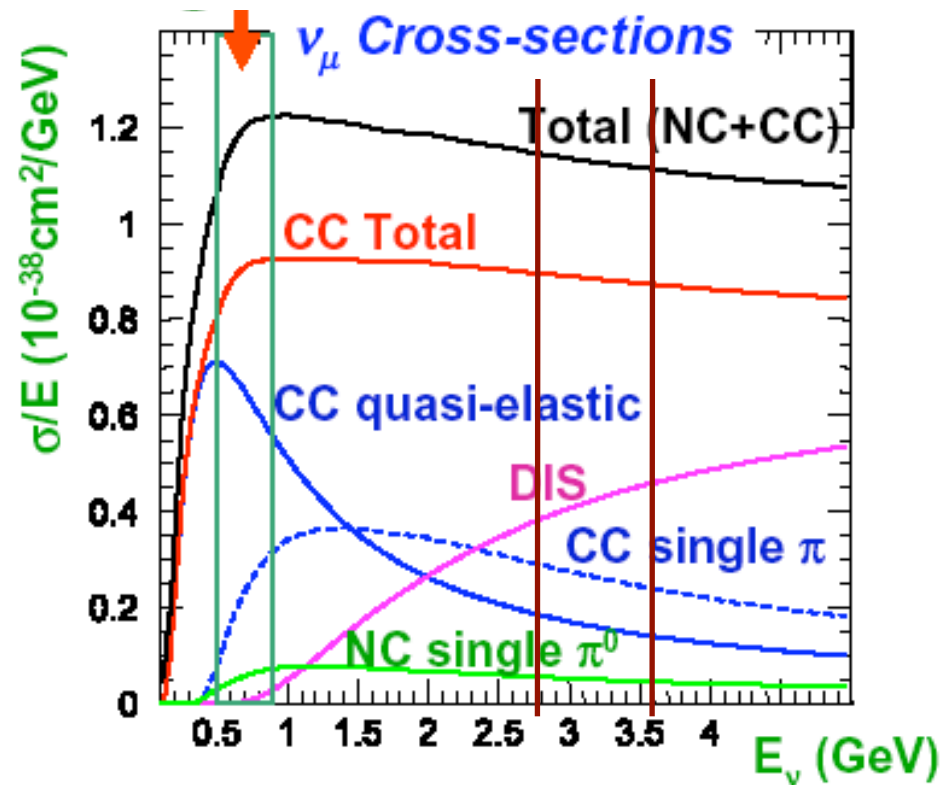
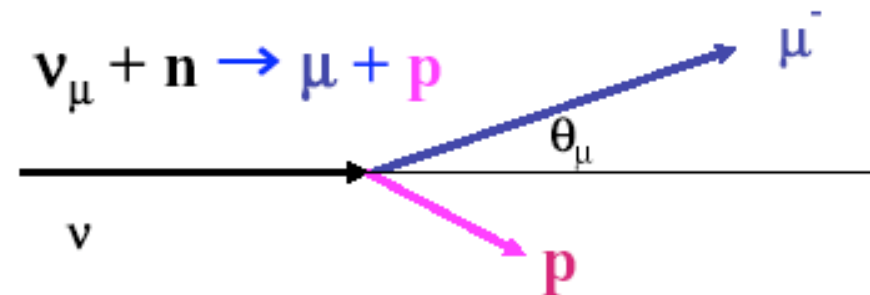
NUMI/  
MINOS

Inelastic scatterings

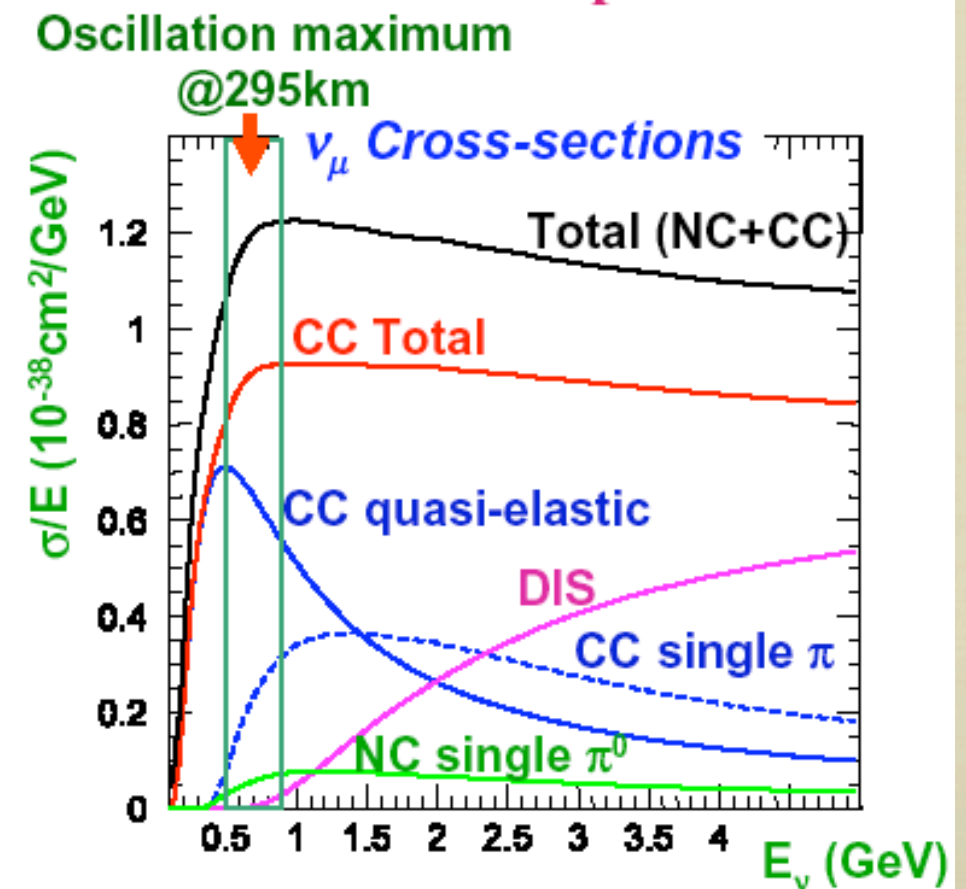


K2K

CC quasi elastic scatterings



OSCILLATION MAXIMUM@750 KM





# Detectors

Far Detector



Near Detector



DIFFEREN

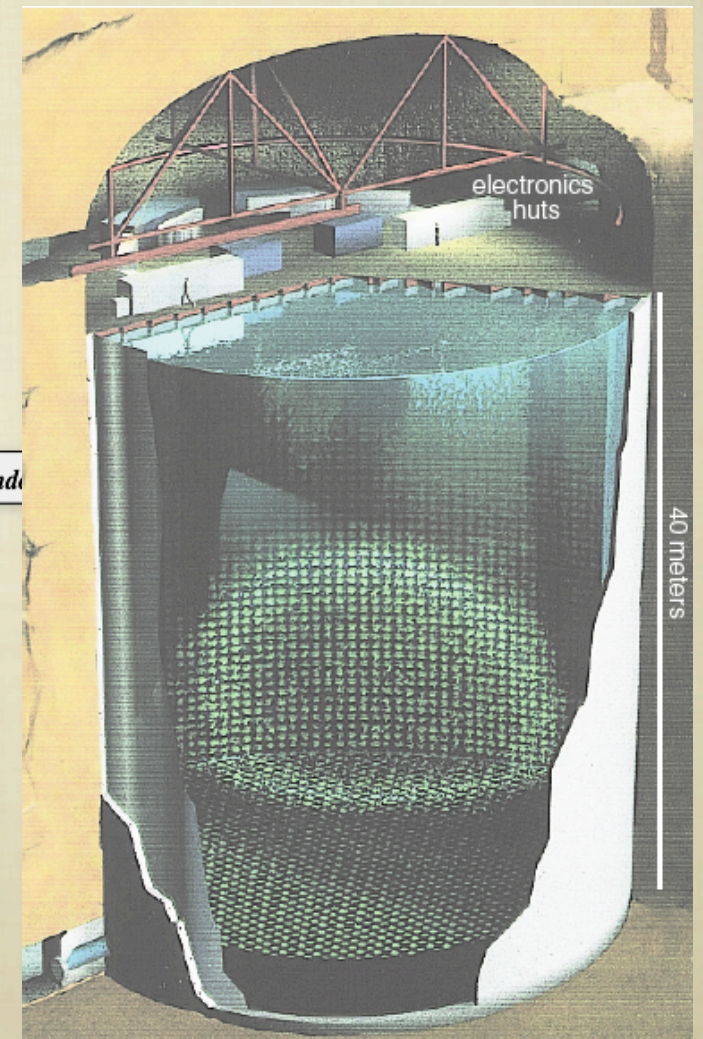
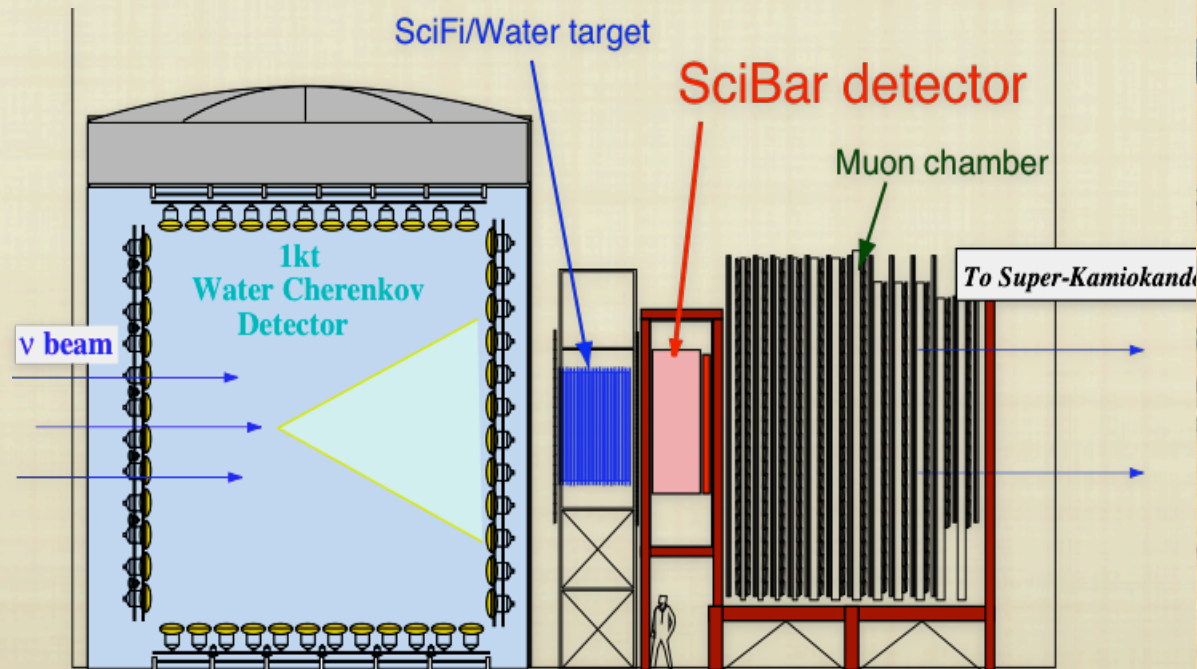
T

TECHNOL

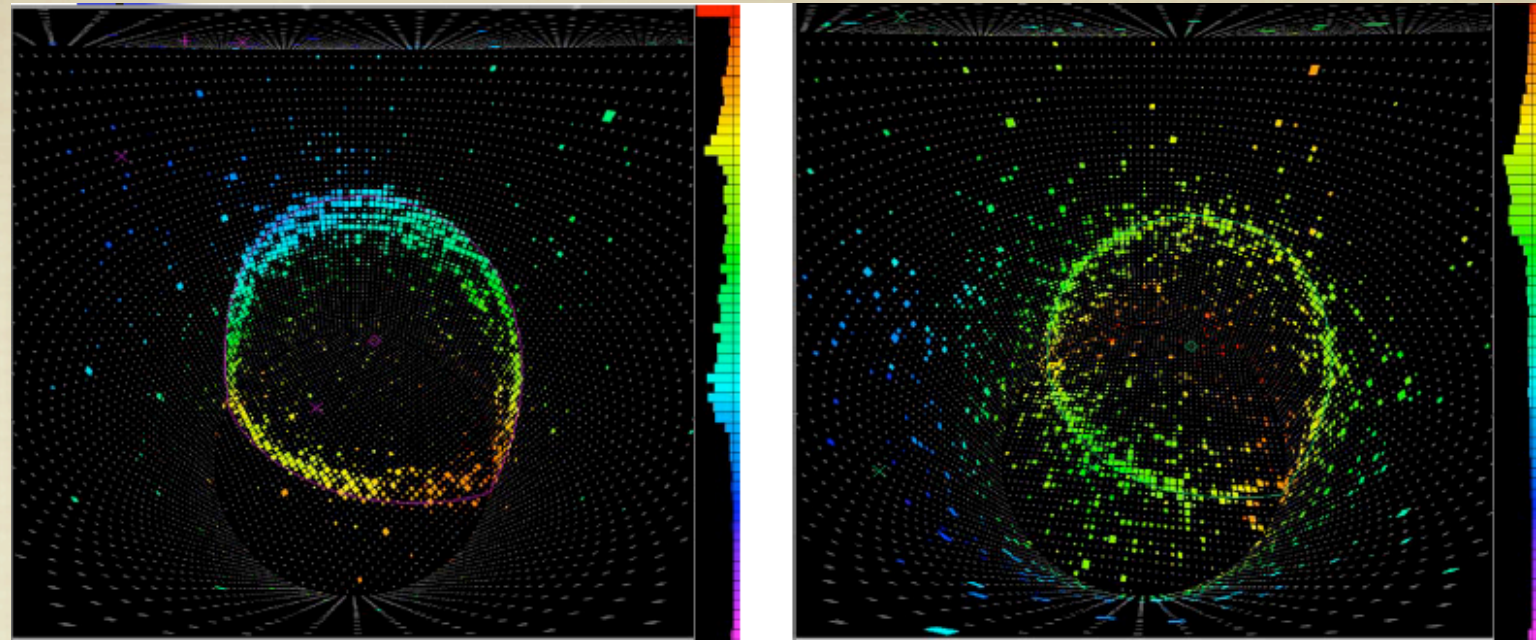
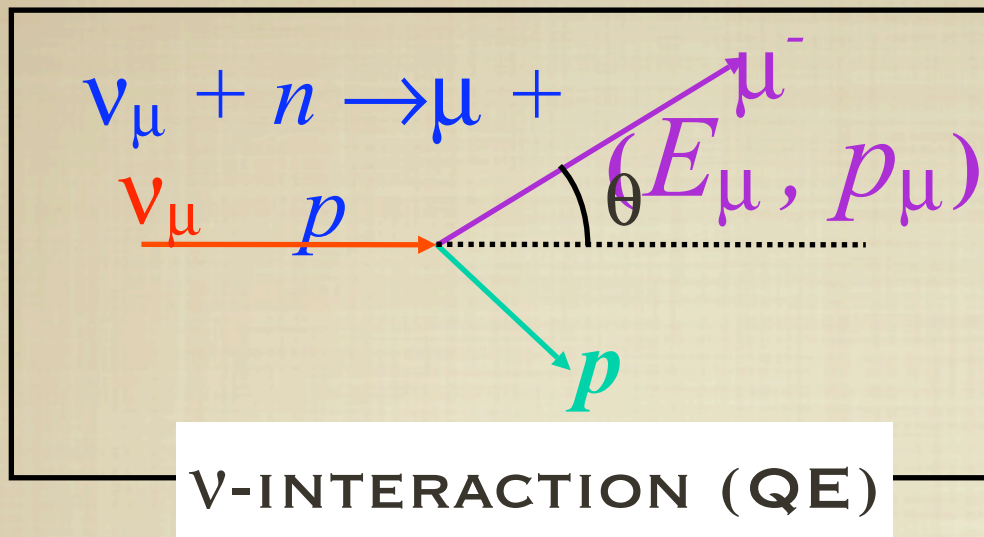
OGIES:

WATER &

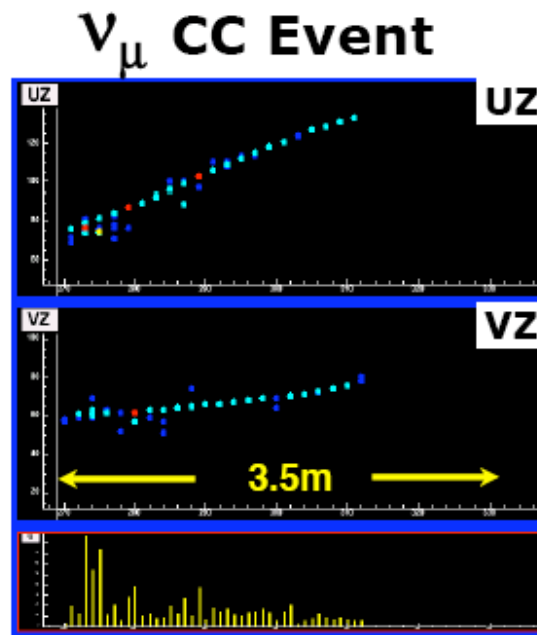
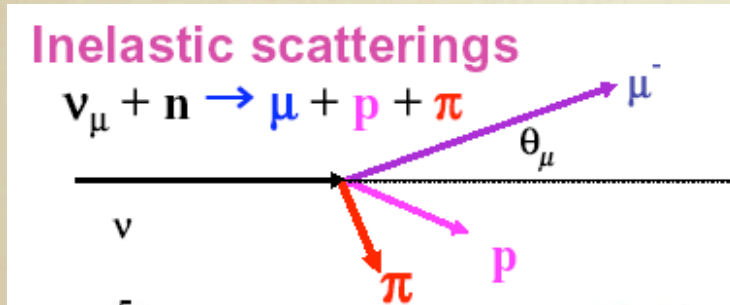
IRON



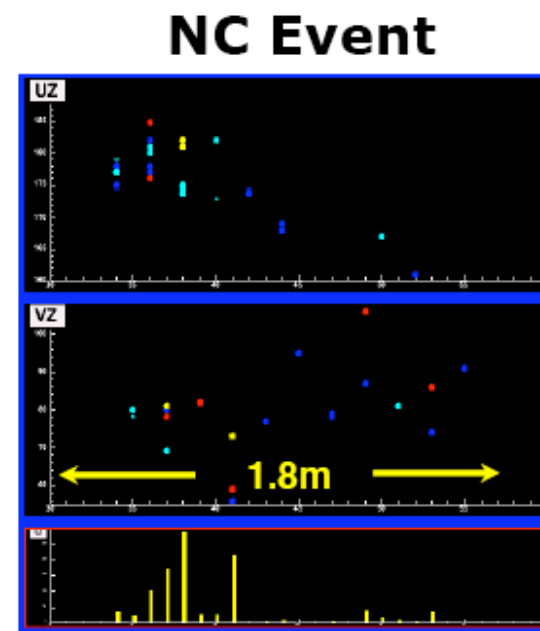




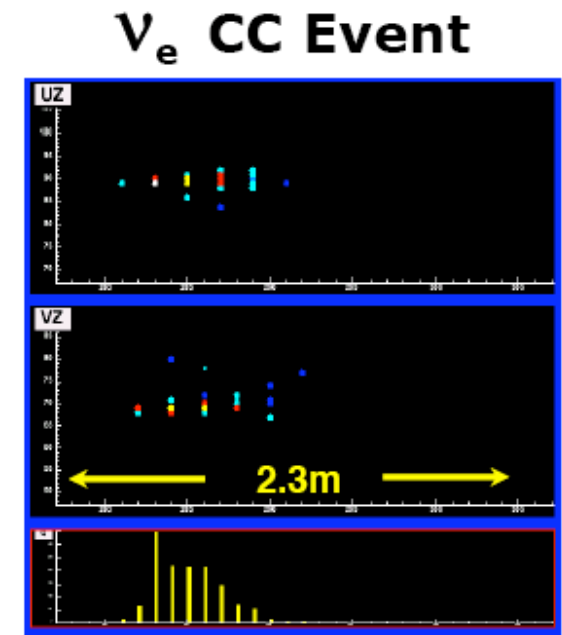
# Neutrinos in water & iron



• long  $\mu$  track+ hadronic activity at vertex



• short event, often diffuse



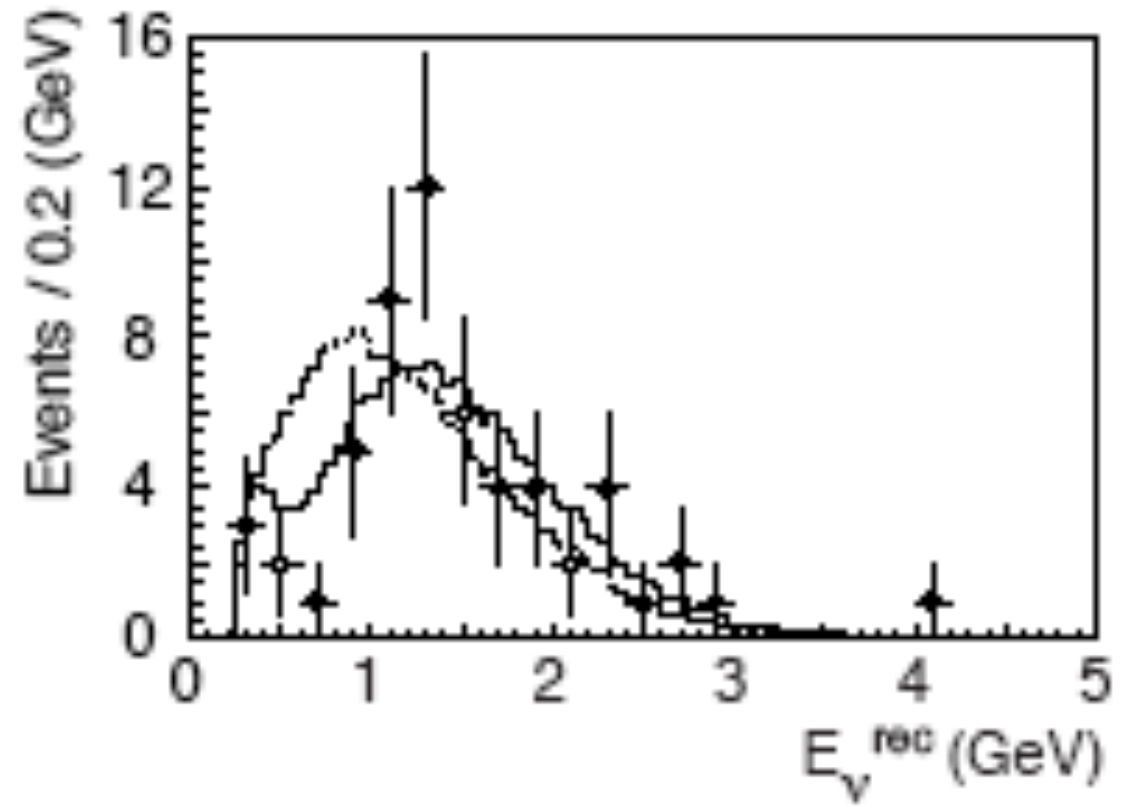
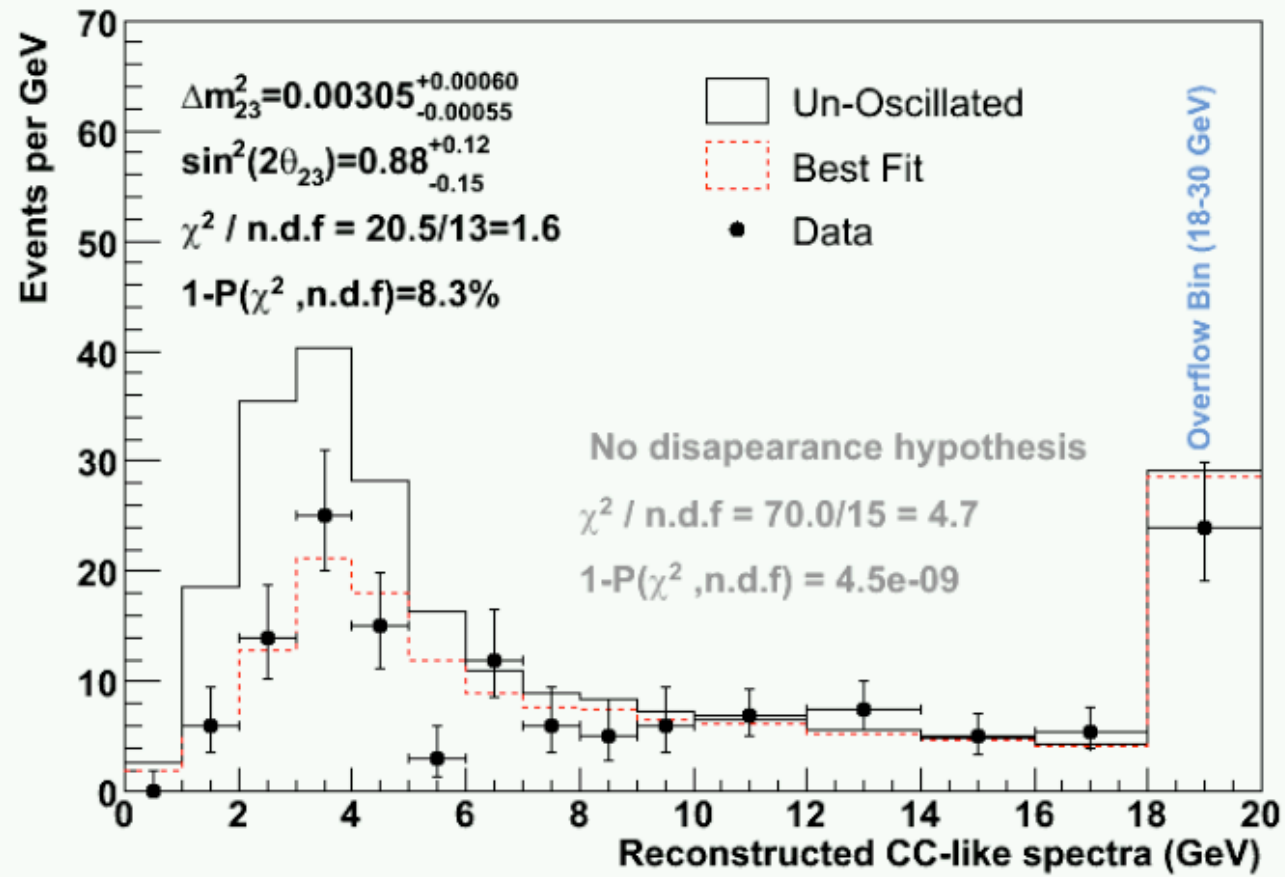
• short, with typical EM shower profile

$$E_\nu = E_{\text{shower}} + P_\mu$$

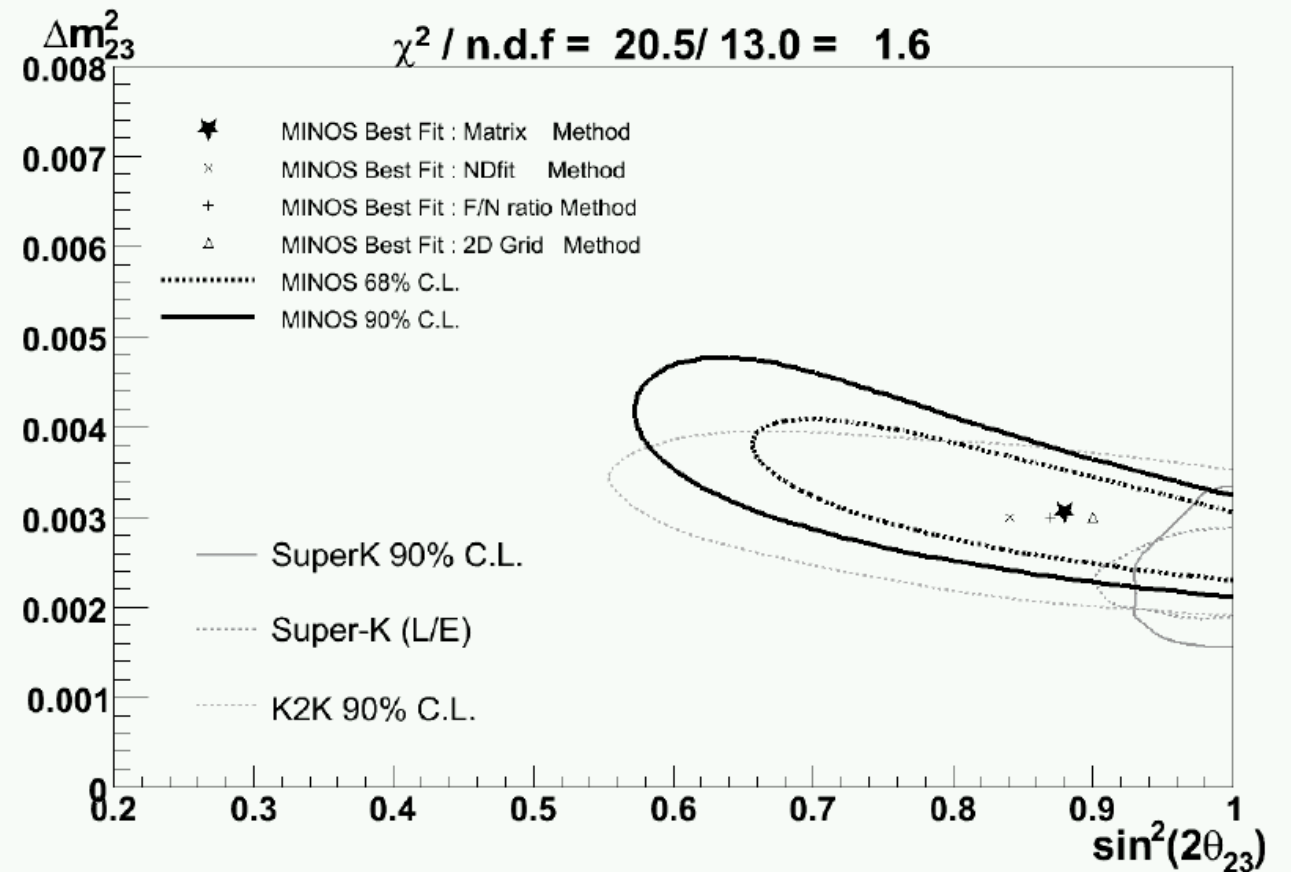
55%/√E

6% range, 10% curvature

# Oscillation Results for 0.93E20 p.o.t

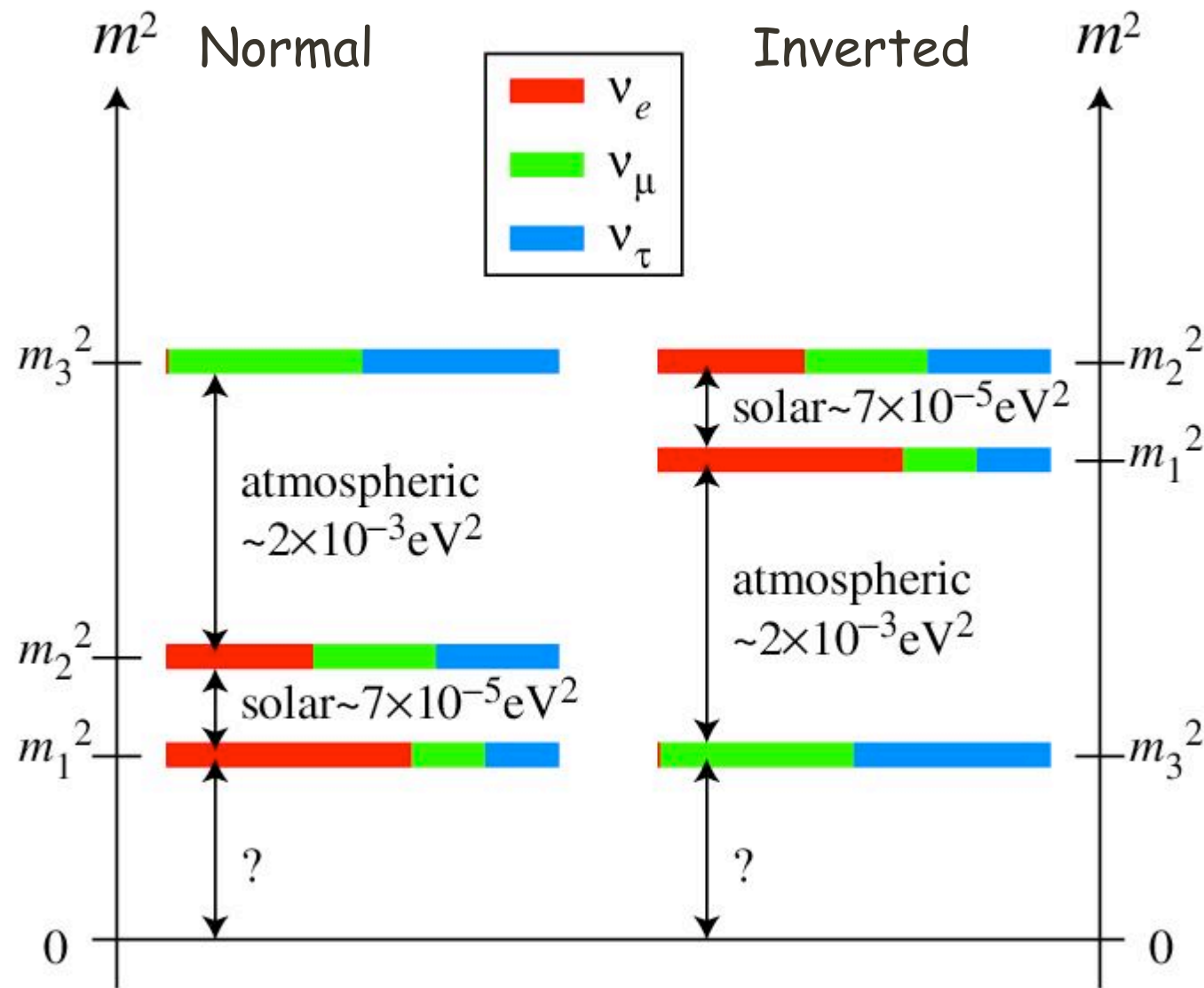


ATMOSPHERIC  
OSCILLATION  
CONFIRMED!

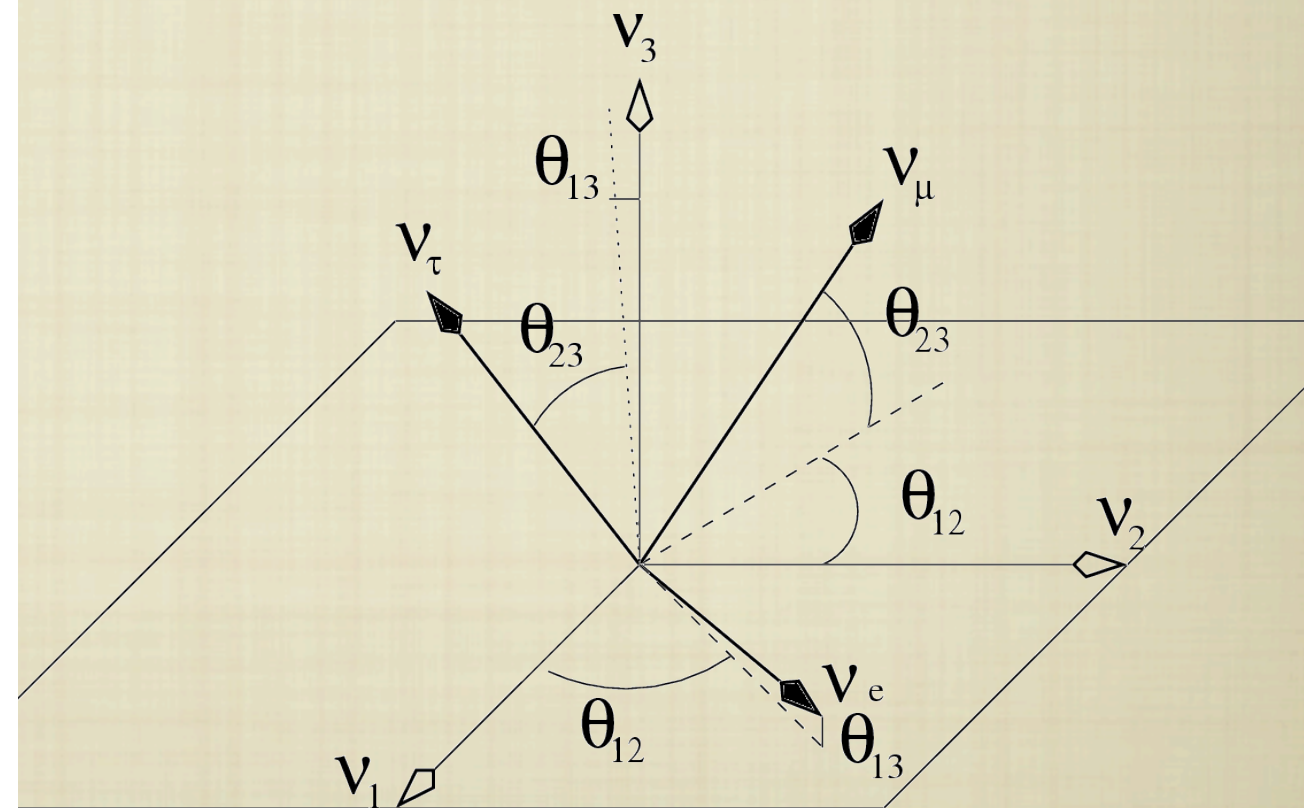




# Oscillations revisited



Parameter	Best-fit value	$3\sigma$ range
$\theta_{12}$	$33.2^\circ$	$28.7^\circ \text{ .. } 38.1^\circ$
$\theta_{23}$	$45.0^\circ$	$35.7^\circ \text{ .. } 55.6^\circ$
$\theta_{13}$	$0.0^\circ$	$0^\circ \text{ .. } 12.5^\circ$



# THE PNMS MATRIX

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Unless the other two angles  $\theta_{13}$  is small  
(experimental upper limit  $\theta_{13} < 10^\circ$ )

If  $\delta \neq 0, \pi, 2\pi \dots$  then weak interactions violate CP symmetry in the lepton sector (as in the quark sector)

atmospheric

CP violation phase

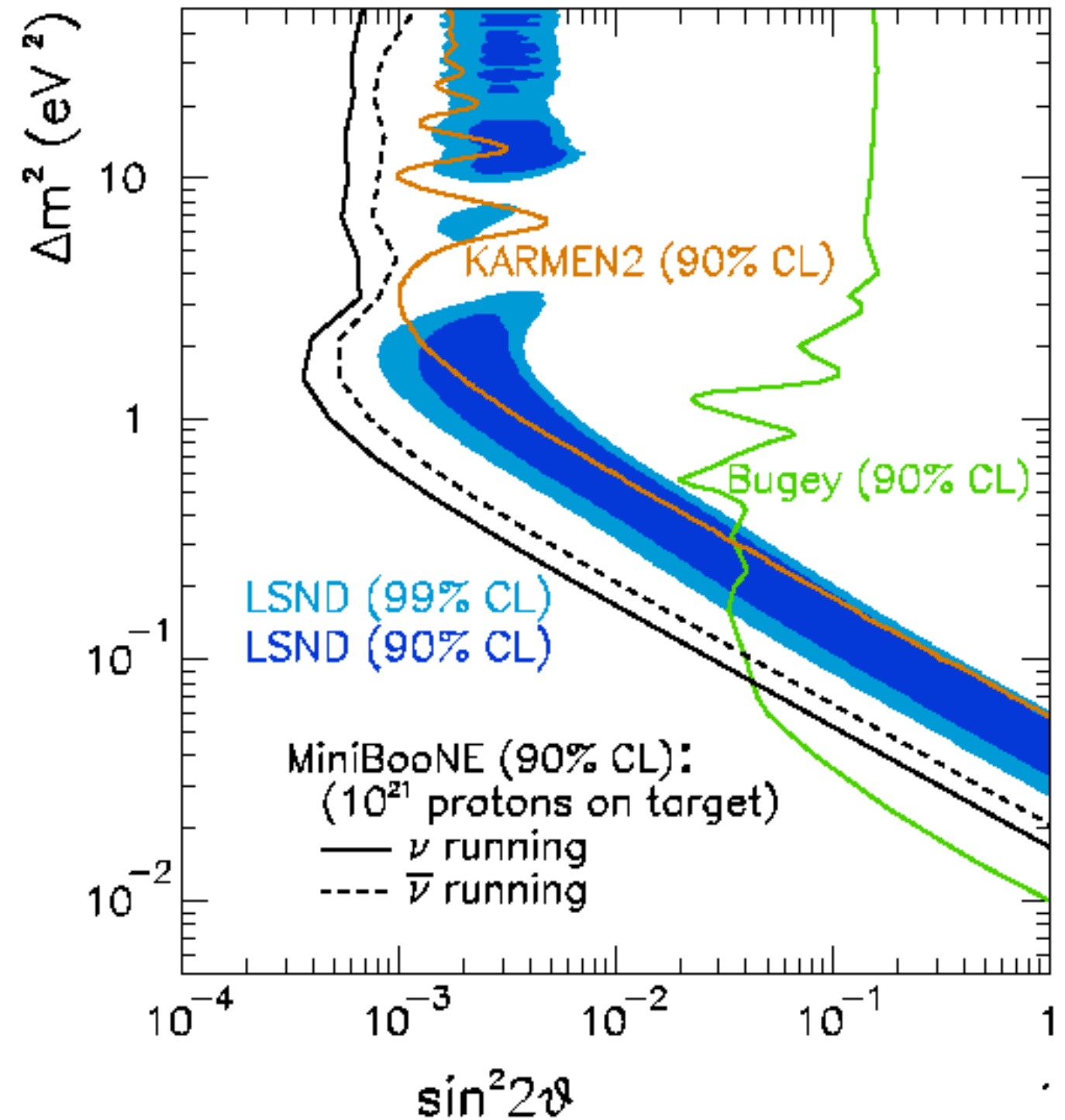
solar

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Links atmospheric & solar sectors

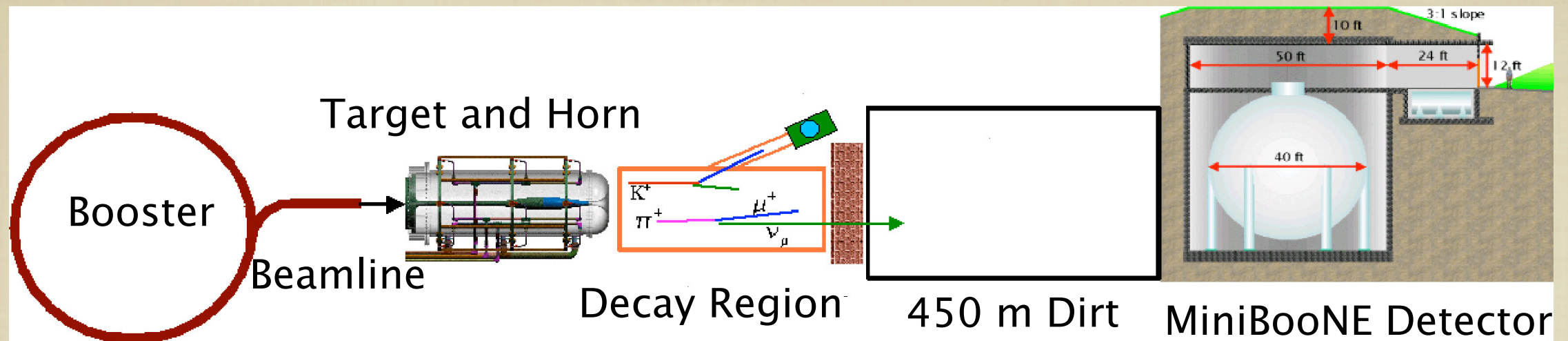


# THE LAST ANOMALY



$$\bar{\nu}_\mu \rightarrow \bar{\nu}_e ?$$

# MINIBOOONE



**Primary Beam**  
(protons)

**Secondary Beam**  
(mesons)

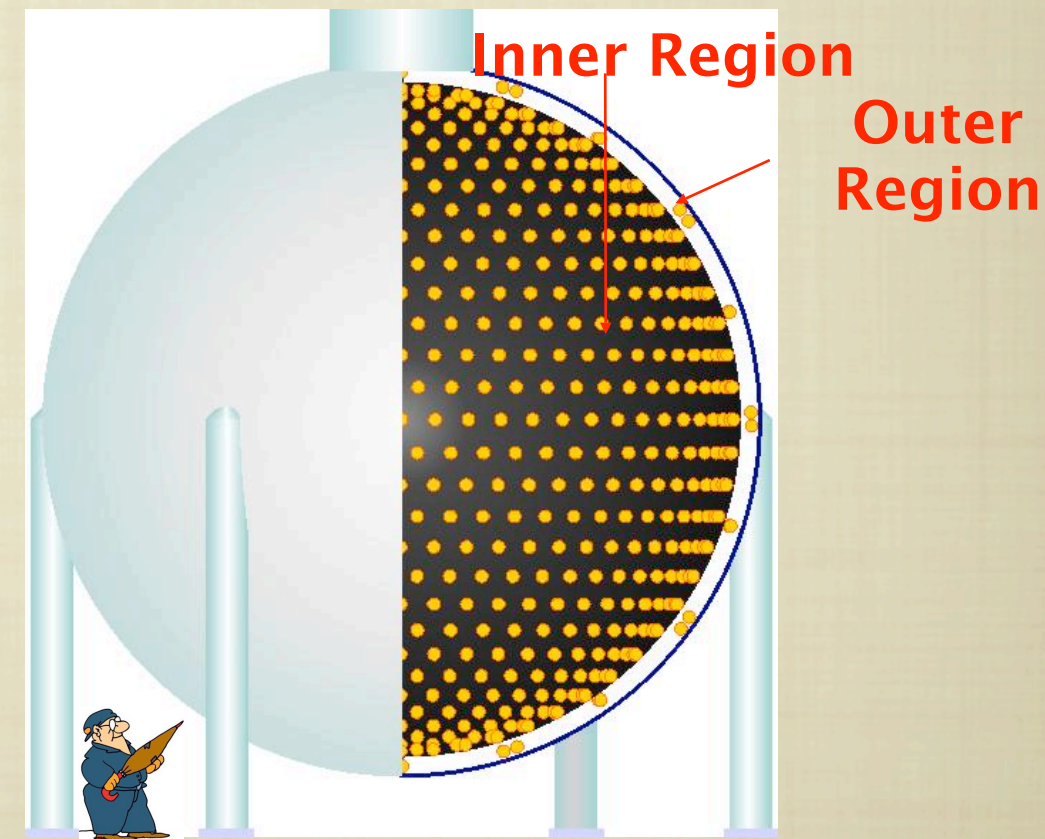
**Neutrino Beam**

## Beam

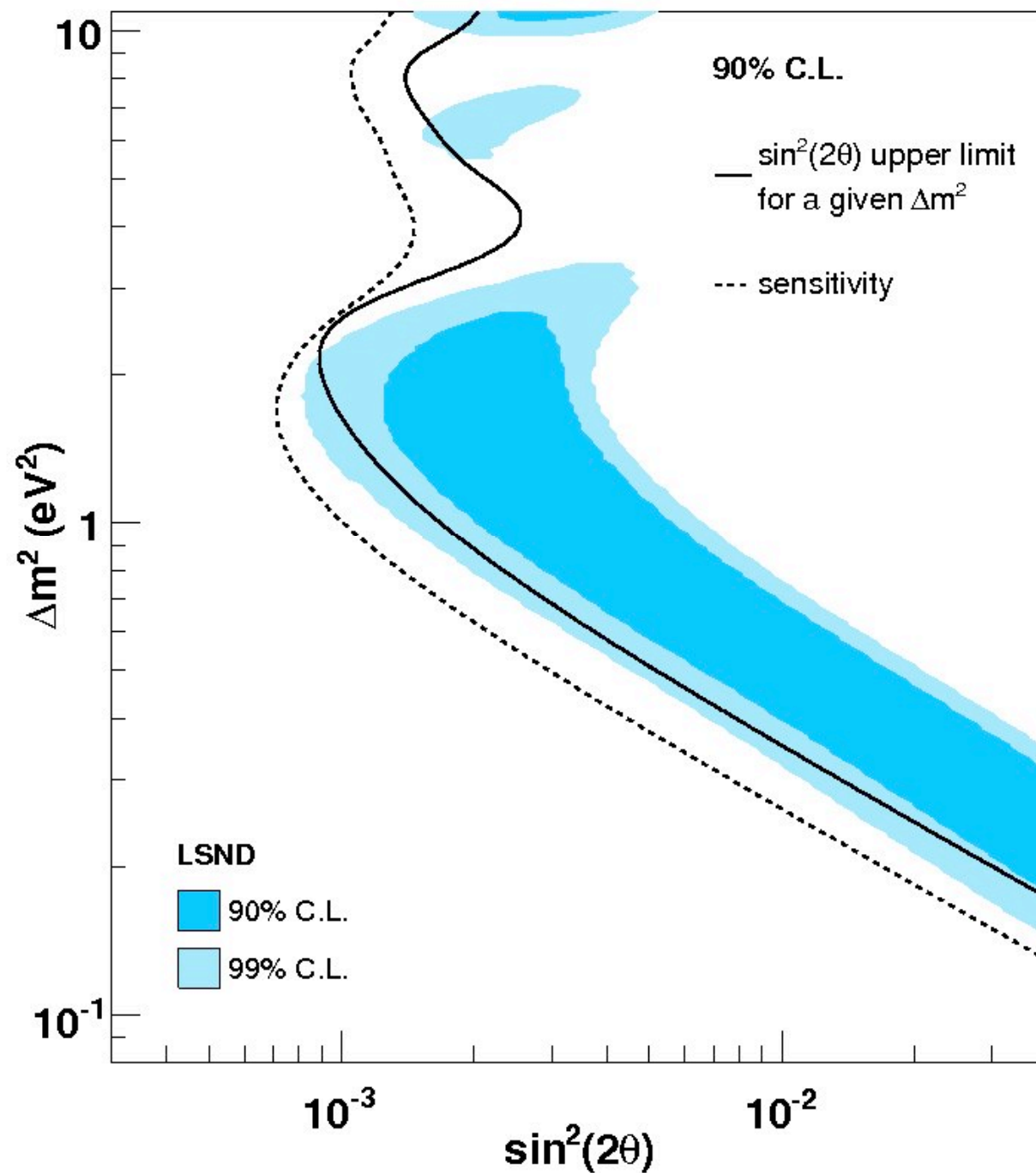
- ~0.5–1 GeV neutrinos or antineutrinos produced from FNAL Booster

## Detector

- 12 m in diameter sphere filled with 800 t of undoped mineral oil
- Neutrino interactions in oil seen via Cherenkov and scintillation light







- No overlap in 90% CL allowed LSND and MiniBooNE regions

- MiniBooNE **excludes** two neutrino appearance-only oscillations as the explanation of the LSND anomaly at **~98% CL**

- Any interpretation of the LSND anomaly that would produce a significant excess for  $E_\nu > 475$  MeV at MiniBooNE is also ruled out