

Detector Basics III

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Summary of lecture II

Cherenkov counters

unsegmented

requires low rates, relatively low multiplicity events

cannot be magnetized

best at ~ 1 GeV and below

solar, super-nova, atmospheric, accelerator neutrinos

e-like/mu-like particle ID

muons, pions, protons potentially below threshold

electron neutrino efficiency 40% (@ 0.7 GeV)

NC rejection 90%

channel counts 0.5 channels/kt [SK]

<cost/kt> \sim \$8M [SK,2009]

Tracking calorimeters

segmented

can operate at high rates and relatively high multiplicity events

can be magnetized

best at ~ 1 GeV and above

super-nova, atmospheric, accelerator, accelerator neutrinos

e/ γ / μ / π /p separation possible

all kinetic energy visible in principle

35% (@ 2 GeV)

99%

20 channels/kt [NOvA]

\sim \$15M [NOvA,2009]

Today: A bit of a mash up of some remaining topics

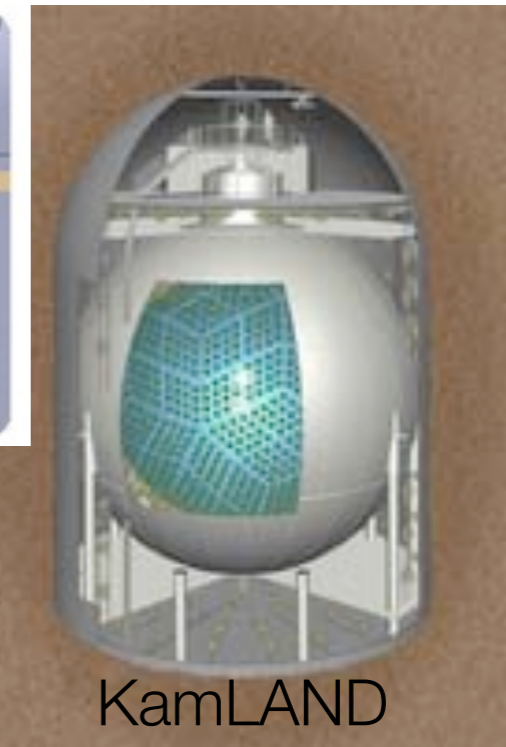
- Low energy scintillator detectors
- Tau neutrino detection
- Time projection chambers

Unsegmented liquid scintillator detectors

- Large volume of liquid scintillator viewed by PMT's
- Anti-electron neutrino detection from reactors at ~ 3.5 MeV
- Electron neutrino detection via elastic scattering from Sun at 0.7 MeV
- Scintillator allows for larger light collection (~ 200 photons/MeV) than water
- Used for detection of anti-neutrinos from reactor experiments (CHOOZ/KamLAND/Double CHOOZ/Daya Bay/Reno) and neutrinos from the Sun (Borexino)
- At these low energies the name of the game is background suppression from naturally occurring radioactive sources



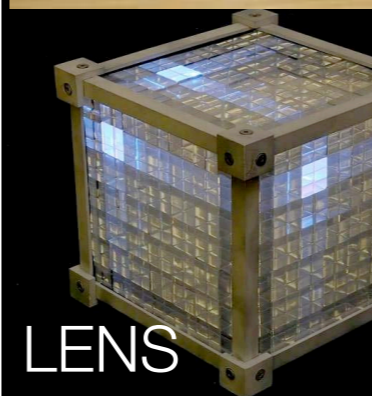
Double CHOOZ



KamLAND



Borexino



LENS



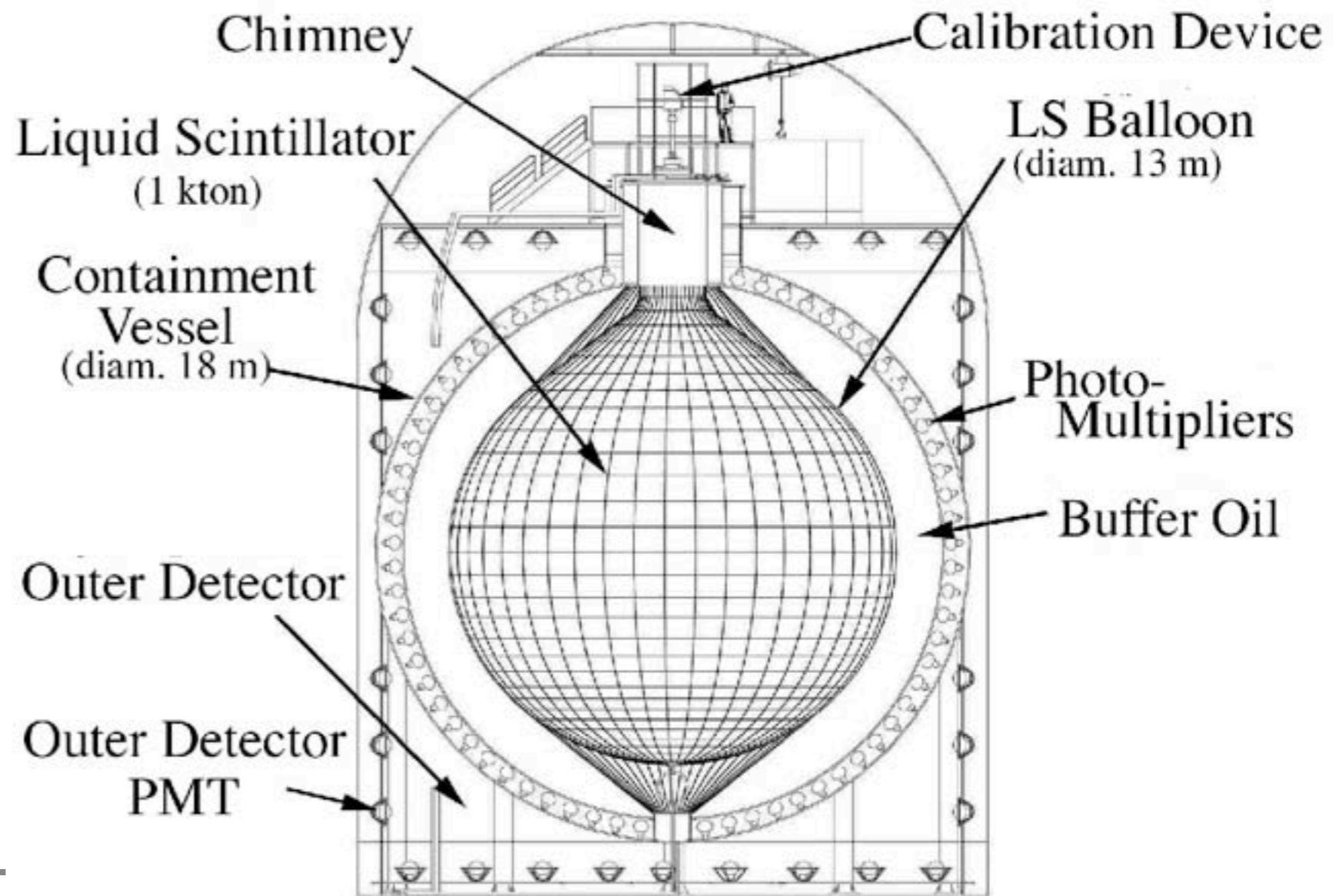
Daya Bay



Reno

Building for low background: Buffer zones

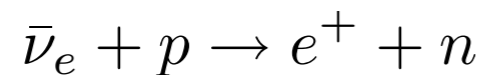
- For these low energy detectors it is common to build the detector in layers of buffer zones with careful control of components that go into the central most zones
- “Dirty” components, for example PMT glass which contains lots of U and Th, are kept away from the central regions.



KamLAND

Background rejection: Coincidence

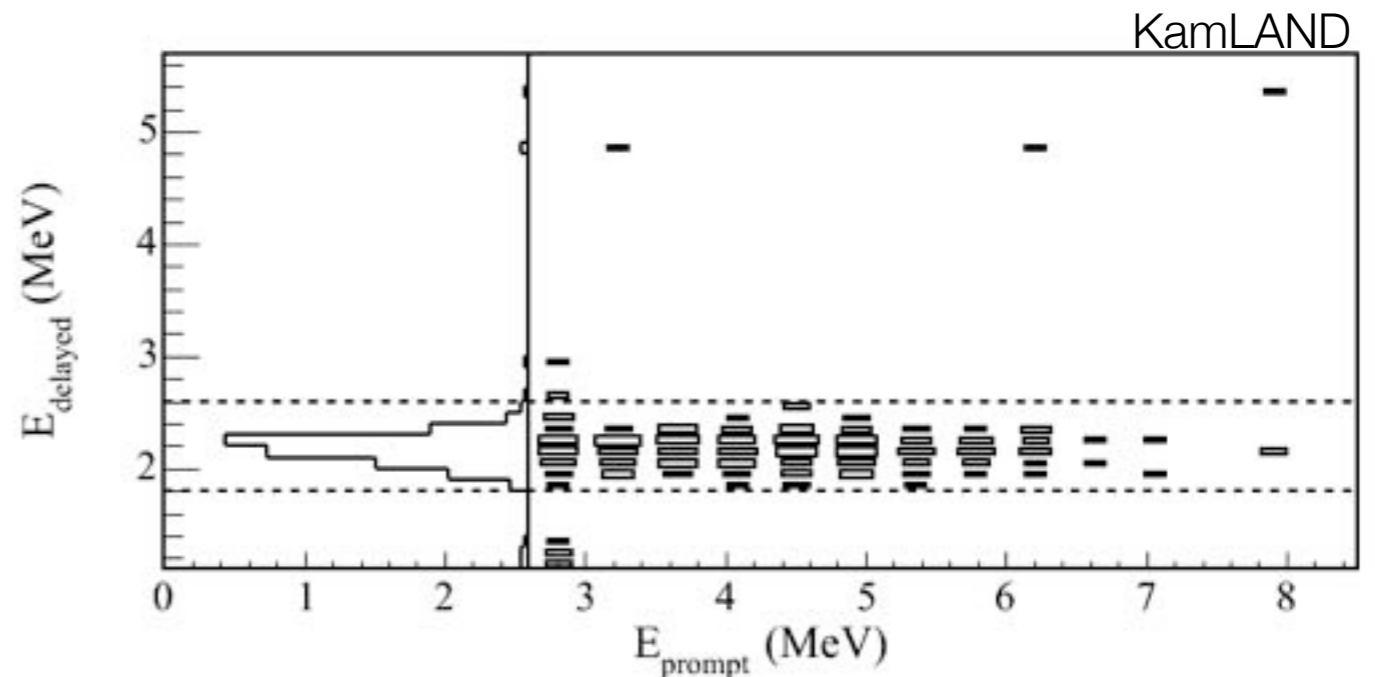
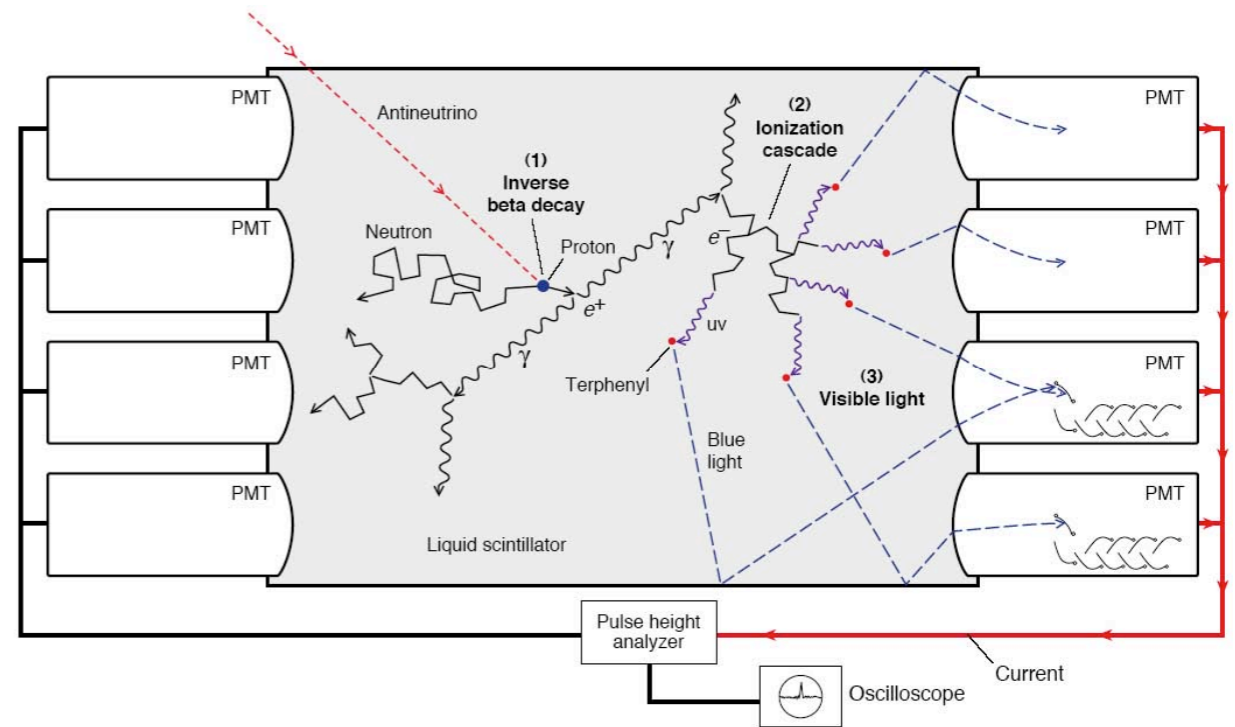
- For reactor neutrino experiment the detection channel is:



- Positron deposits its kinetic energy and annihilates promptly

$$E_{e^+} = E_{\bar{\nu}_e} - (M_n - M_p) + m_e$$

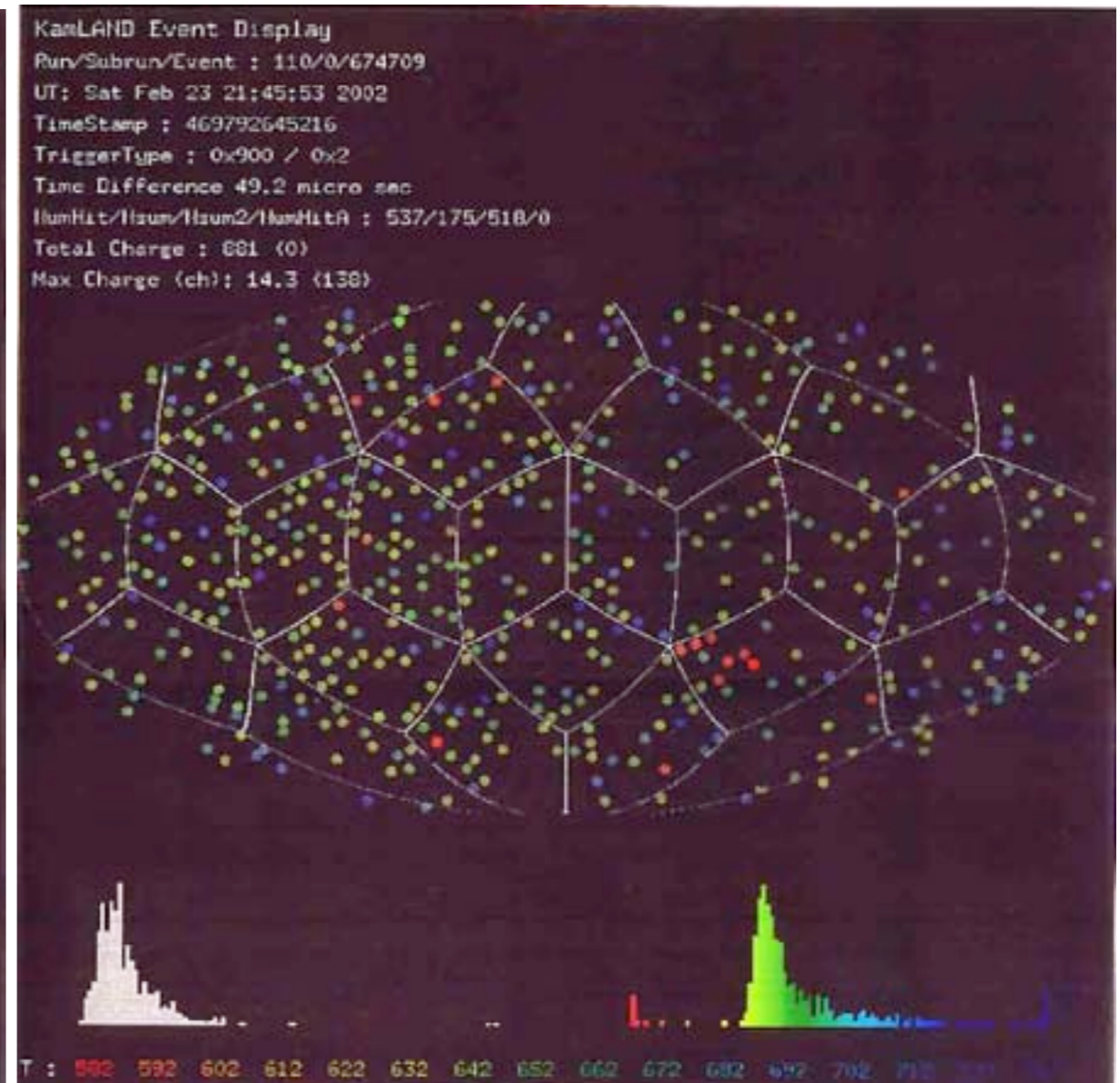
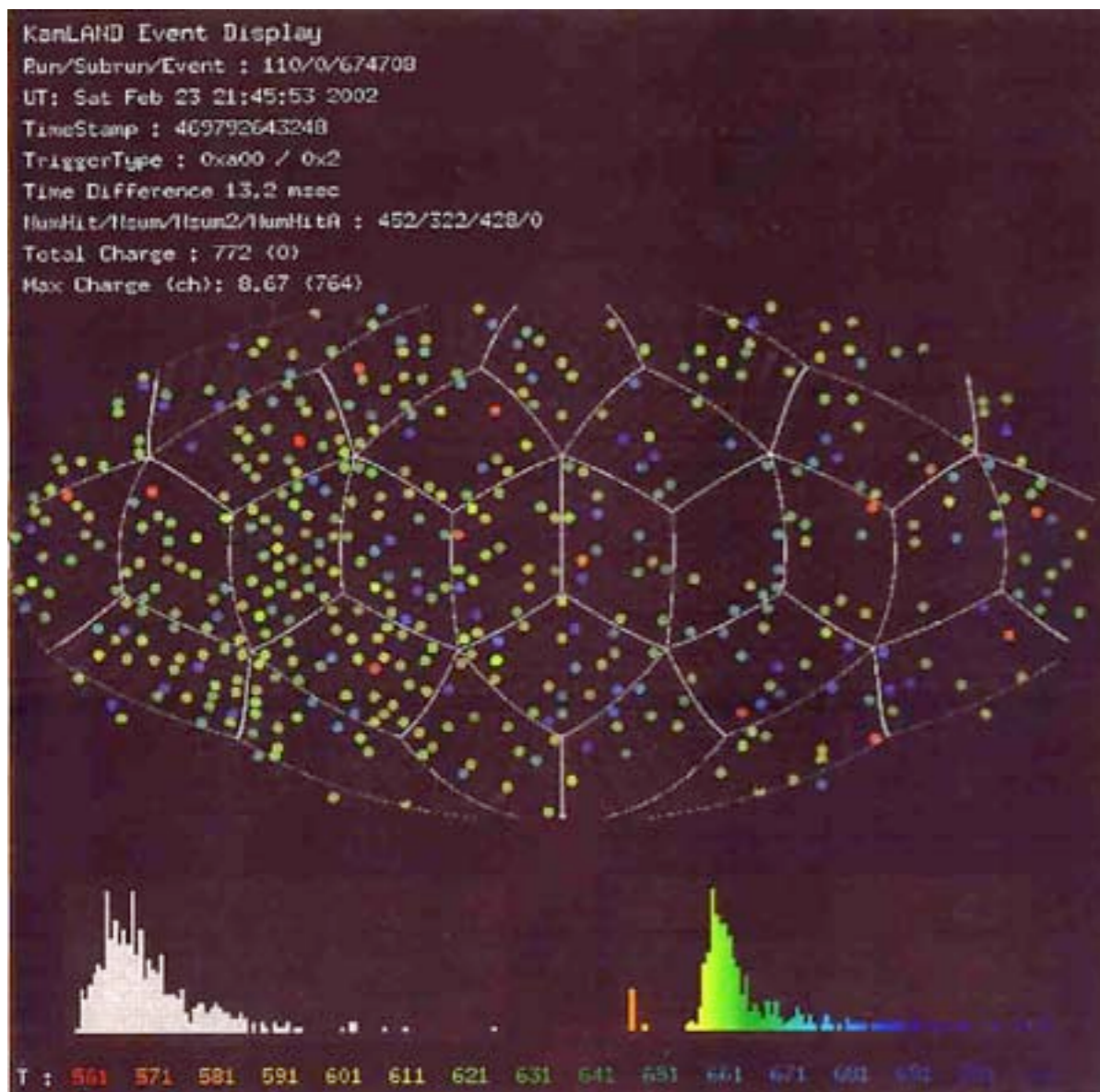
- Neutron wanders around for ~5 ms and is captured about 5 cm away from interaction vertex (on Gd or Cd dissolved in scintillator) releasing 2.2 MeV in gamma rays (in case of KamLAND shown here)
- This energy-time double coincidence signal dramatically beats the background down



KamLAND's signal

Feb 23 21:45:32 2002 469792643248

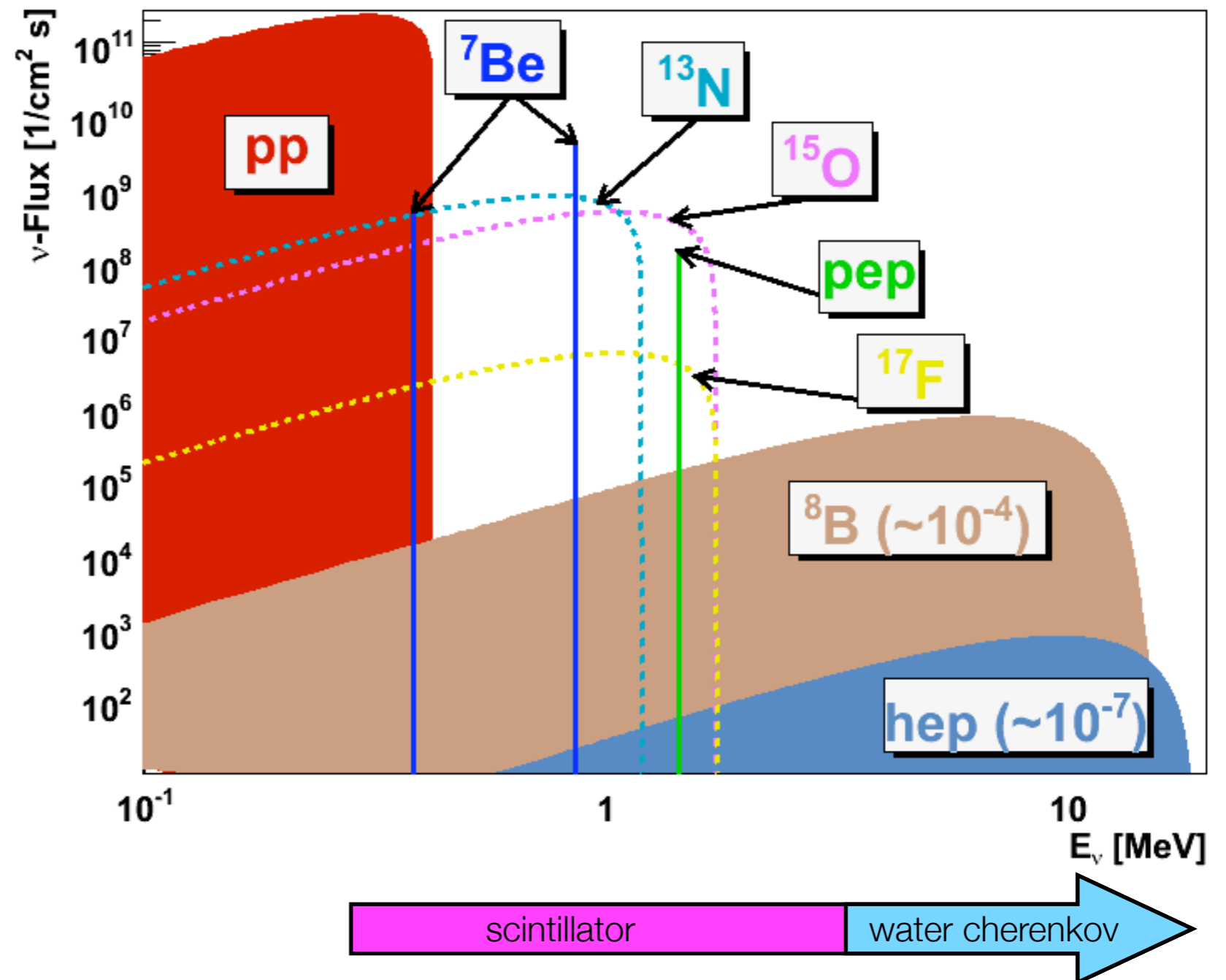
Feb 23 21:45:32 2002 469792645216



Times differ by ~ 2000 ticks of a 25 ns clock

Low energy solar neutrinos

- Water Cherenkov runs out of light around 4 MeV
- Switch to scintillator at lower energies
- Scintillation light is emitted isotropically. Gives up directional information - can't rely on pointing events to Sun
- Unlike inverse-beta decay, elastic scattering does not provide an energy-time coincidence signal.

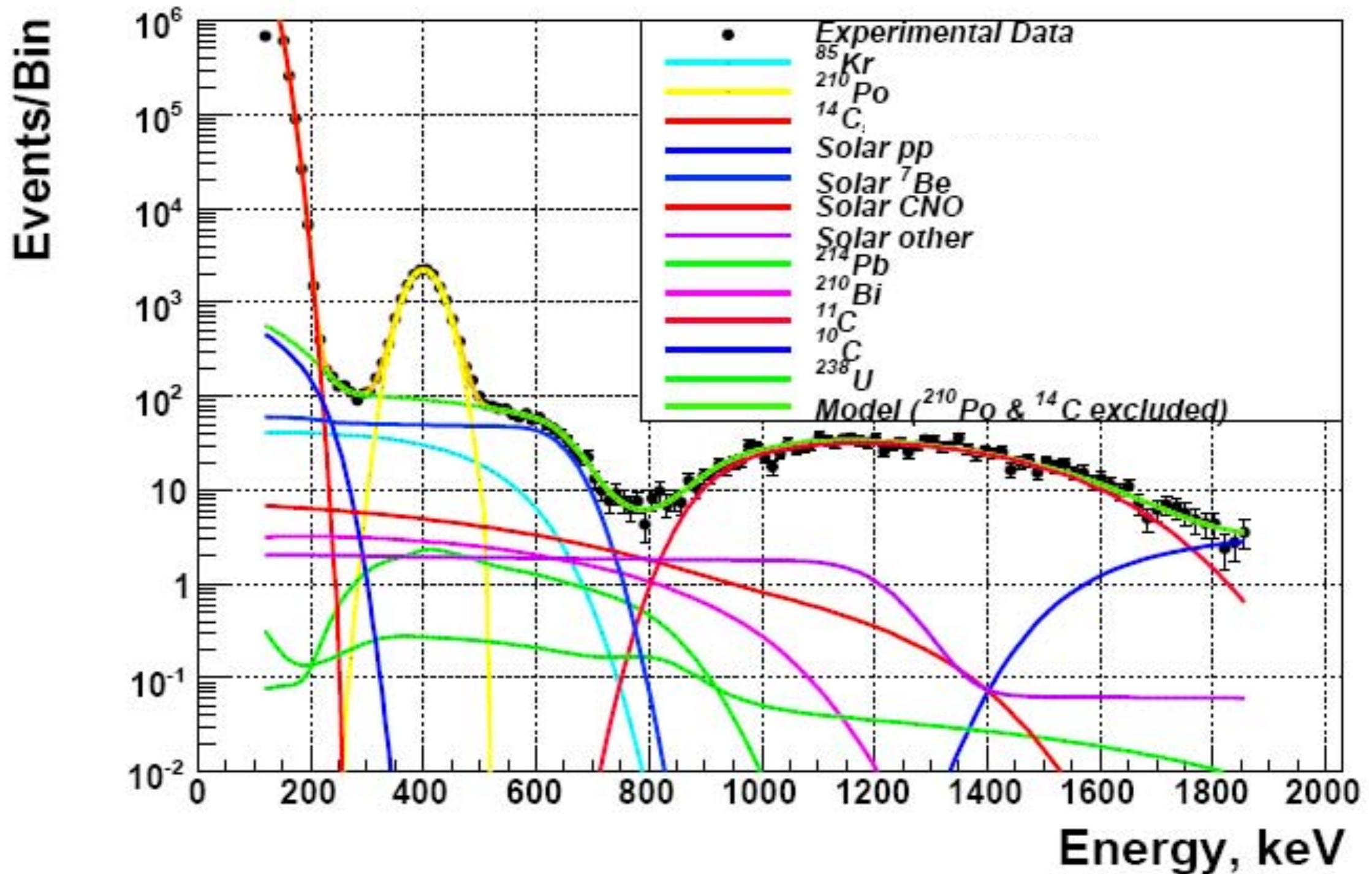


Radio-pure scintillator: Borexino

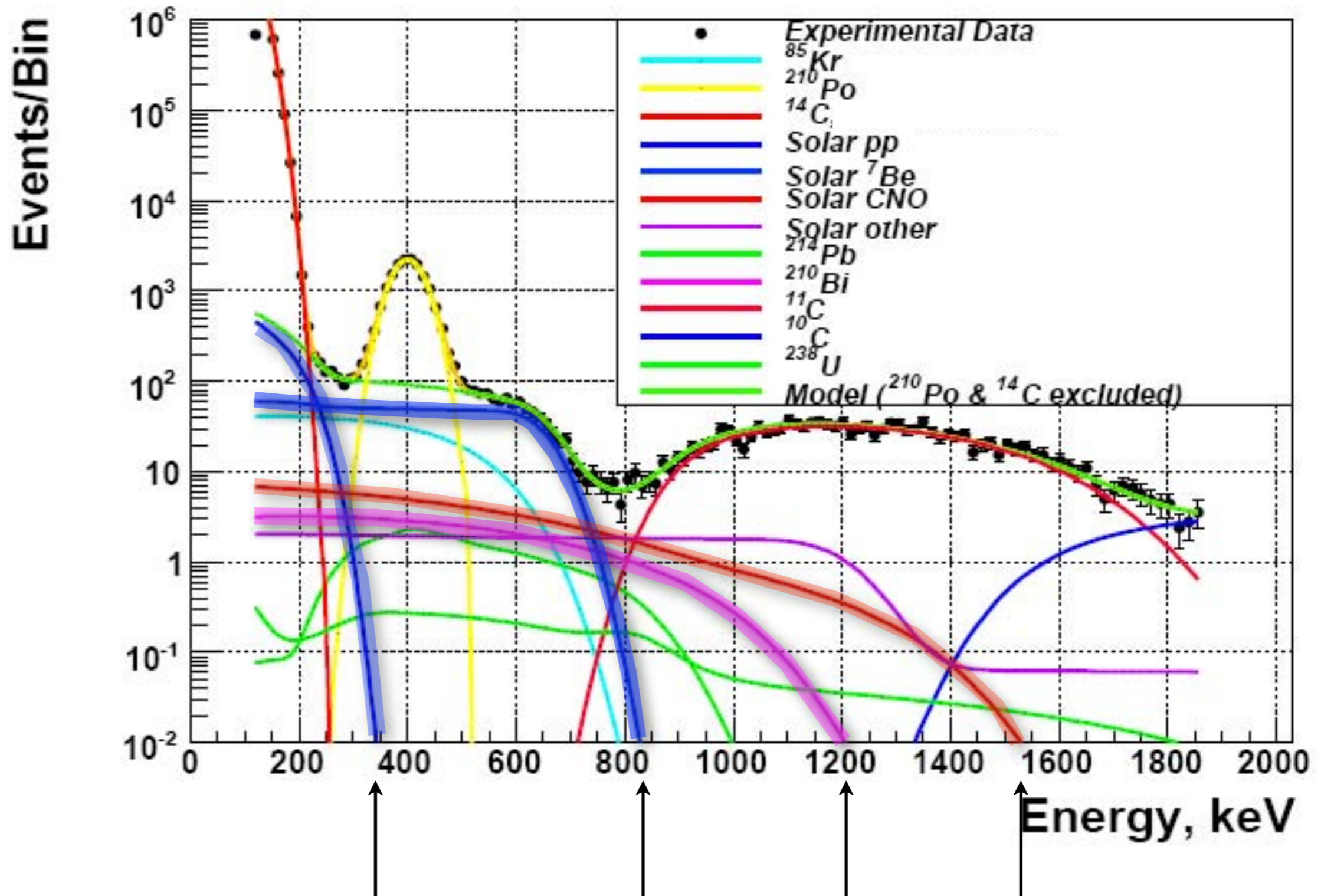
Background	Typical abundance (source)	Borexino goals	Borexino measured
$^{14}\text{C}/^{12}\text{C}$	10^{-12} (cosmogenic) g/g	10^{-18} g/g	$\sim 2 \cdot 10^{-18}$ g/g
^{238}U (by ^{214}Bi - ^{214}Po)	$2 \cdot 10^{-5}$ (dust) g/g	10^{-16} g/g	$(1.6 \pm 0.1) \cdot 10^{-17}$ g/g
^{232}Th (by ^{212}Bi - ^{212}Po)	$2 \cdot 10^{-5}$ (dust) g/g	10^{-16} g/g	$(5 \pm 1) \cdot 10^{-18}$ g/g
^{222}Rn (by ^{214}Bi - ^{214}Po)	100 atoms/cm ³ (air) emanation from materials	10^{-16} g/g	$\sim 10^{-17}$ g/g (~ 1 cpd/100t)
^{210}Po	Surface contamination	~ 1 c/d/t	May 07 : 70 c/d/t Sep08 : 7 c/d/t
^{40}K	$2 \cdot 10^{-6}$ (dust) g/g	$\sim 10^{-18}$ g/g	$< 3 \cdot 10^{-18}$ (90%) g/g
^{85}Kr	1 Bq/m ³ (air)	~ 1 c/d/100t	(28 ± 7) c/d/100t (fast coinc.)
^{39}Ar	17 mBq/m ³ (air)	~ 1 c/d/100t	$\ll ^{85}\text{Kr}$

Elastic scattering does not provide a double coincidence signal. No choice but to make extremely radio-pure scintillator

Radio-pure scintillator: Borexino

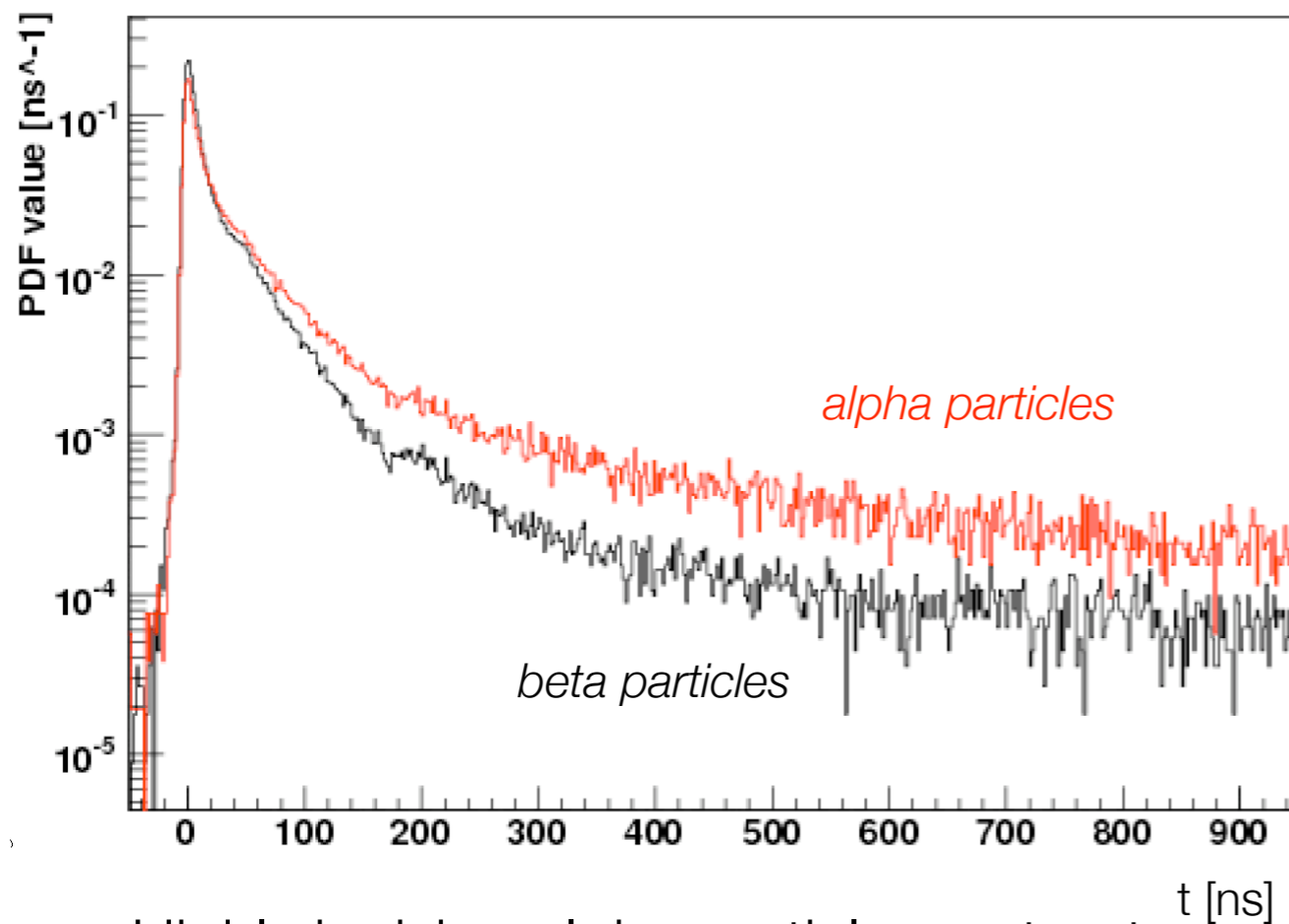


Radio-pure scintillator: Borexino

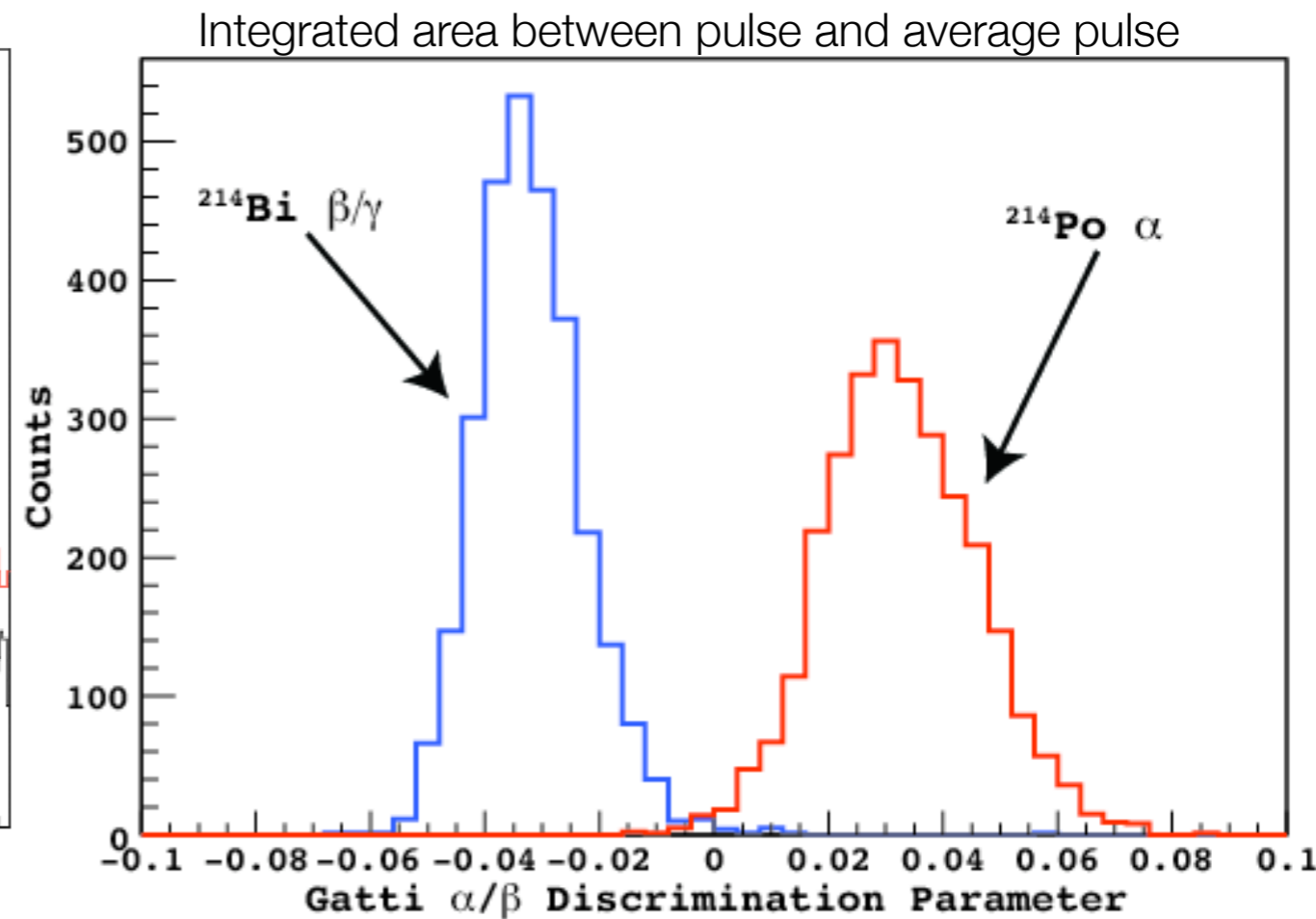


Alpha/Beta discrimination: Borexino

- Can remove events due to alpha emitters based on pulse shape analysis

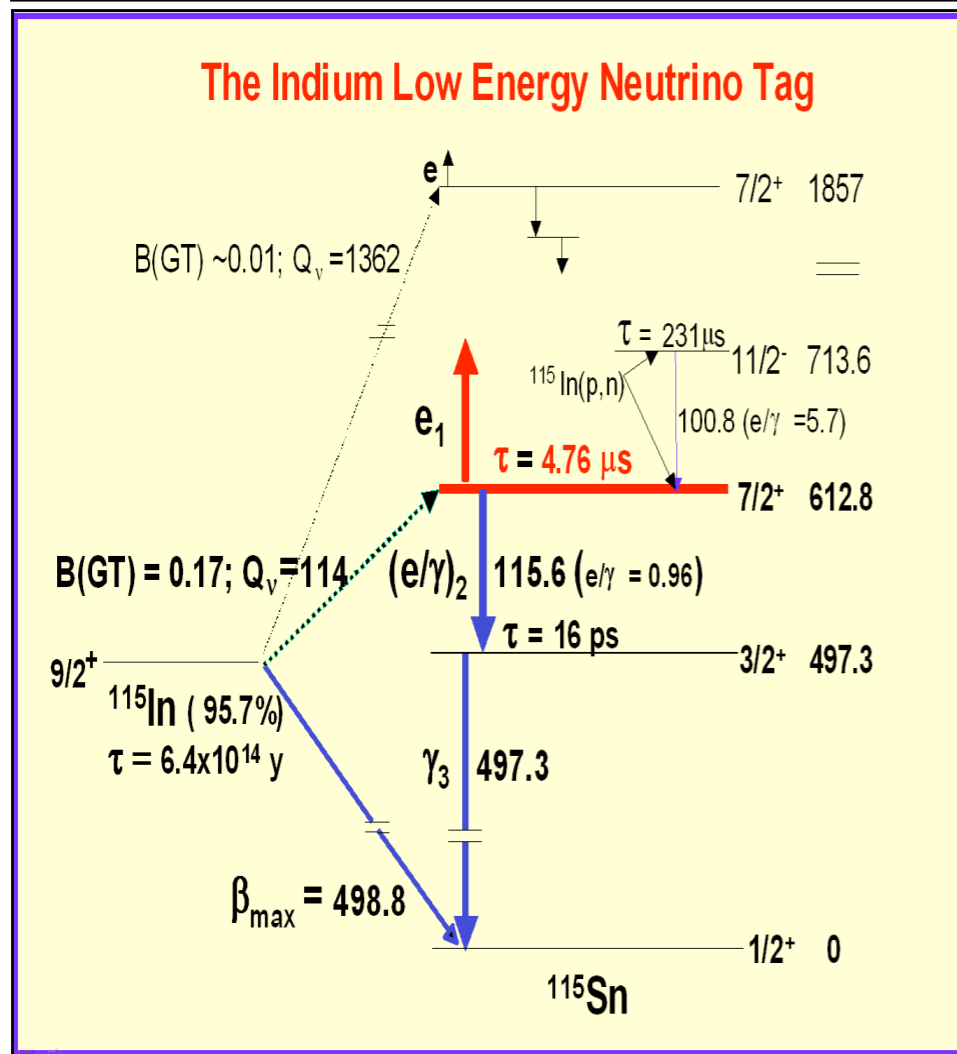
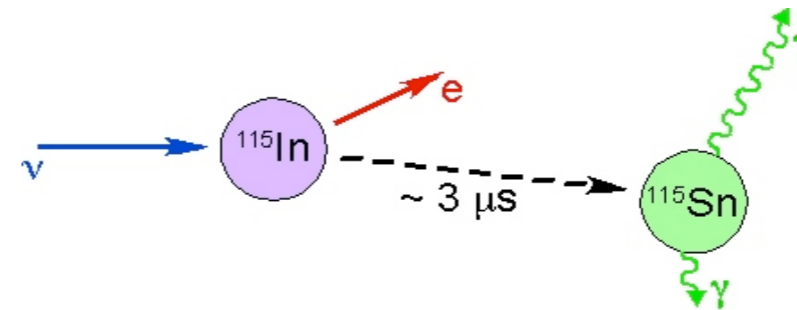
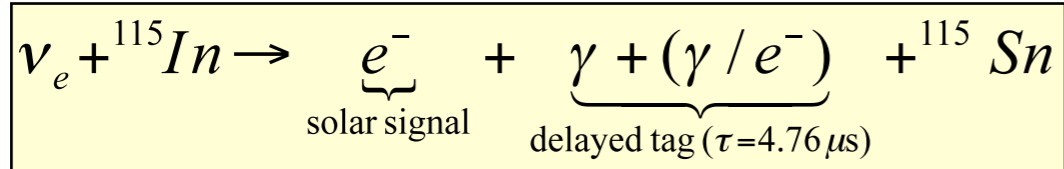


Highly ionizing alpha particles saturate more of the scintillator and have a longer decay time



Even lower? LENS

- The LENS experiment proposed to use scintillator loaded with Indium¹¹⁵

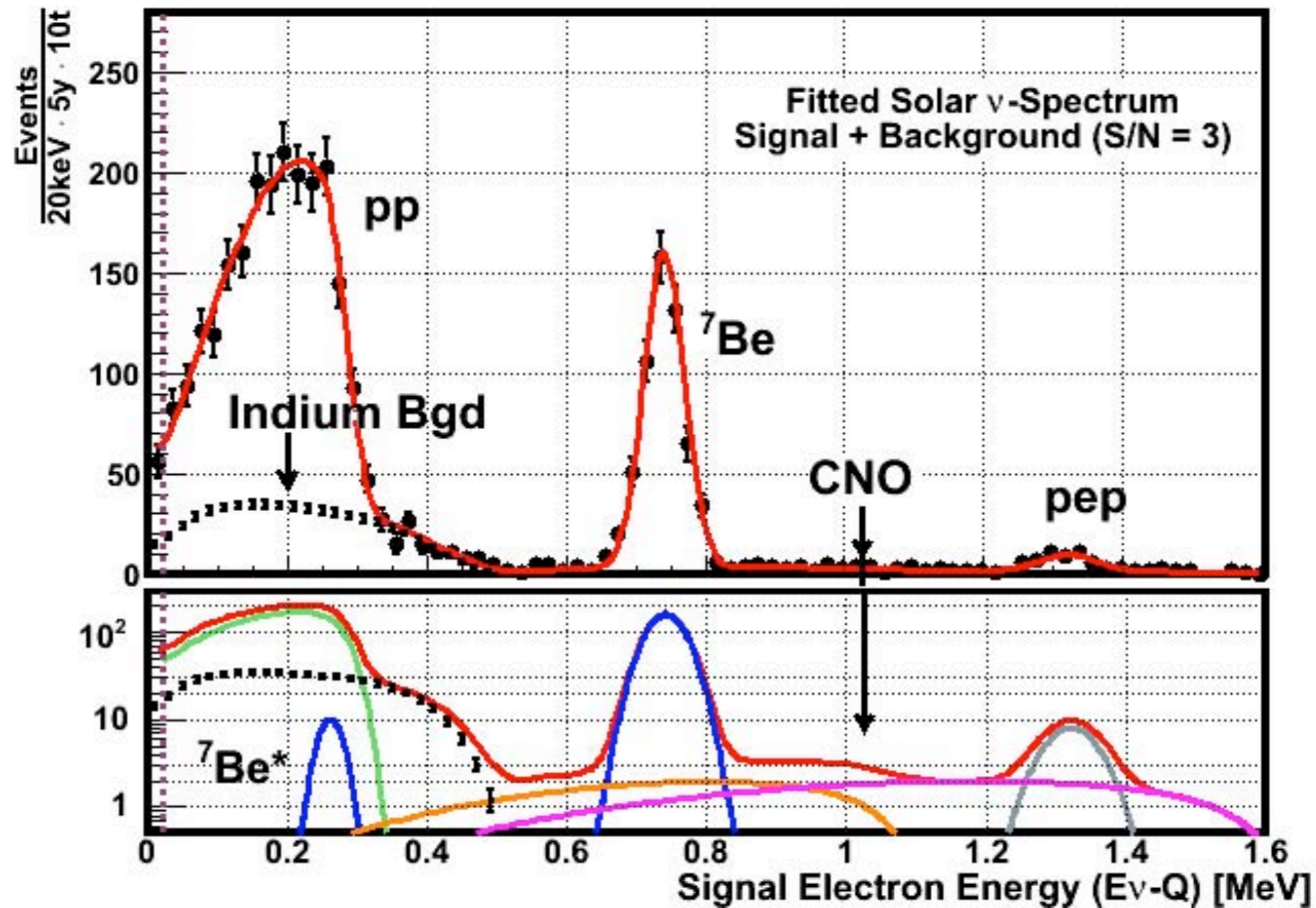


Signature is a triple coincidence in space and time



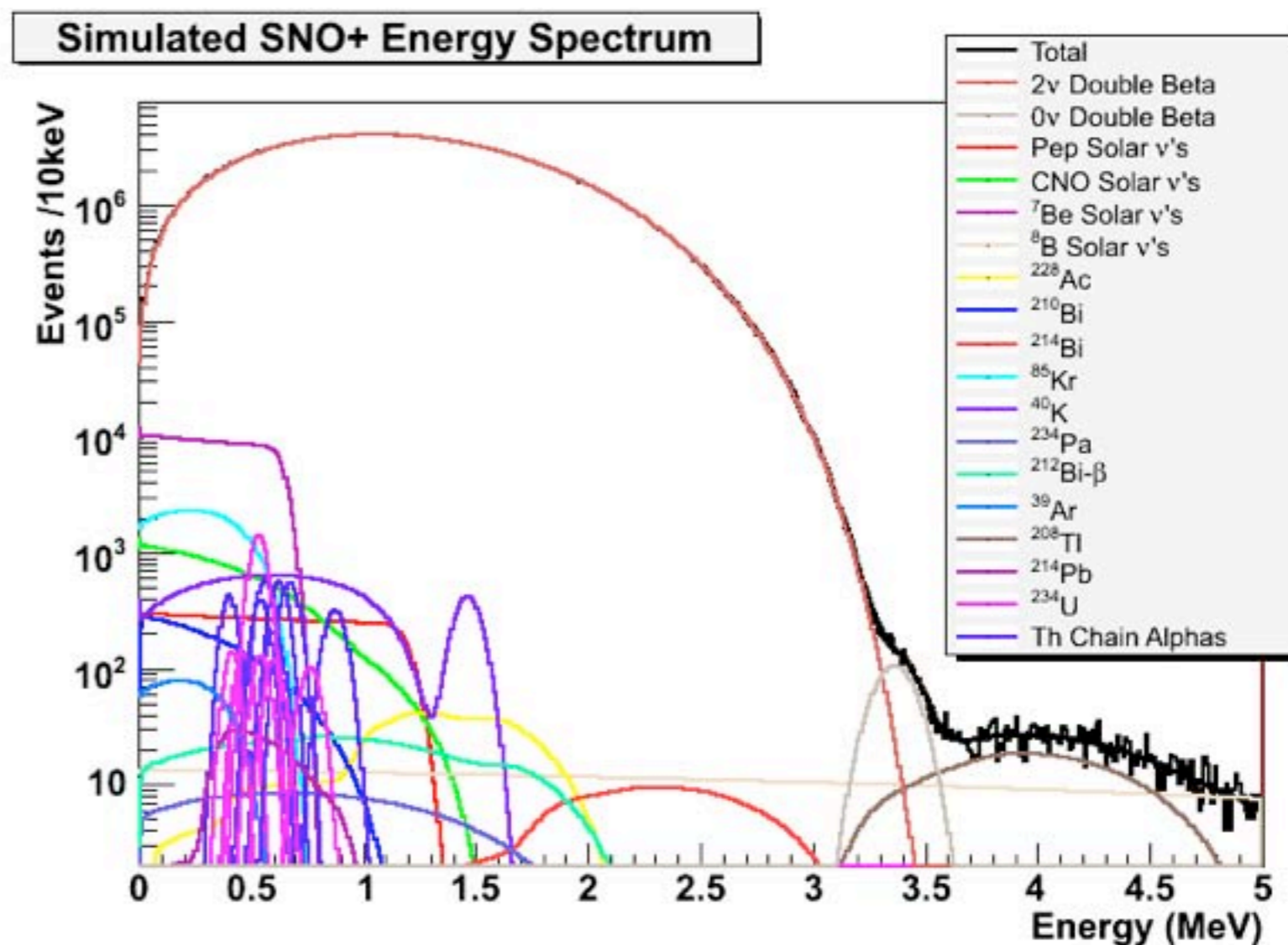
Segment detector and guide light to walls via complete internal reflection

Simulated solar neutrino spectrum in LENS



Liquid scintillator for $0\nu\beta\beta$: SNO+

- Reuse SNO cavern but using liquid scintillator instead of D_2O
- Load scintillator with ^{150}Nd
- Overwhelm relatively poor energy resolution with statistics. SNO+ could hold as much as 500 kg of source. Compare to current largest $0\nu\beta\beta$ experiment which uses roughly 10 kg of isotope.

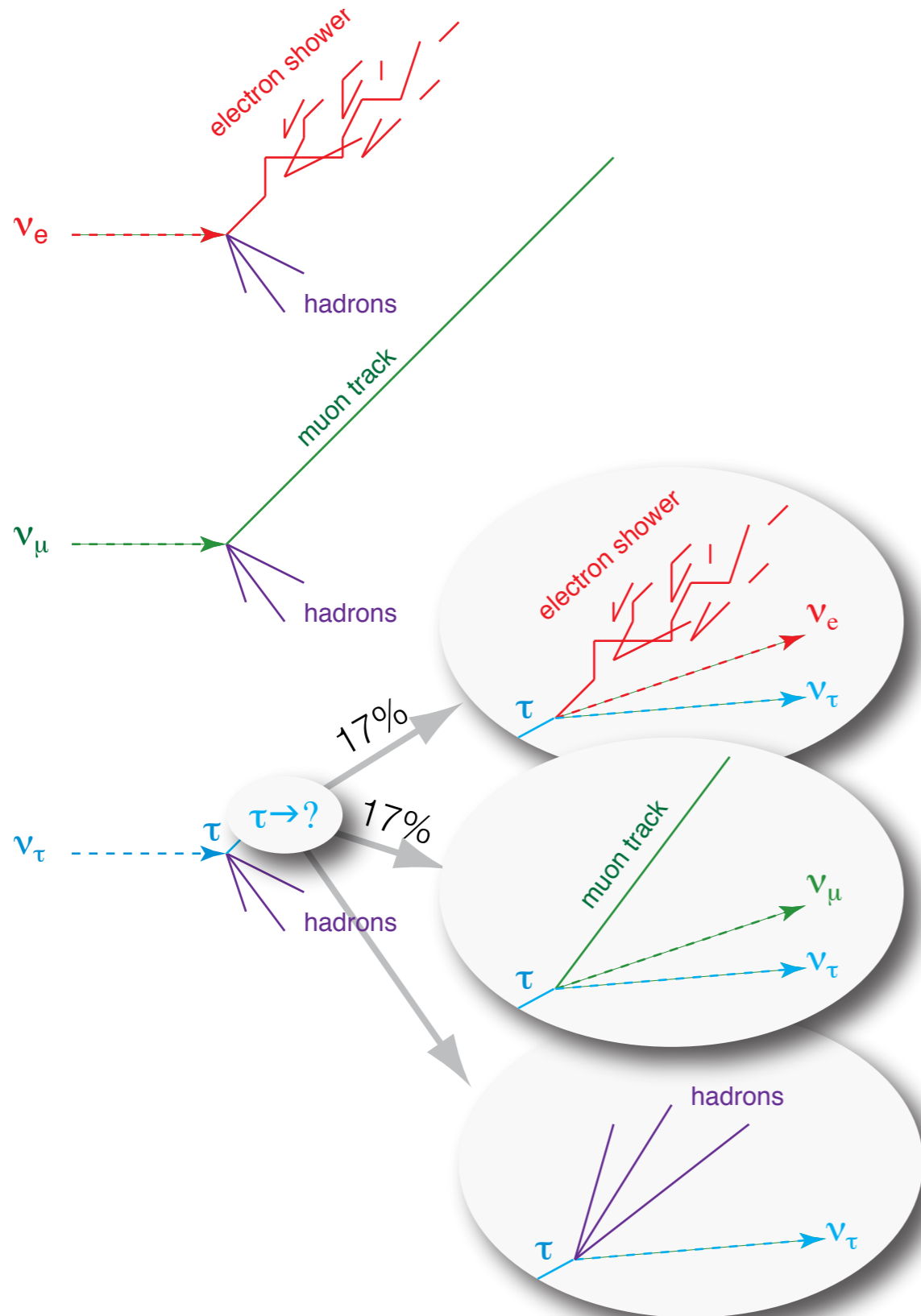


$0\nu\beta\beta$ signal and backgrounds at lower edge of reported observation by Klapdor-Kleingrothaus et al.

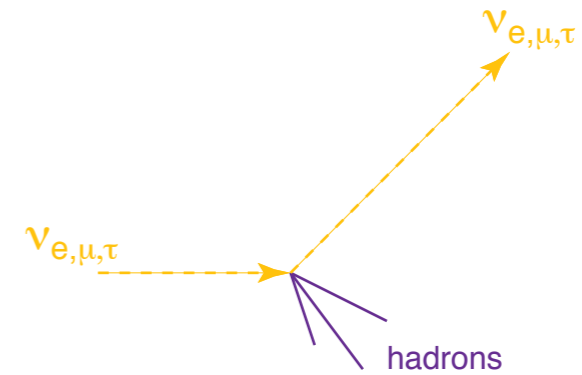
Tau Neutrino Detection

Neutrino detection channels

Charged-current



Neutral-current



- In charged-current (CC) events outgoing lepton tags incoming neutrino flavor.
 - ▶ In the case of ν_τ , the presence of a τ must be deduced from the τ decay products
- In CC events nearly all the neutrino energy is deposited in the detector
- In neutral-current events, only hadrons are present and no information about the incident neutrino flavor is available
- CC rates are affected by oscillations
- NC rates are not affected by oscillations
 - ▶ In only a few analyses are NC events considered to be signal. In most cases NC events are backgrounds to the CC processes

Tau neutrinos

- Tau neutrinos are difficult to observe
 - *They are difficult to produce.* First direct observation (DONUT) was via decays of charmed particles in a beam dump.
 - *They are difficult to make interact.* Threshold for tau production is 3.5 GeV. This puts them above the oscillation maximum for most beams designed to study oscillations at the atmospheric mass-squared scale. For example, for $L=735$ km, $E_{\text{max}} = 1.5$ GeV, which is below threshold
 - *They are difficult to detect.* The lifetime of the tau is 291 fs; Even when highly boosted, decay length is only a few mm. Required a very finely segmented vertex region.

$$\begin{aligned}c\tau_{\text{LAB}} &= (3 \times 10^{11} \text{ m/s})(290 \times 10^{-15} \text{ s})\left(\frac{E}{1.7 \text{ GeV}}\right) \\ &= 0.5 \text{ mm} \times E [\text{GeV}]\end{aligned}$$

- Tau neutrinos produce backgrounds to electron neutrino searches:

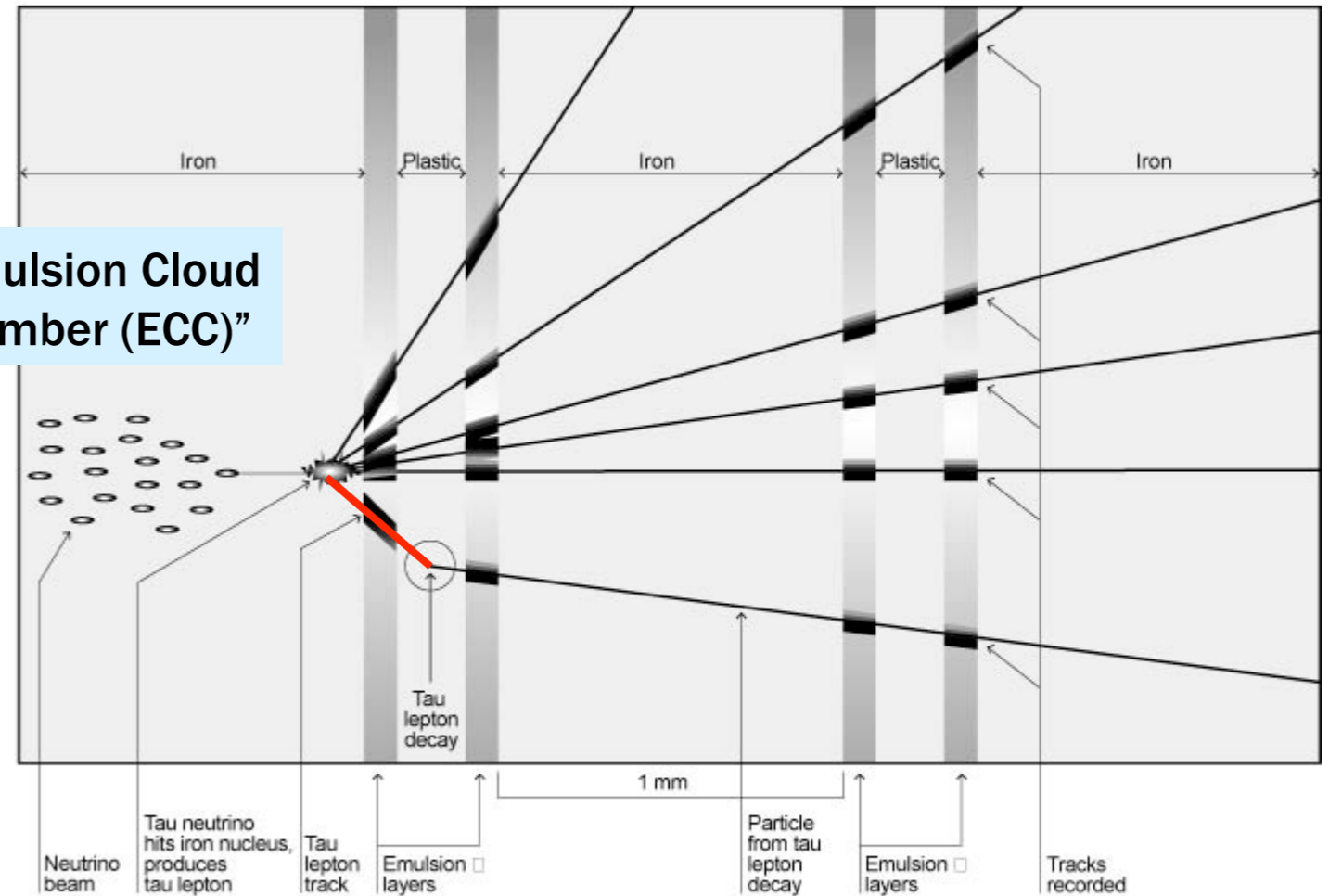
$$\tau \rightarrow e \nu_e \nu_\tau$$

Tau Neutrino Detection

- Several experiments look for tau neutrinos
- Observed by DONUT experiment
- Sought from oscillations by CHORUS and OPERA
- All of the above experiments have used thin films of photographic emulsions placed between target layers
- Use of emulsion allows for resolution of short tau track and search for its decay either through a track kink or to multi-prongs
- Emulsion target followed by other detectors which provide tracking and tell you where you had a neutrino interaction and which emulsions you should develop

Detecting a Tau Neutrino

“Emulsion Cloud Chamber (ECC)”



Of one million million tau neutrinos crossing the DONUT detector, scientists expect about one to interact with an iron nucleus.

Layout of the DONUT experiment

Tracking using scintillating fibers interspersed in emulsion target and drift chambers used to locate vertices in emulsion with \sim mm accuracy

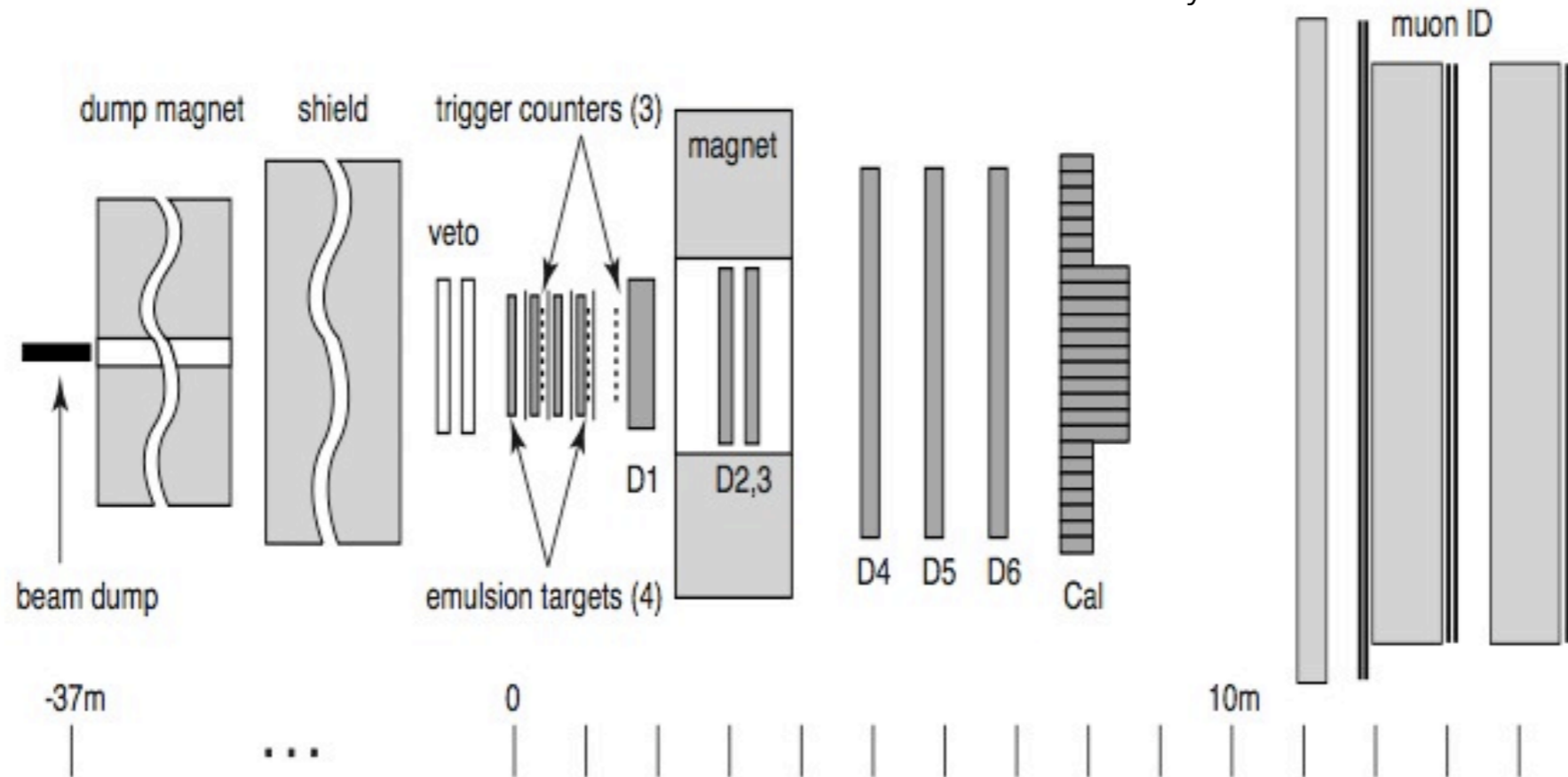


Fig. 1. Experimental beam and spectrometer. At the left, 800 GeV protons were incident on the beam dump, which was 36.5m from the first emulsion target. Muon identification was done by range in the system at the right (downstream).

Tau Neutrino Detection by DONUT Collaboration

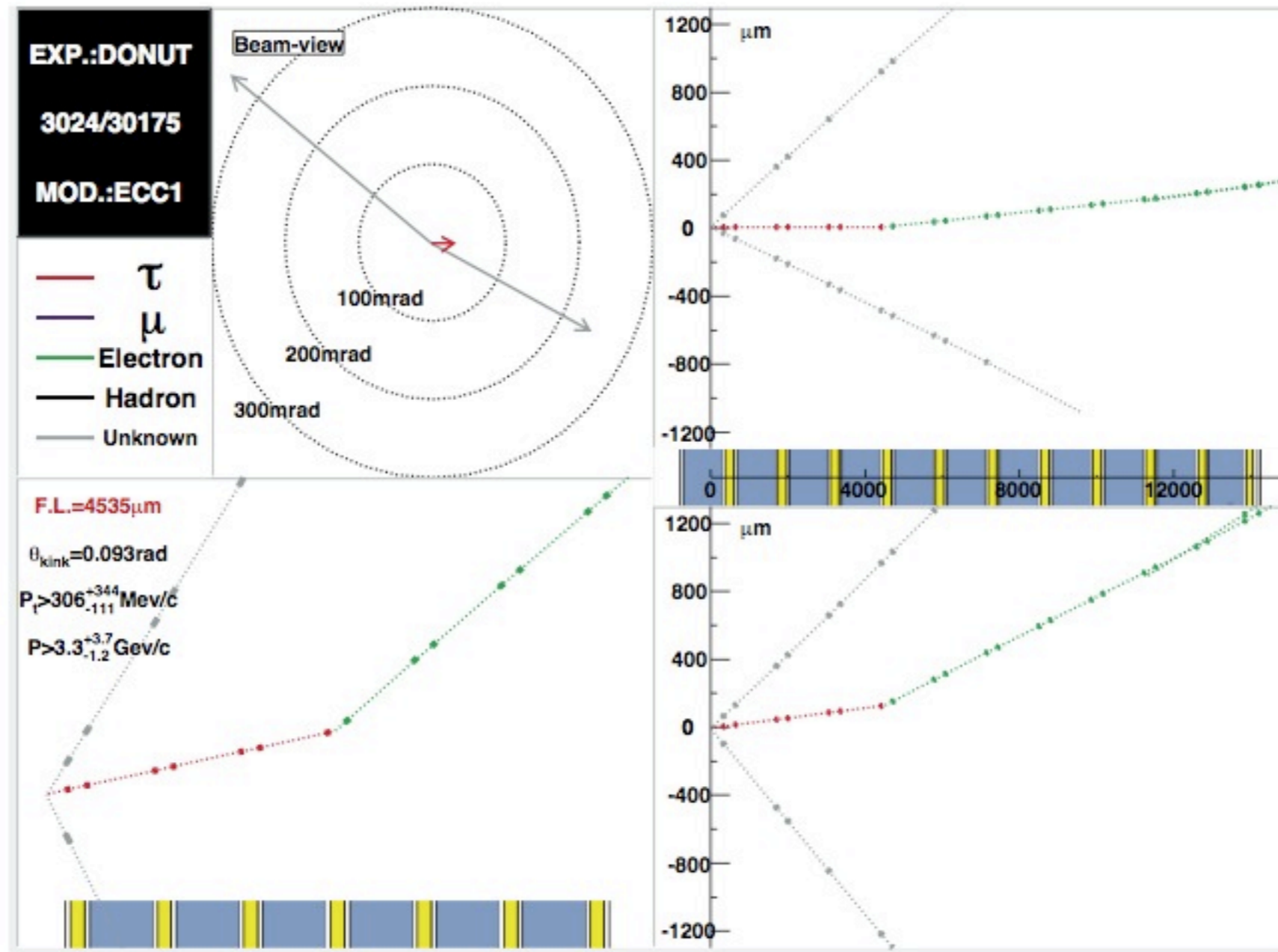
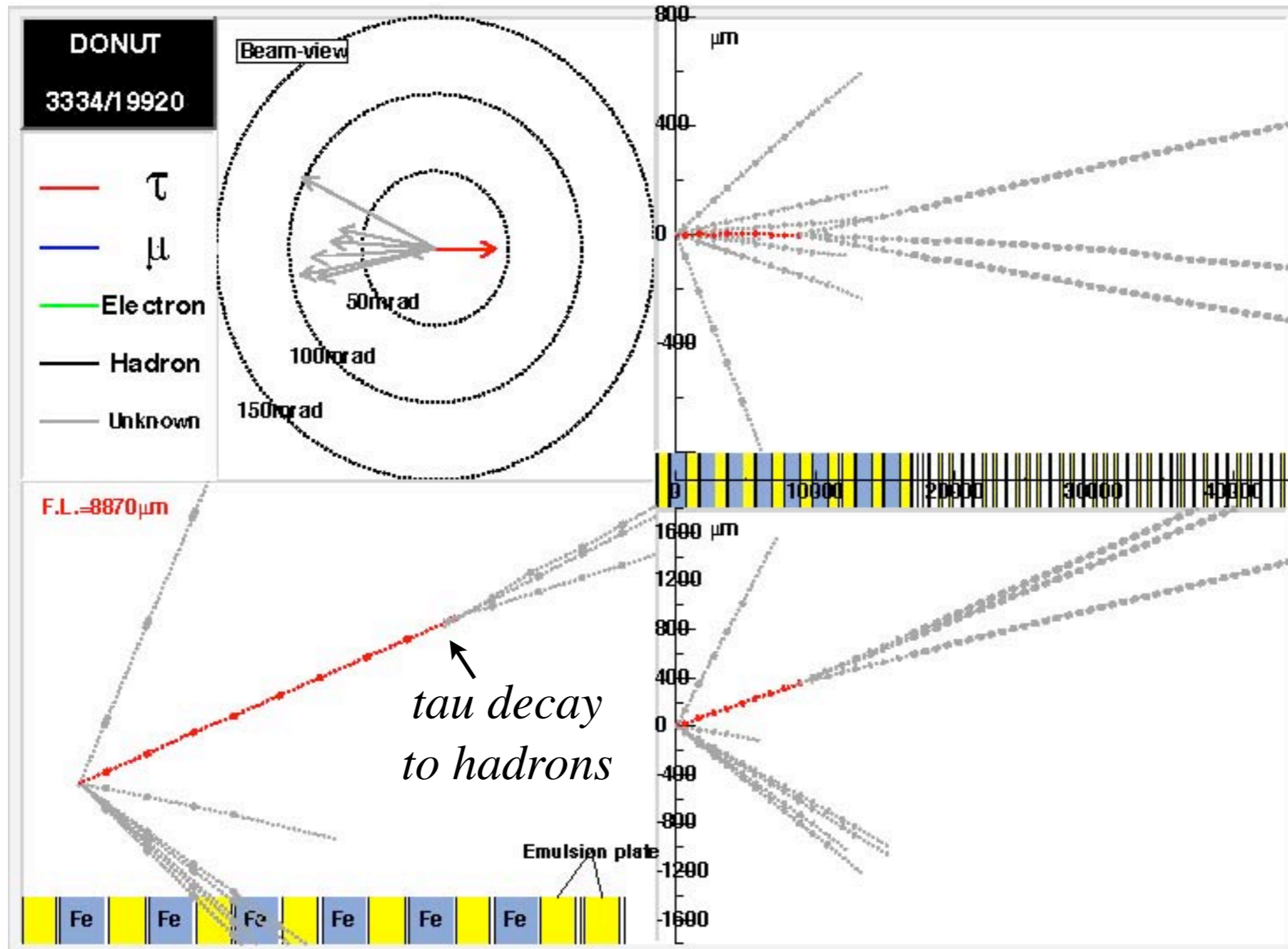


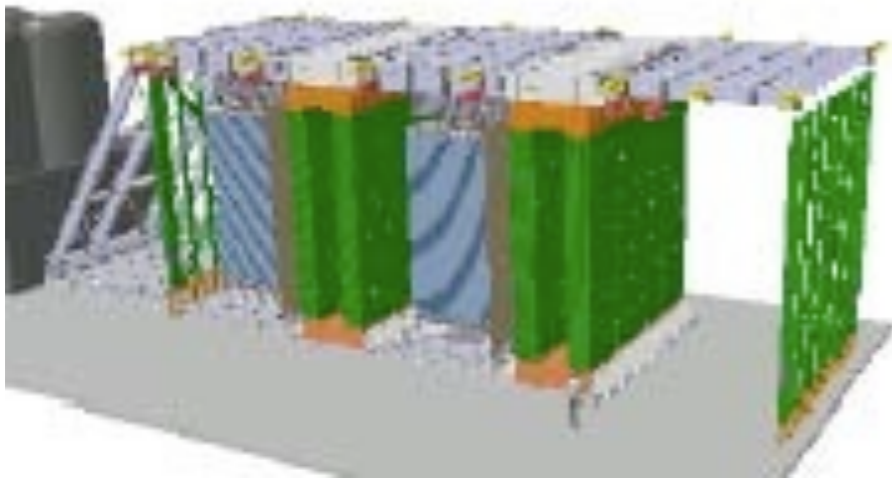
Fig. 18. Event 3024-30175: $\tau \rightarrow e + \nu_{\tau} + \nu_e$.

Tau Neutrino Detection by DONUT Collaboration



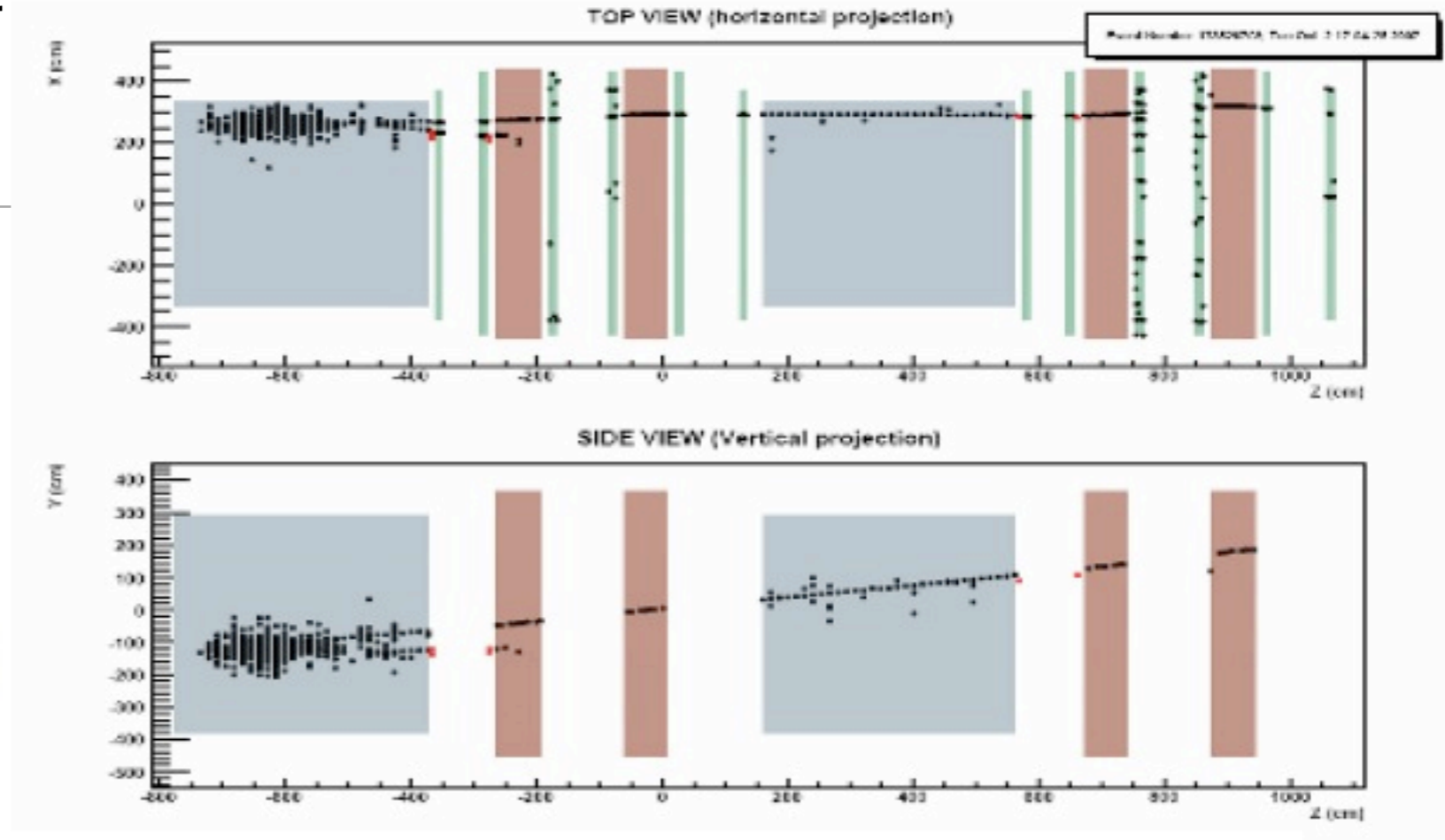
OPERA Experiment In CNGS beam

OPERA uses bricks of lead/
emulsion embedded in a solid
scintillator-based tracking
system + downstream muon
spectrometer

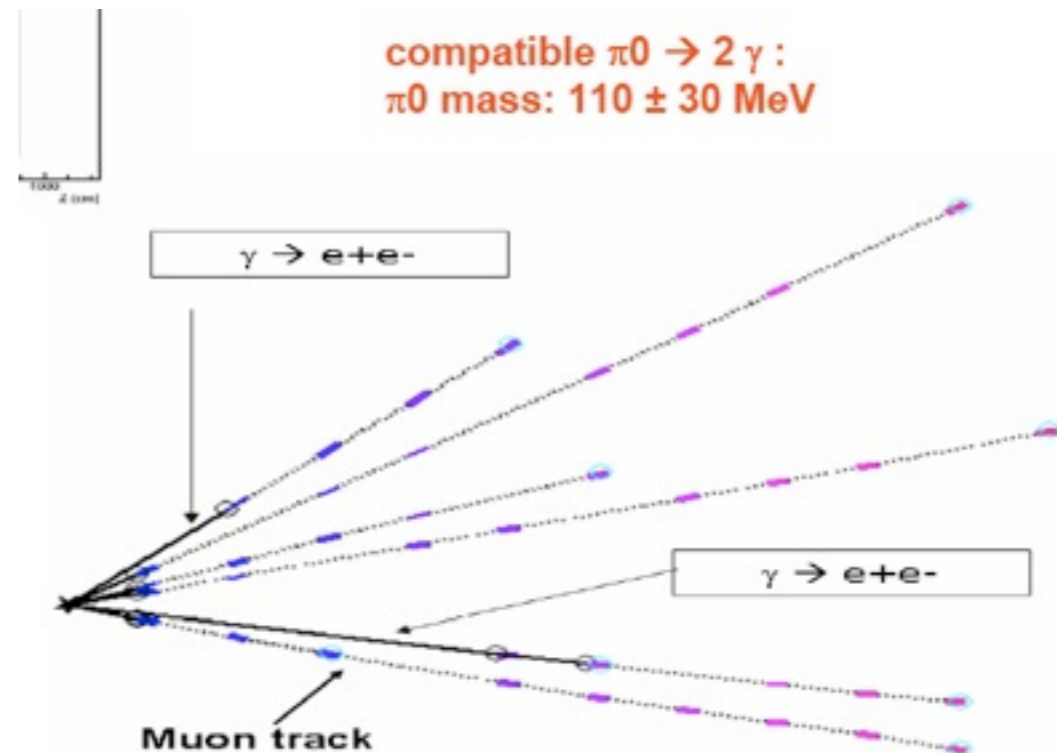


τ^- decay channels	Signal $\div \Delta m^2$ (Full mixing)		Background
	2.5×10^{-3} (eV ²)	3.0×10^{-3} (eV ²)	
$\tau^- \rightarrow \mu^-$	2.9	4.2	0.17
$\tau^- \rightarrow e^-$	3.5	5.0	0.17
$\tau^- \rightarrow h^-$	3.1	4.4	0.24
$\tau^- \rightarrow 3h$	0.9	1.3	0.17
ALL	10.4	15.0	0.76

First event!

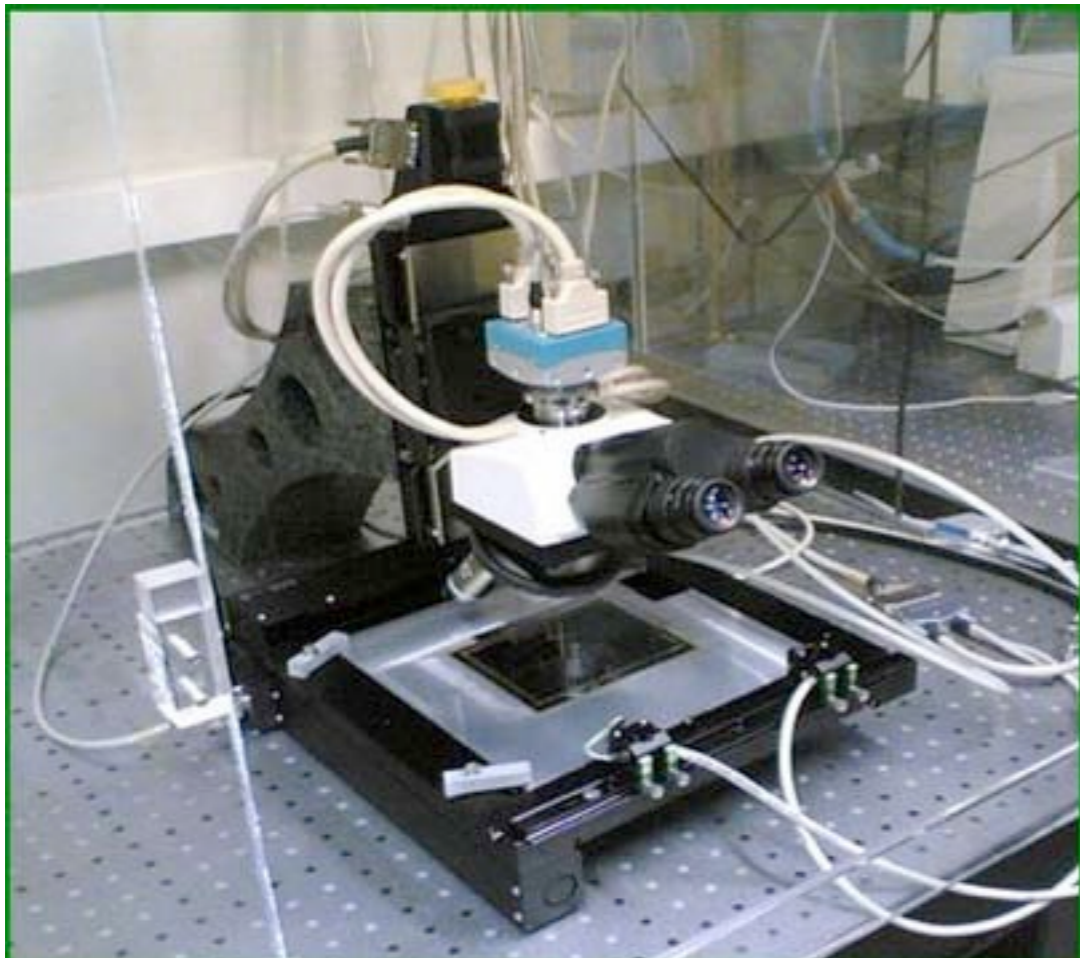


compatible $\pi^0 \rightarrow 2\gamma$:
 π^0 mass: 110 ± 30 MeV

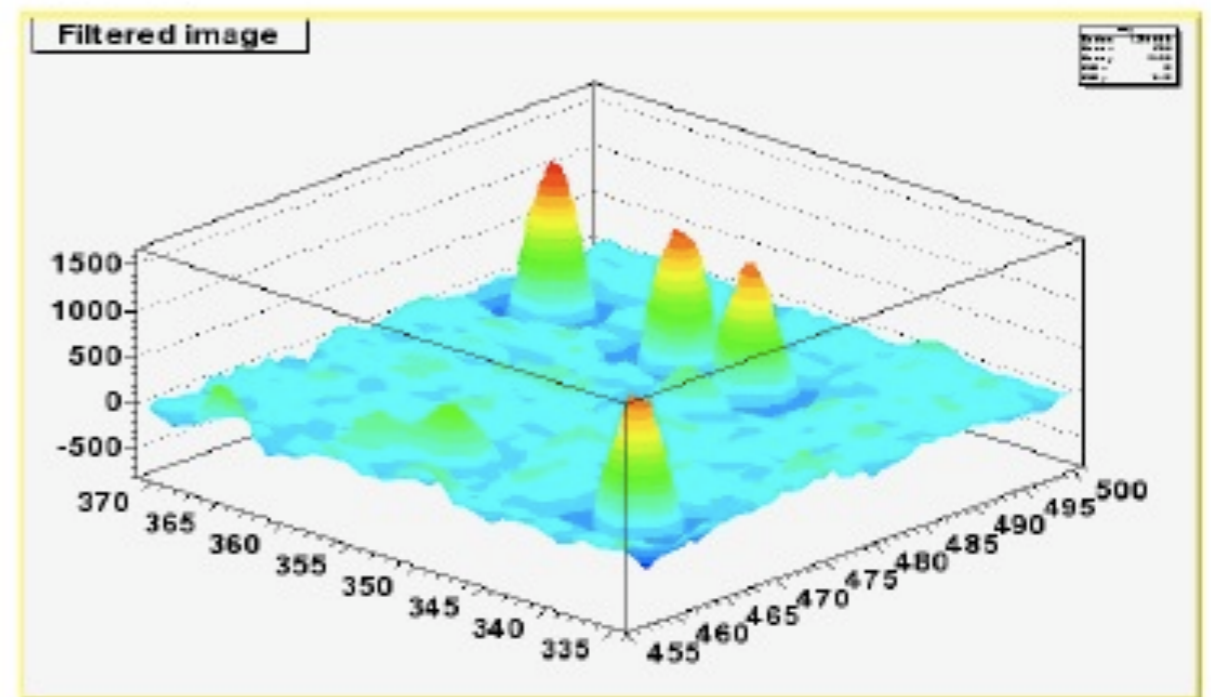


Scanning emulsion layers

Tracking system is used to identify bricks with interesting events. Those bricks are removed from the detector and the emulsion layers are scanned and digitized.



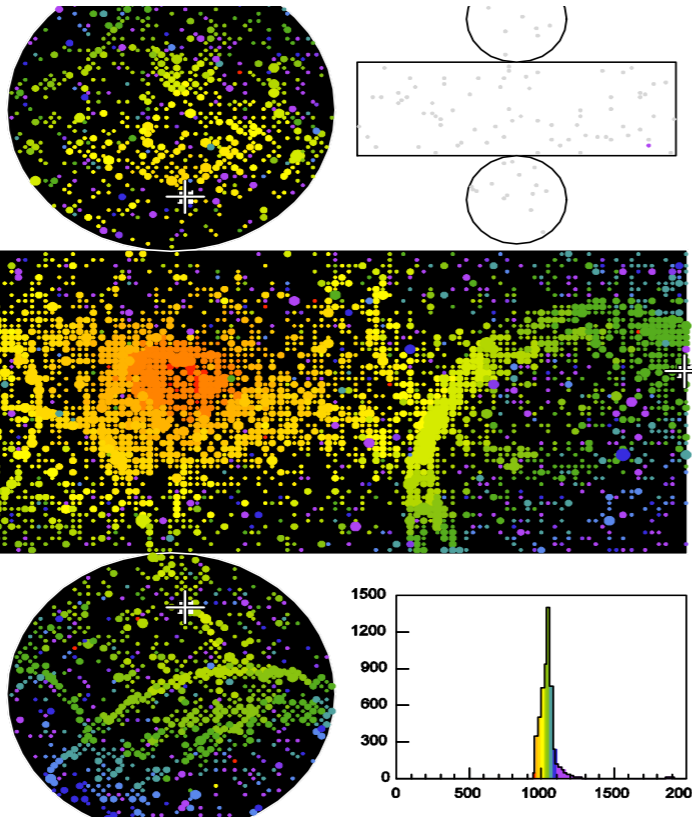
Deep inside view - image



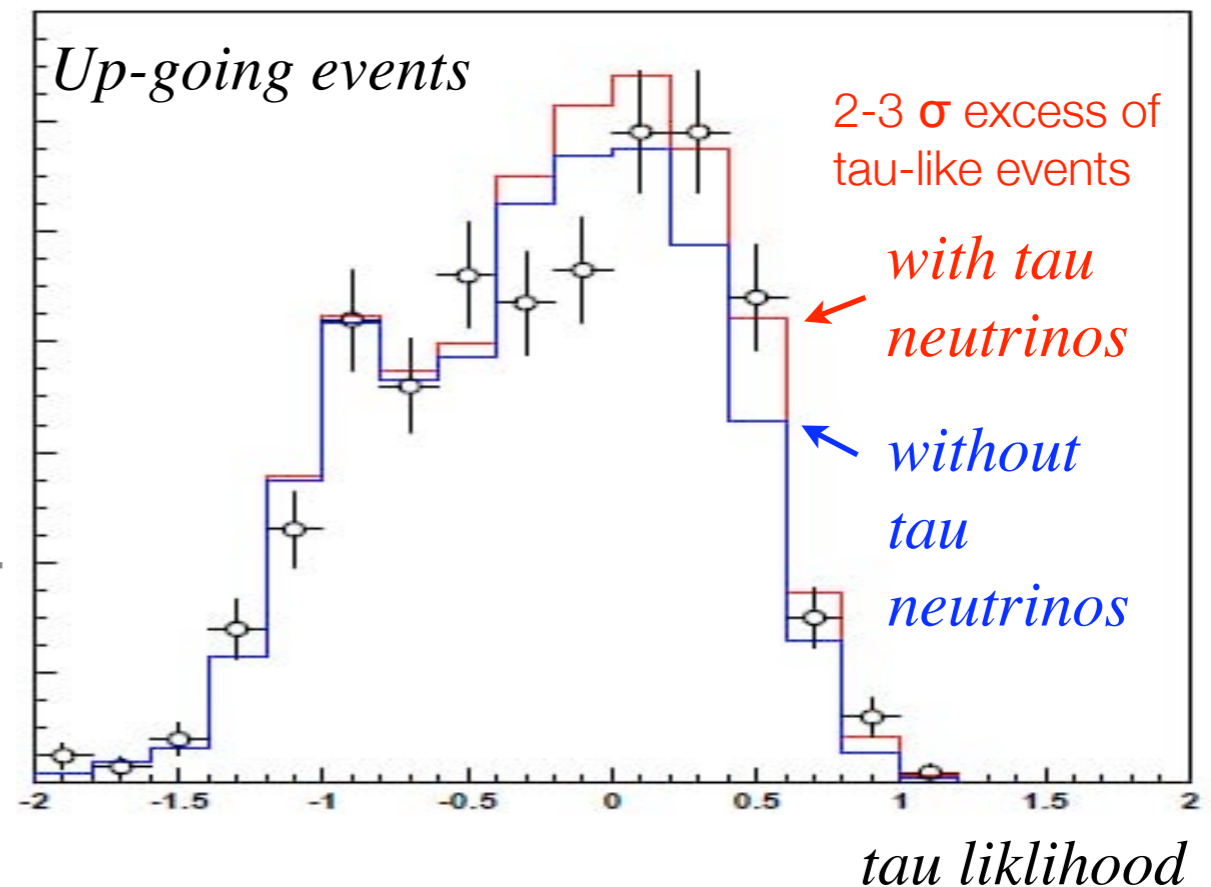
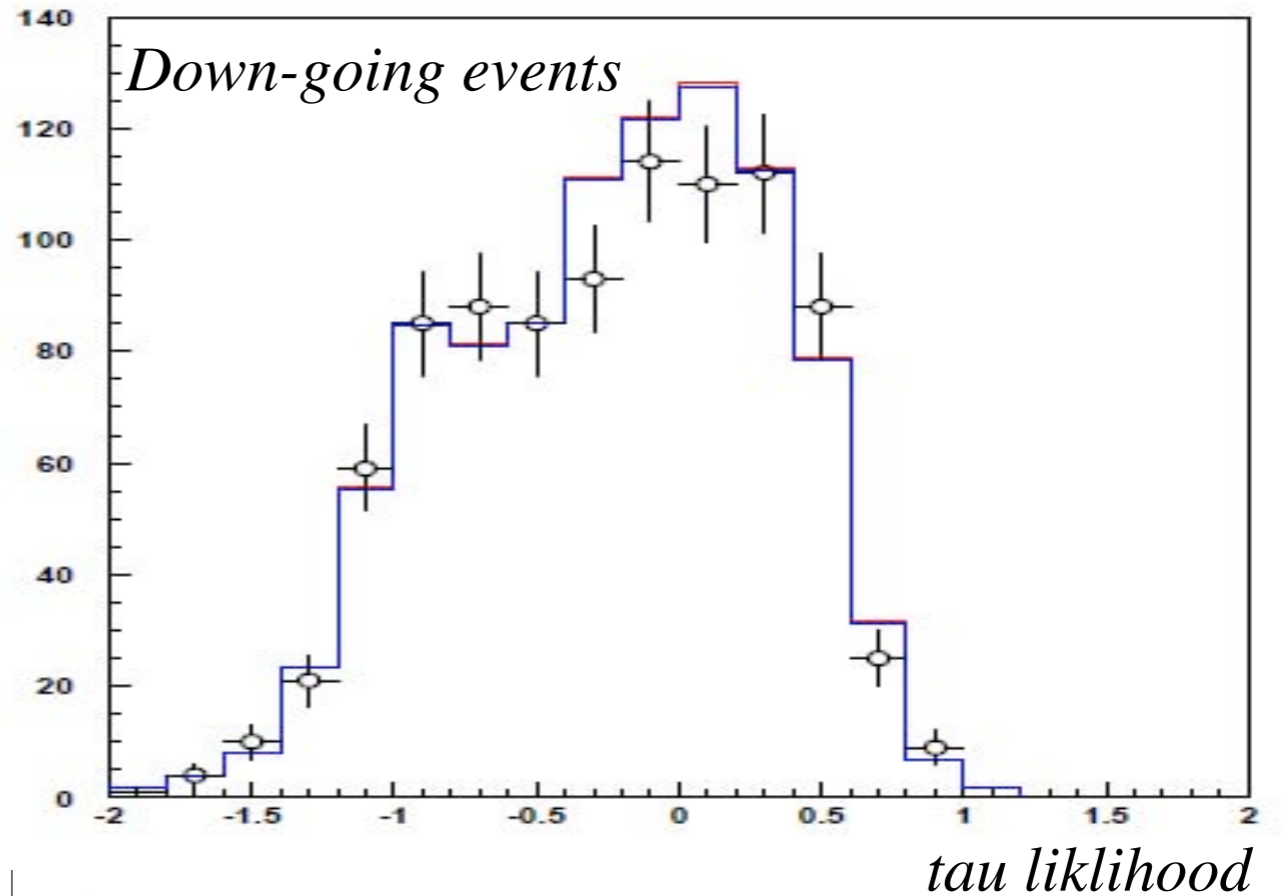
Four tracks in single emulsion layer

Statistical Tau Appearance

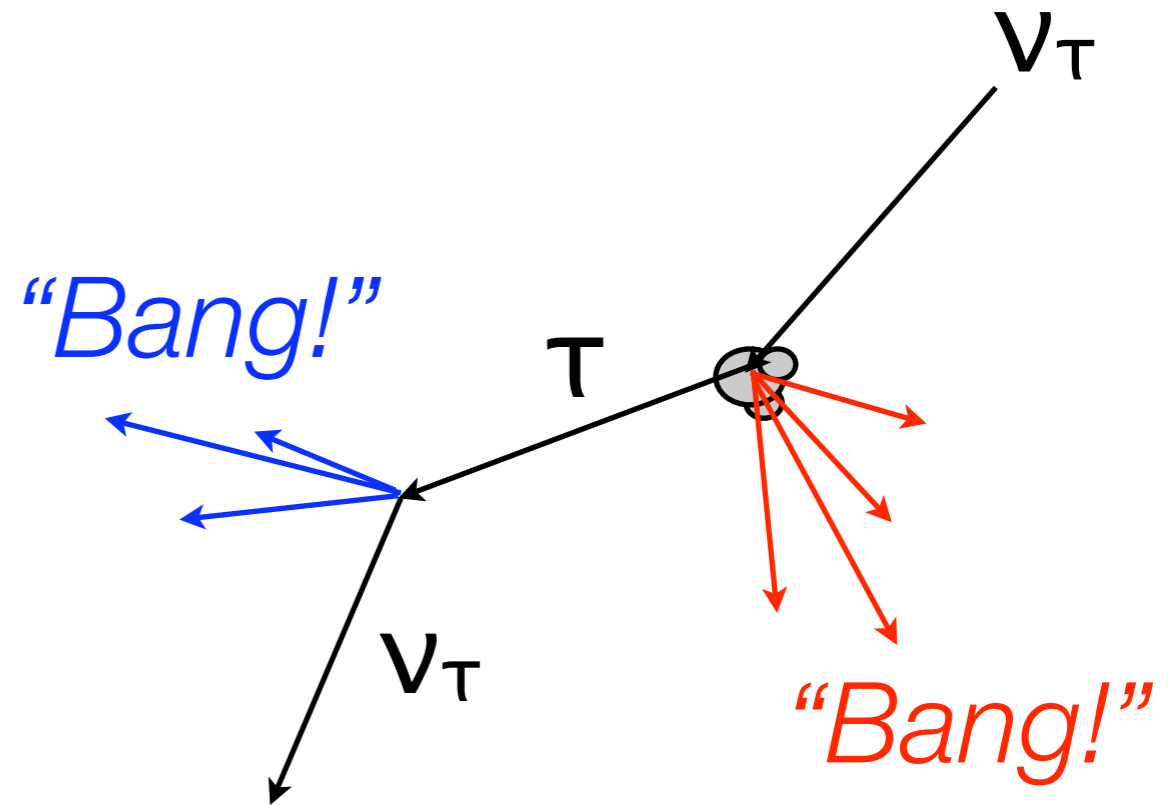
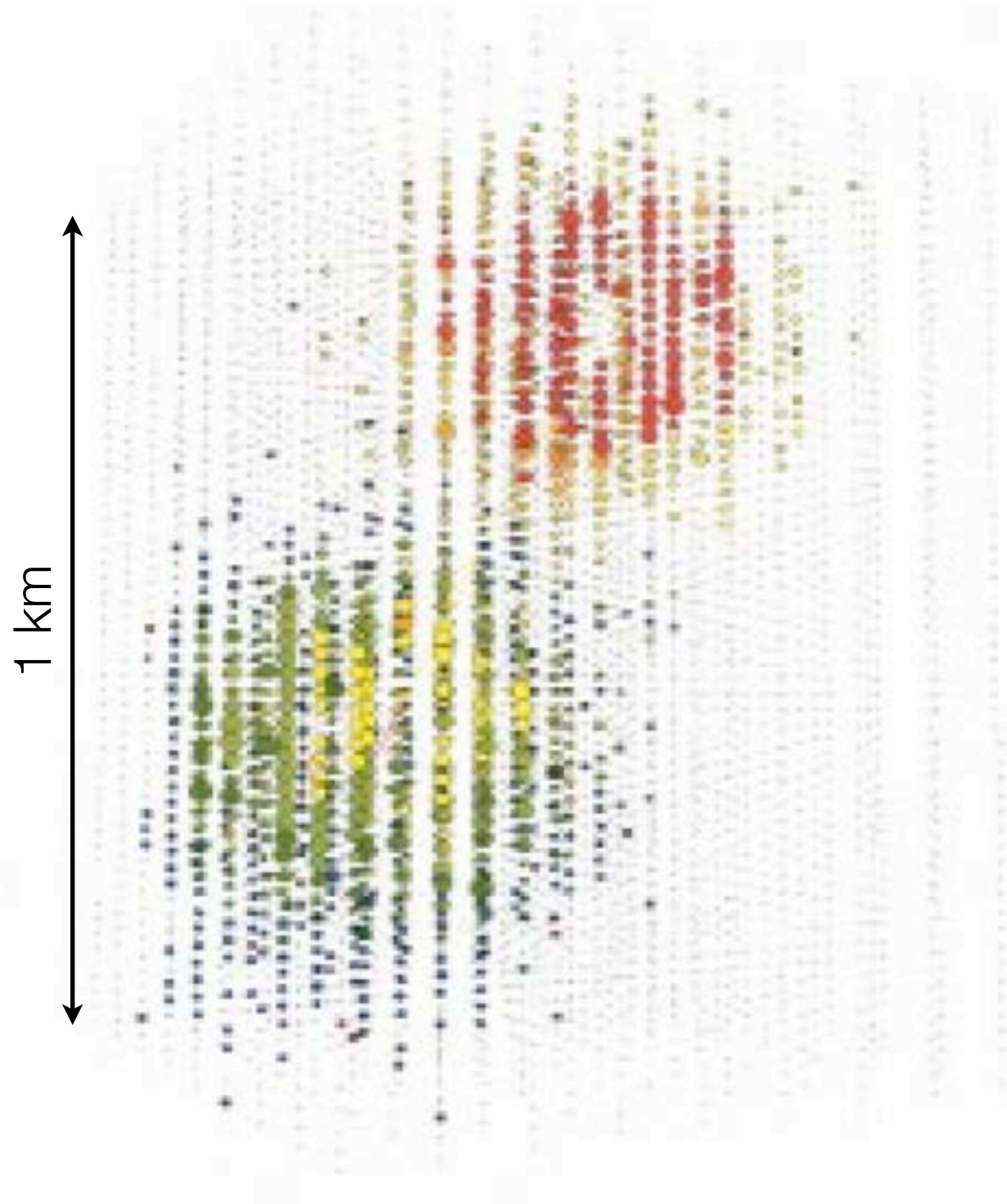
```
Event 30
9:03
its, 14223 pE
s, 0 pE (in-time)
κ03
ed
```



While large detectors may not be able to identify tau neutrino events one-by-one, they may be able to separate tau neutrino events from other events statistically. In SK hadrons from tau decay travel about 1 interaction length (80 cm) and scatter. Produces large number of thin rings.



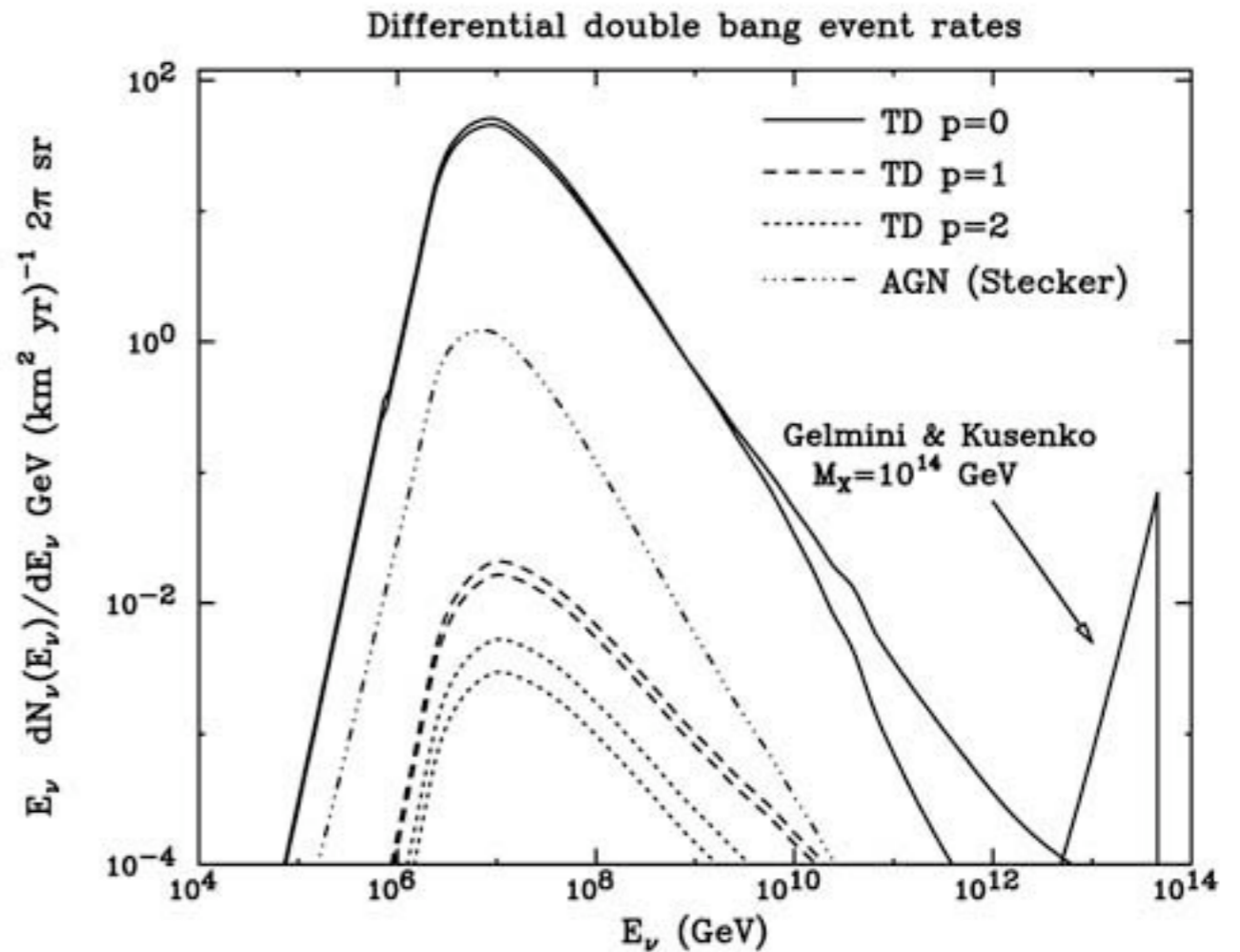
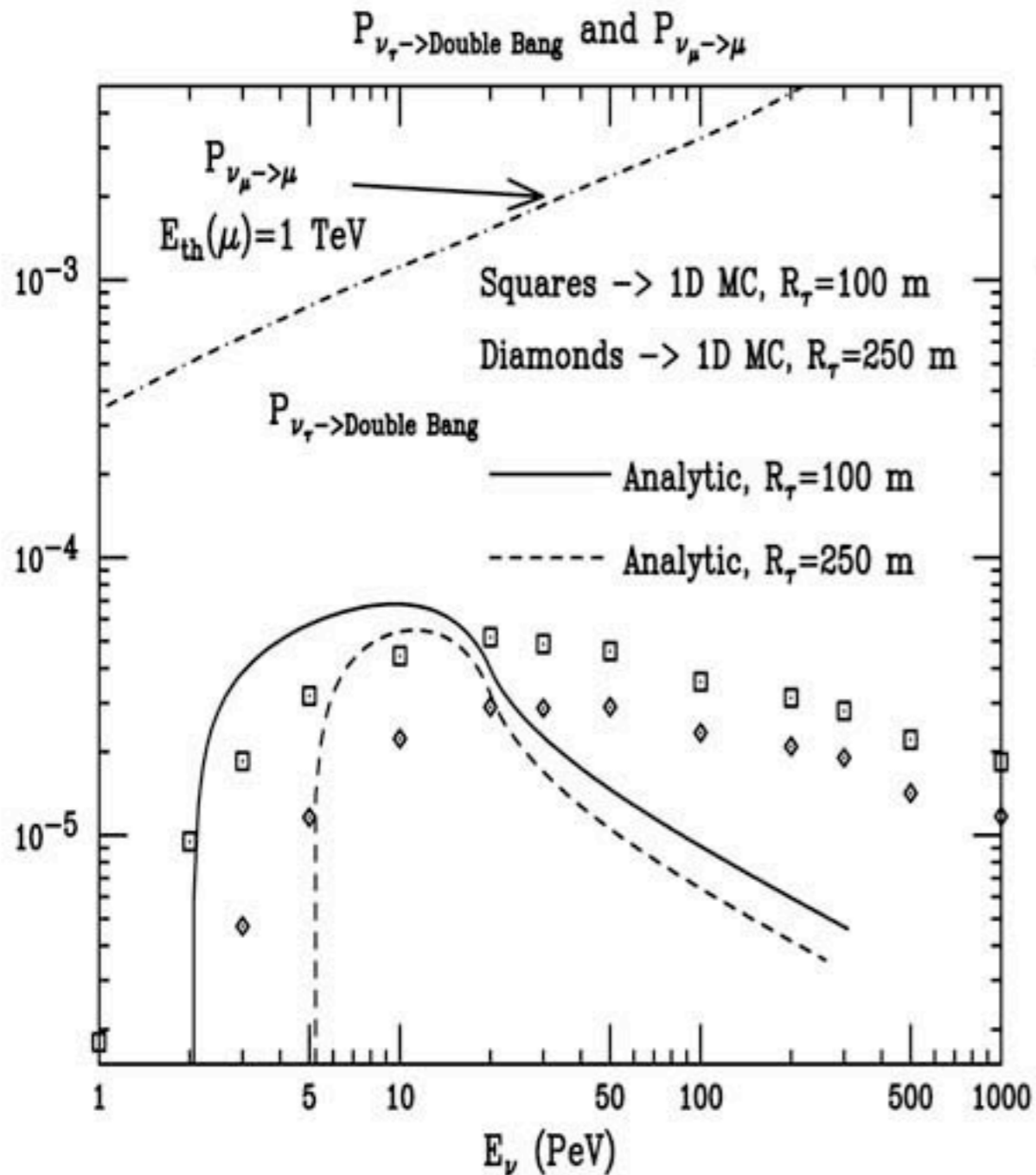
Tau neutrinos in Ice Cube: Double bang



Tau decay length:

$$\begin{aligned} c\tau_{\text{LAB}} &= (3 \times 10^8 \text{ m/s})(290 \times 10^{-15} \text{ s})\left(\frac{E}{1.7 \text{ GeV}}\right) \\ &= 5 \times 10^{-5} \text{ m} \times E[\text{GeV}] \\ &= 50 \text{ m for } E = 1 \text{ PeV} \end{aligned}$$

Tau neutrinos in Ice Cube: Double bang



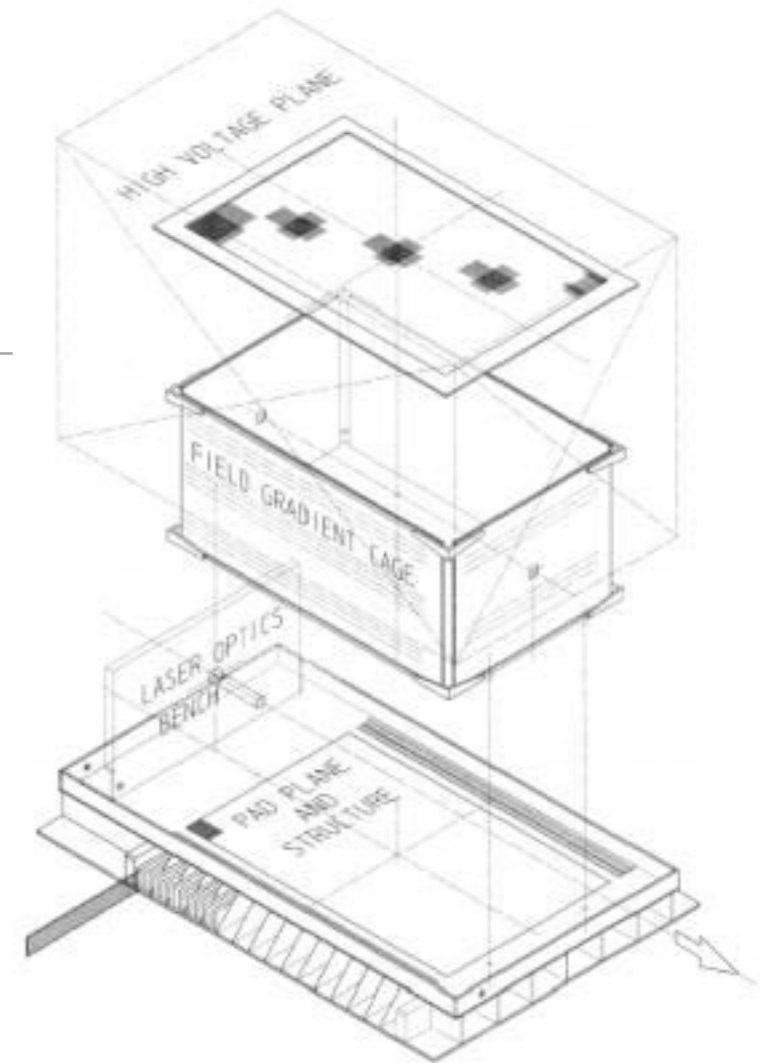
Integrates to ~ 0.01 or 100 events per year for models of topological defects.
 ~ 1 / year from AGNs.

Probability for tau neutrino to produce detectable “double bang” event

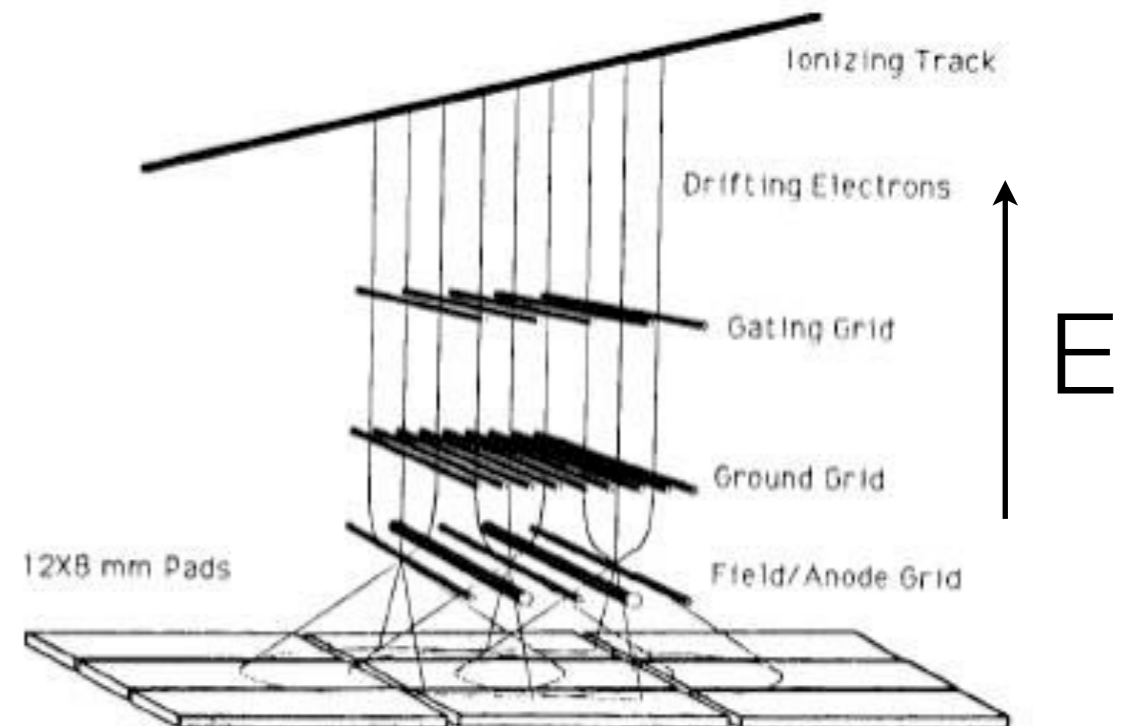
Time Projection Chambers

Time projection chambers (TPCs)

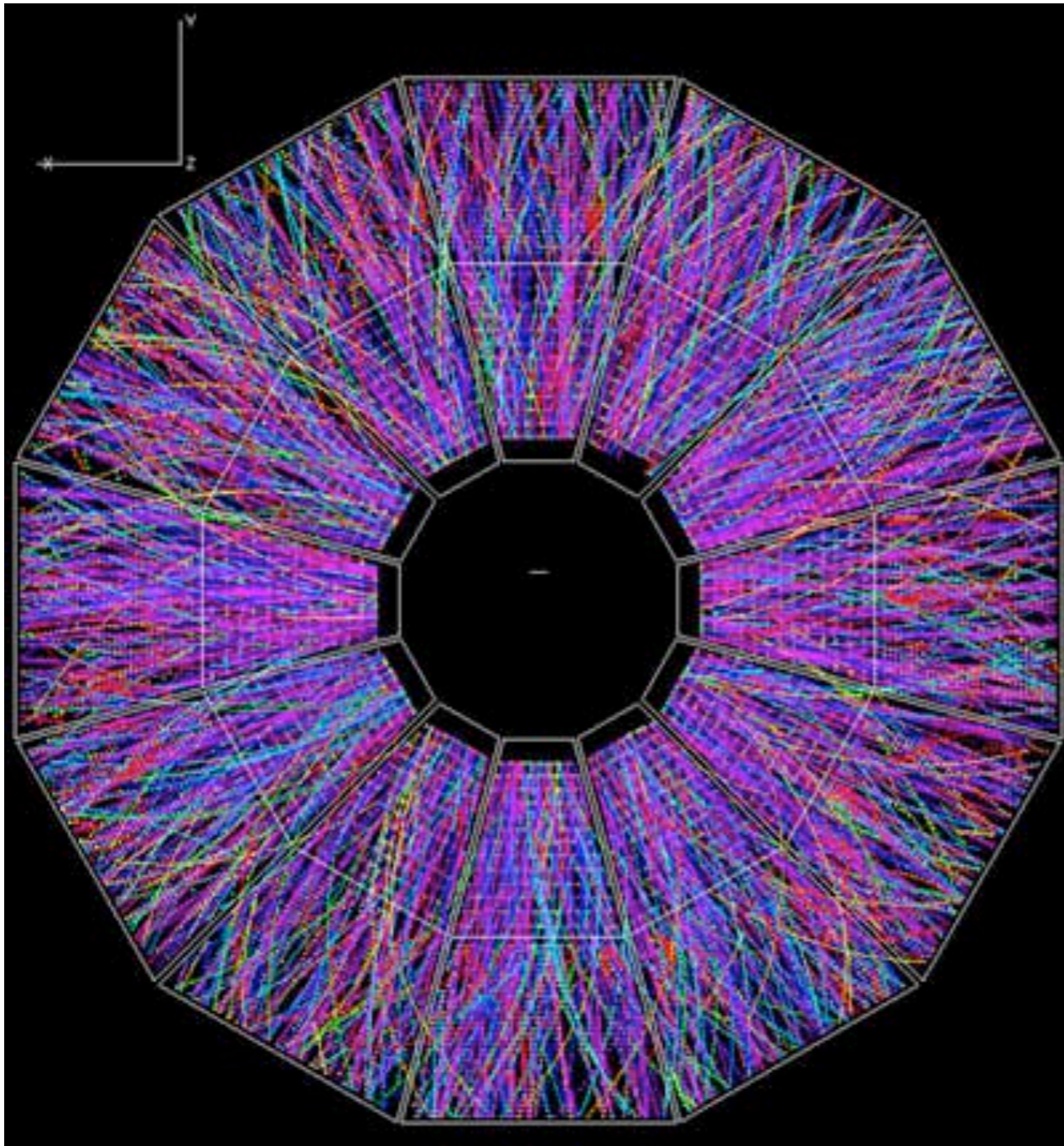
- Gas TPC's have been widely used by a number of high energy and nuclear experiments
- Provide 3D tracking with ~mm resolution.
- Particle ID possible below 1 GeV using $\langle dE/dx \rangle$
- A very common gas Argon-Methane gas mixture. Nobel gas allows for long electron drifts and methane boosts ionization yield.
- Electric fields typically 200 V/cm
- Electron clouds reach terminal velocity at about 5 cm/usec.
- Pads or wire chambers provide 2-D track projection. Arrival times provide 3rd dimension.



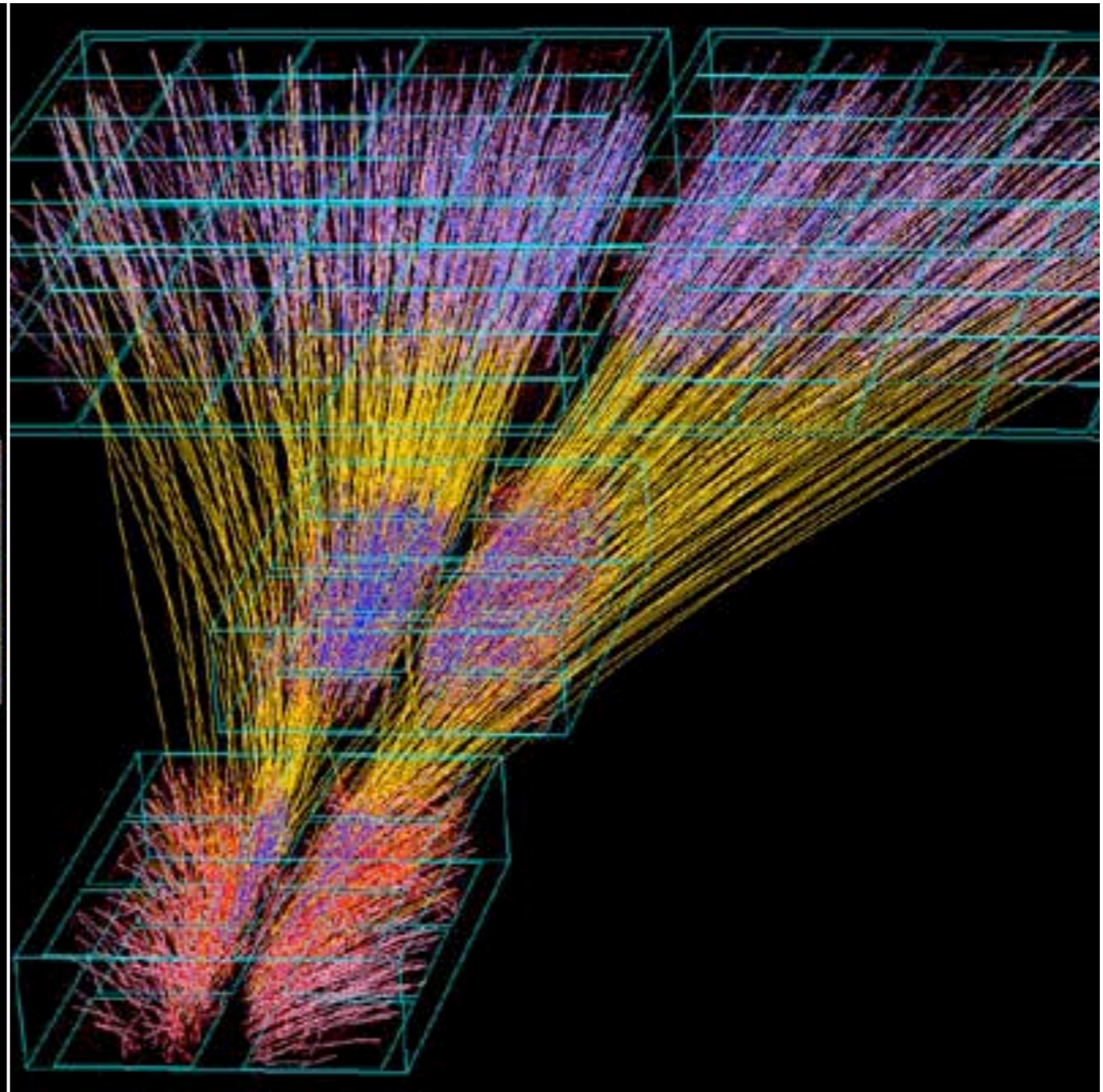
HISS TPC PAD PLANE



TPCs are well suited to high multiplicity environments

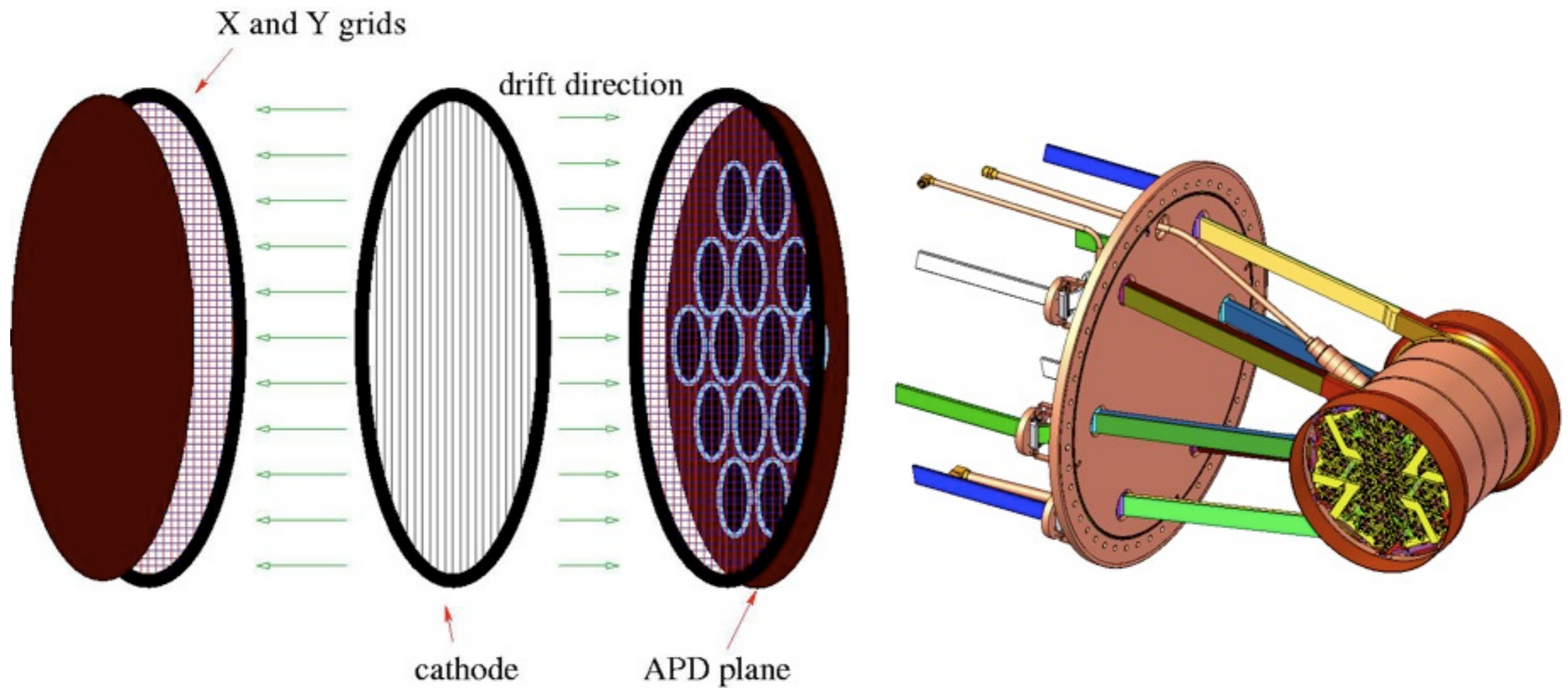


Au + Au at 130 GeV in STAR @ BNL



Pb + Pb at 17 GeV in NA49 @ CERN

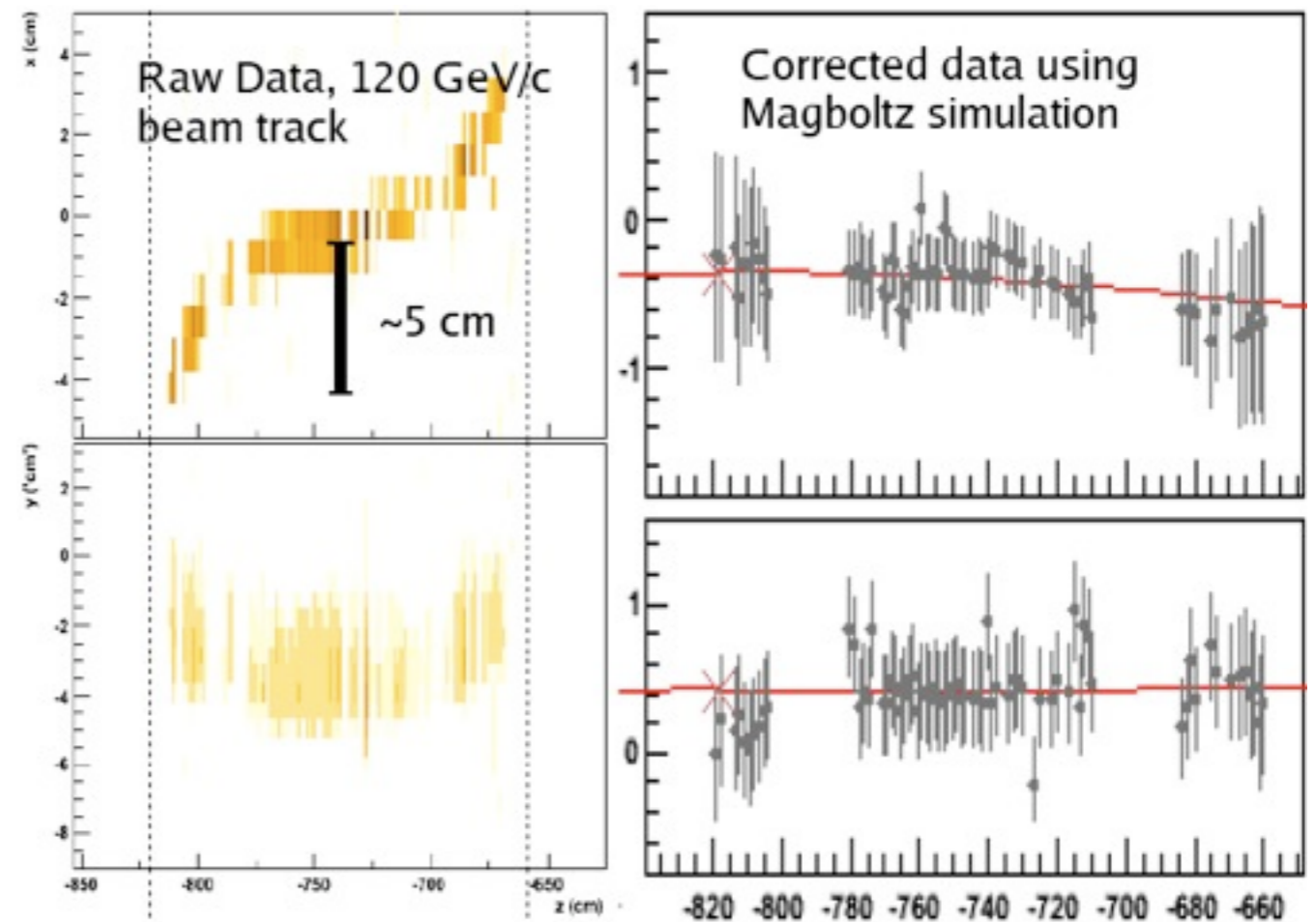
Liquid Xe TPC : EXO



TPC for $0\nu\beta\beta$

Operation of TPC in magnetic field

- Very common to operate TPC's in magnetic fields
- Need to minimize $\vec{E} \times \vec{B}$ or electron drifts become extremely complicated
- In general not a problem if size of TPC \ll aperture of magnet. But, if these are comparable, B field will inevitably have bends.
- If corrections are large, do not trust analytical calculations (eg. Rolandi and Blum) which assume linear "drag" term on electron drift. Rather, learn to use the MAGBOLTZ program.



5 cm distortion over ~1 m drift in a gas Ar TPC (MIPP) due to $E \times B$ effects

Liquid Argon TPC: Concept

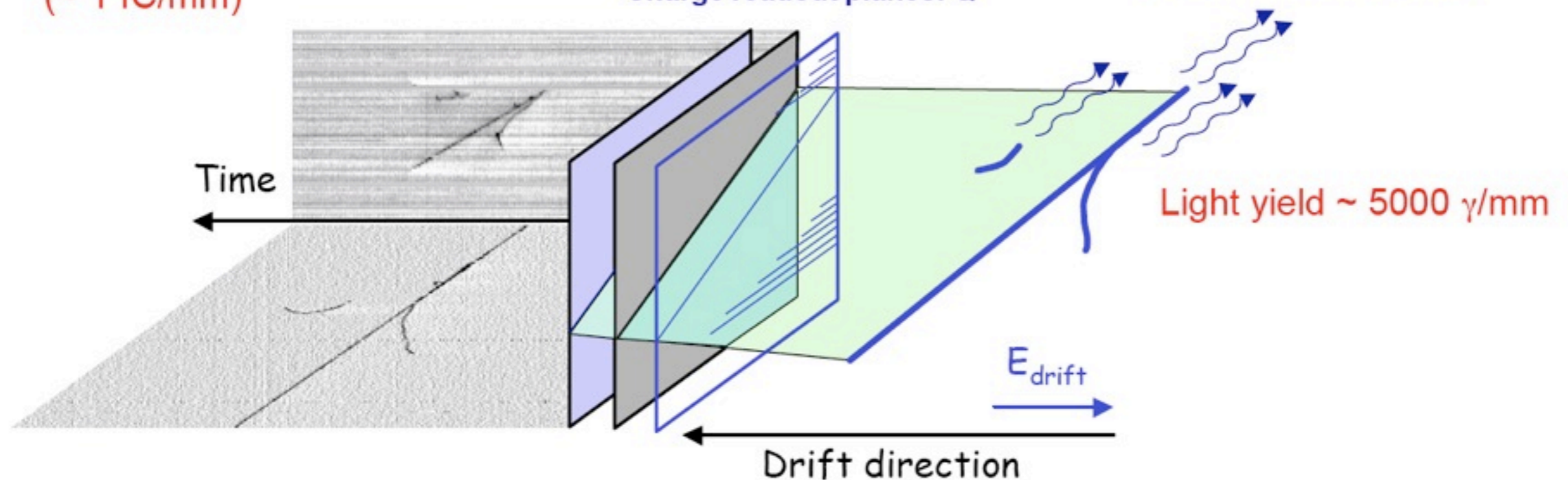
To be applicable to neutrino experiments higher density is required. Use liquid Ar instead of gas. Has potential to reach very large masses (100 kt) with ~mm granularity.

- Boiling point: 87 K (compare to N₂ 77 K)
- Density 1.4 g/cc
- Interaction length: 114 cm
- Radiation length: 14 cm
- Moliere radius: 7 cm

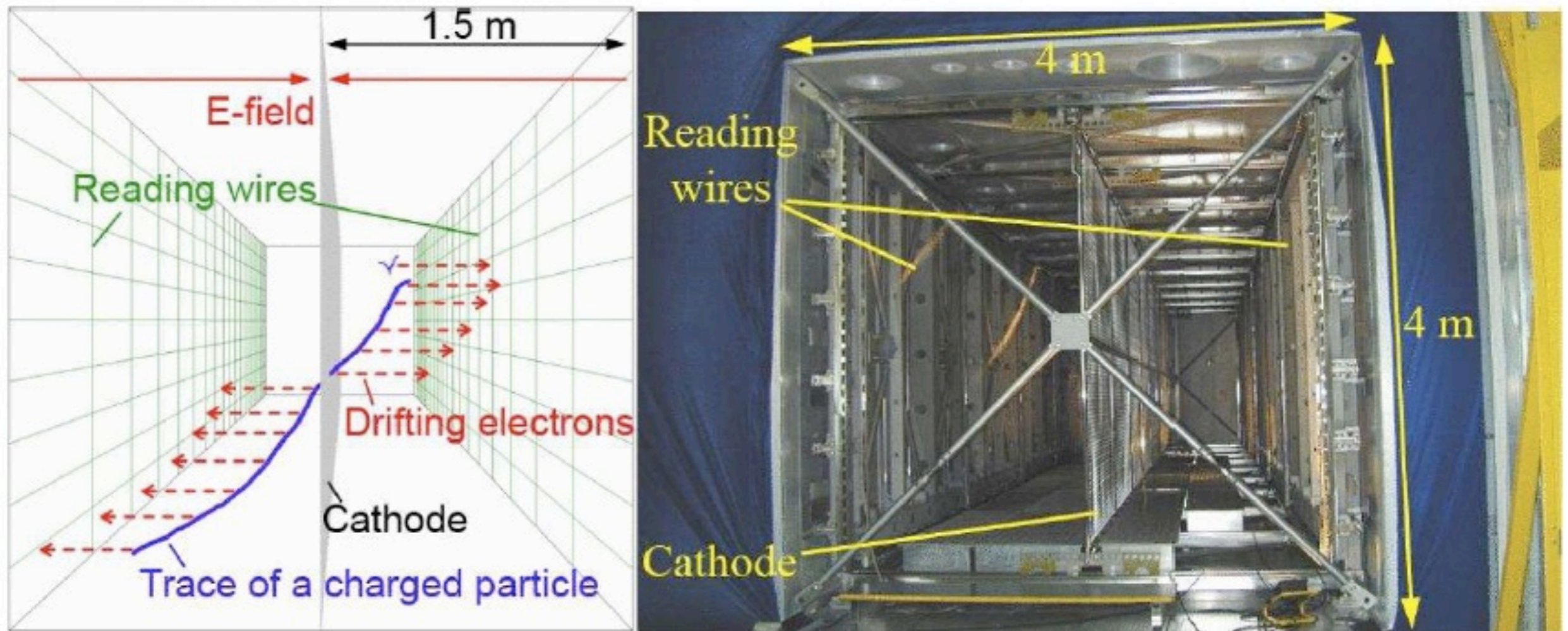
Charge yield ~ 6000 electrons/mm
(~ 1 fC/mm)

Charge readout planes: Q

used to trigger and set to
UV Scintillation Light: L



The ICARUS LqAr Detector



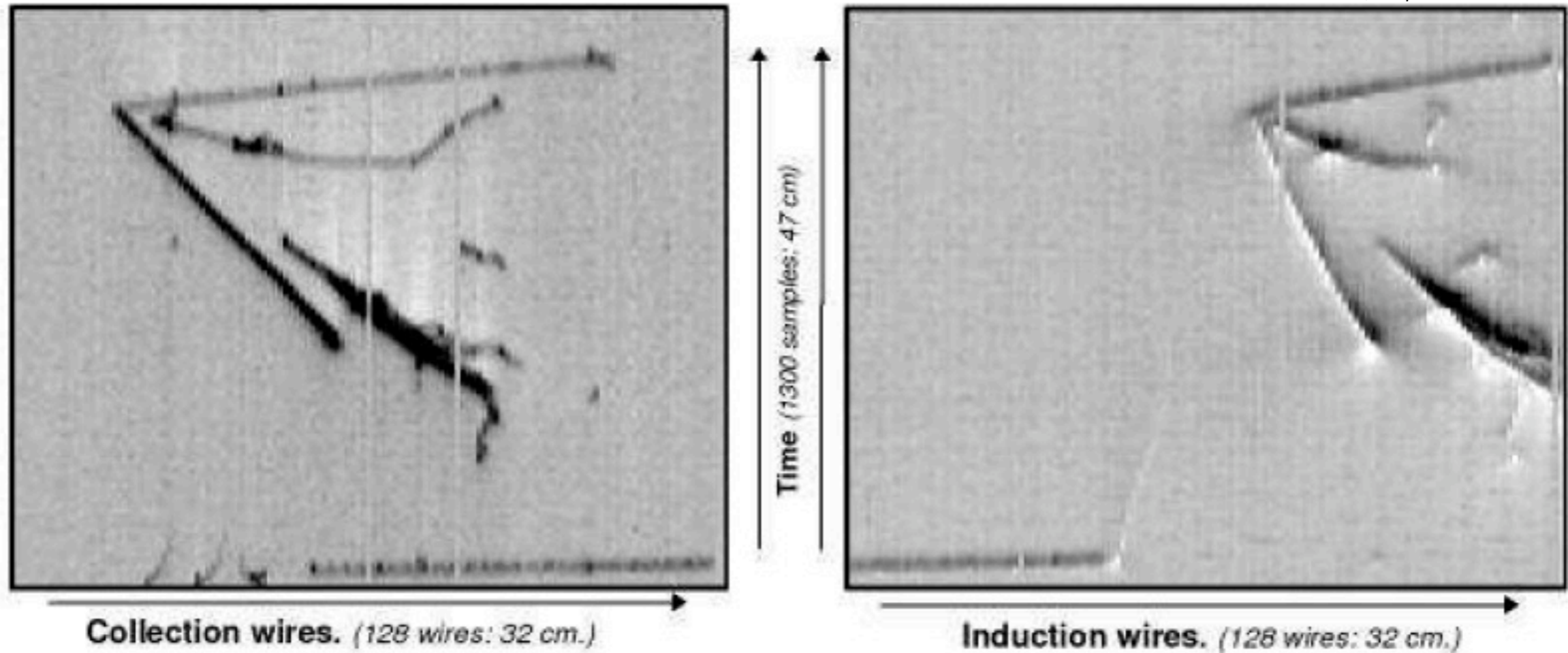
A.M. de la Ossa Romero, hep-ex/0703026

Figure 2.4: Picture of the open T300 ICARUS module during assembly.

What's going on in this event?

Recorded by 50L LqAr detector in WANF beam

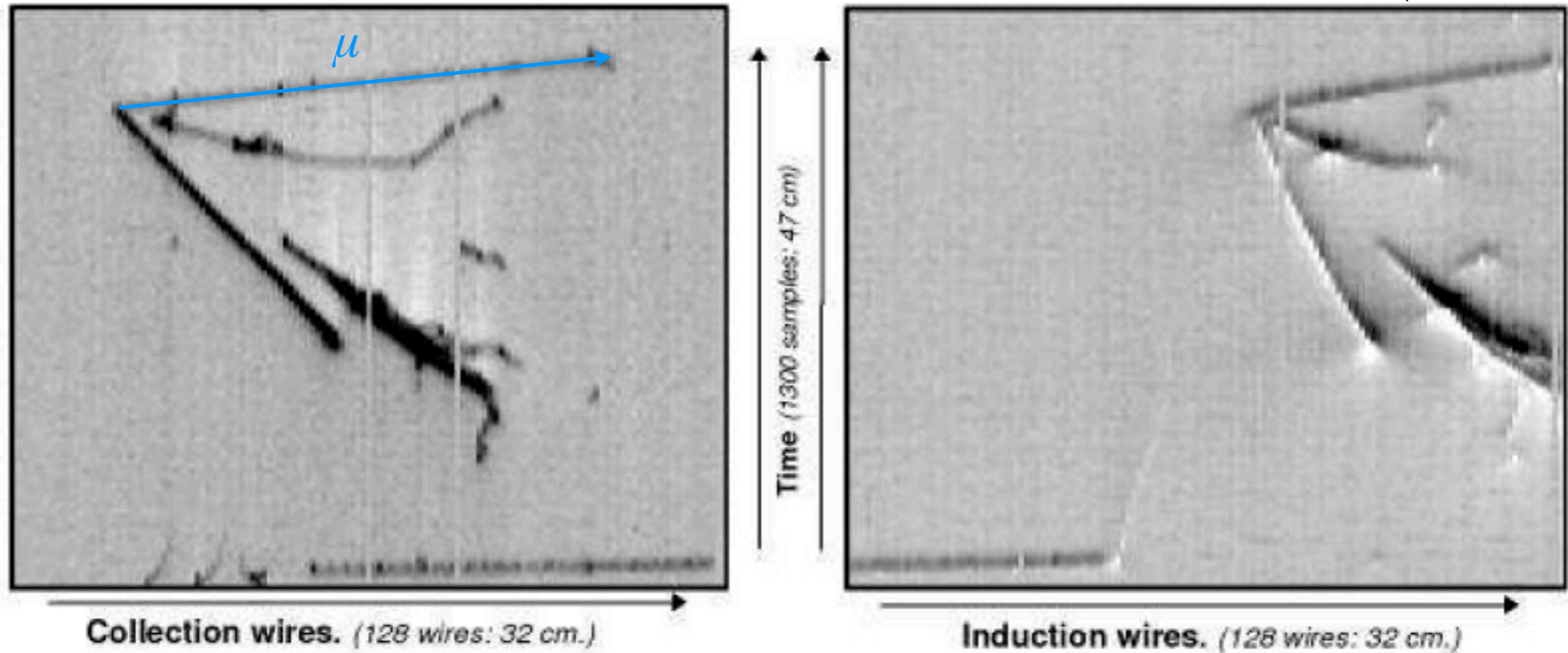
A.M. de la Ossa Romero, hep-ex/0703026



What's going on in this event?

Recorded by 50L LqAr detector in WANF beam

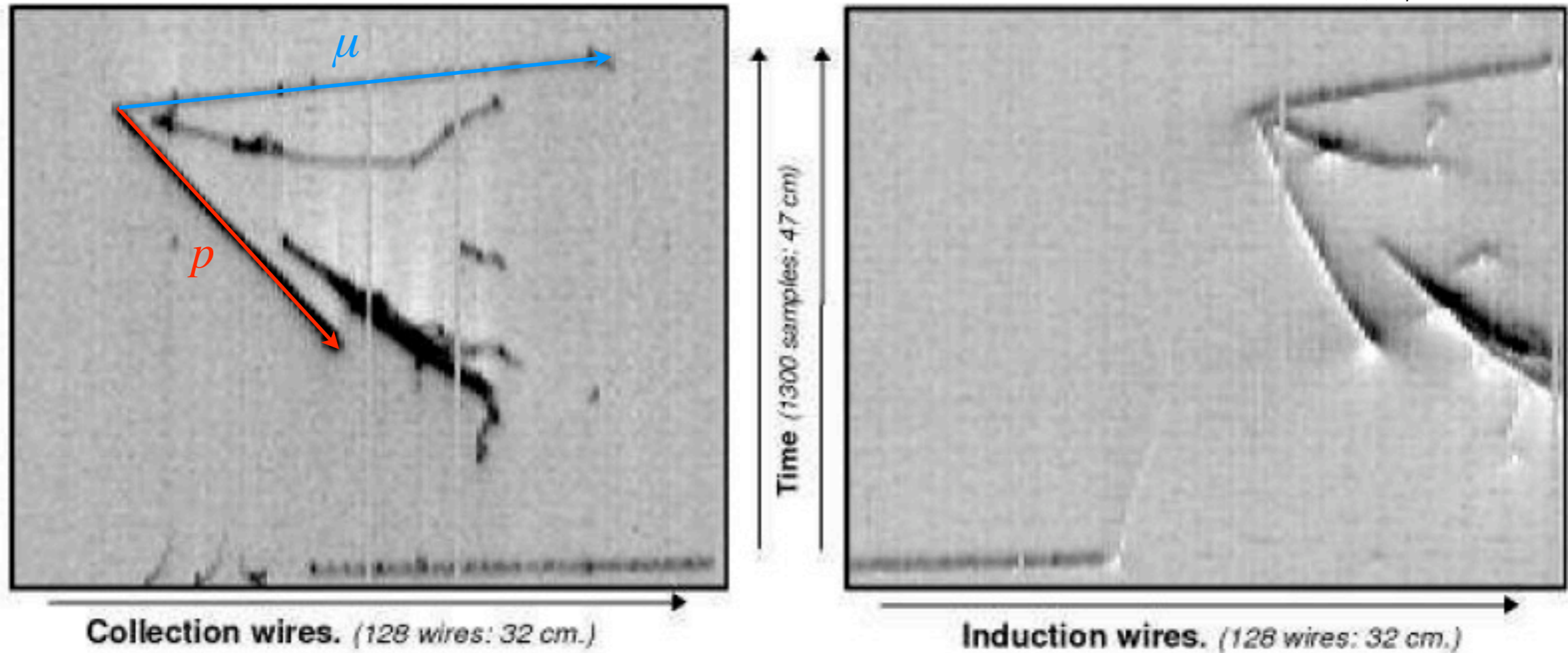
A.M. de la Ossa Romero, hep-ex/0703026



What's going on in this event?

Recorded by 50L LqAr detector in WANF beam

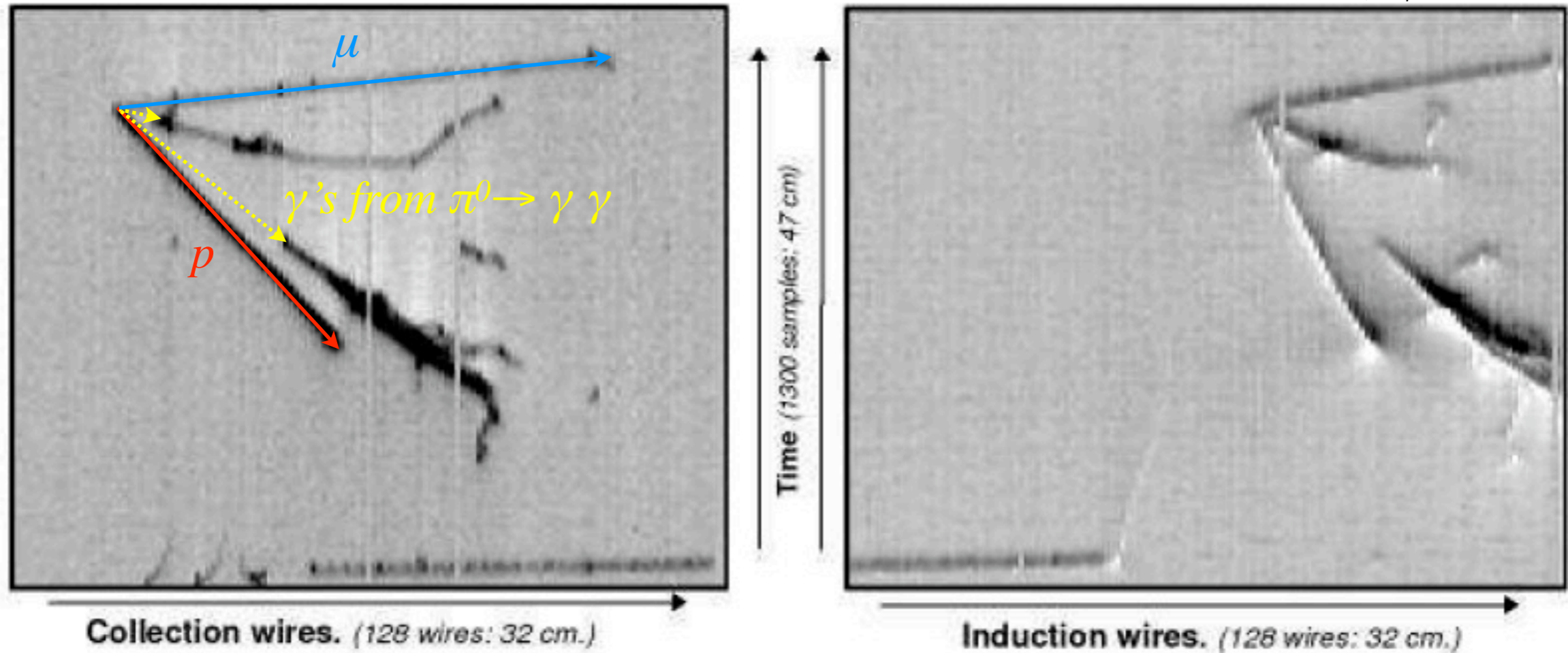
A.M. de la Ossa Romero, hep-ex/0703026



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A.M. de la Ossa Romero, hep-ex/0703026

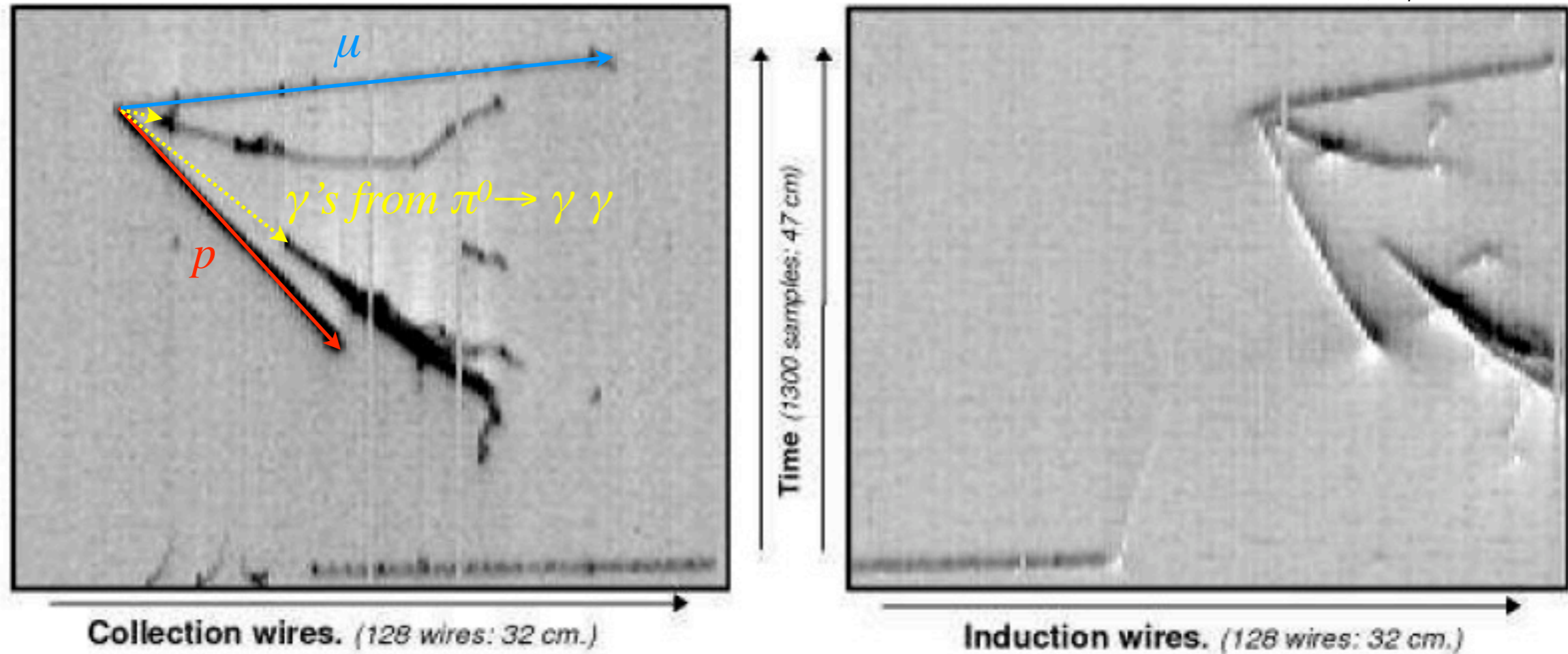


Figure 5.21: The raw image of a low multiplicity real event in the collection (left) and induction plane (right). The event is reconstructed as $(\nu_{\mu} n \rightarrow \mu^{-} \Delta^{+} \rightarrow \mu^{-} p \pi^{0})$ with a mip leaving the chamber, an identified stopping proton and a pair of converted photons from the π^{0} decay. When these photons escape from the chamber, the event is tagged as a *golden* event.

What's going on in this event?

Recorded by 50L LqAr detector in WANF beam

A.M. de la Ossa Romero, hep-ex/0703026

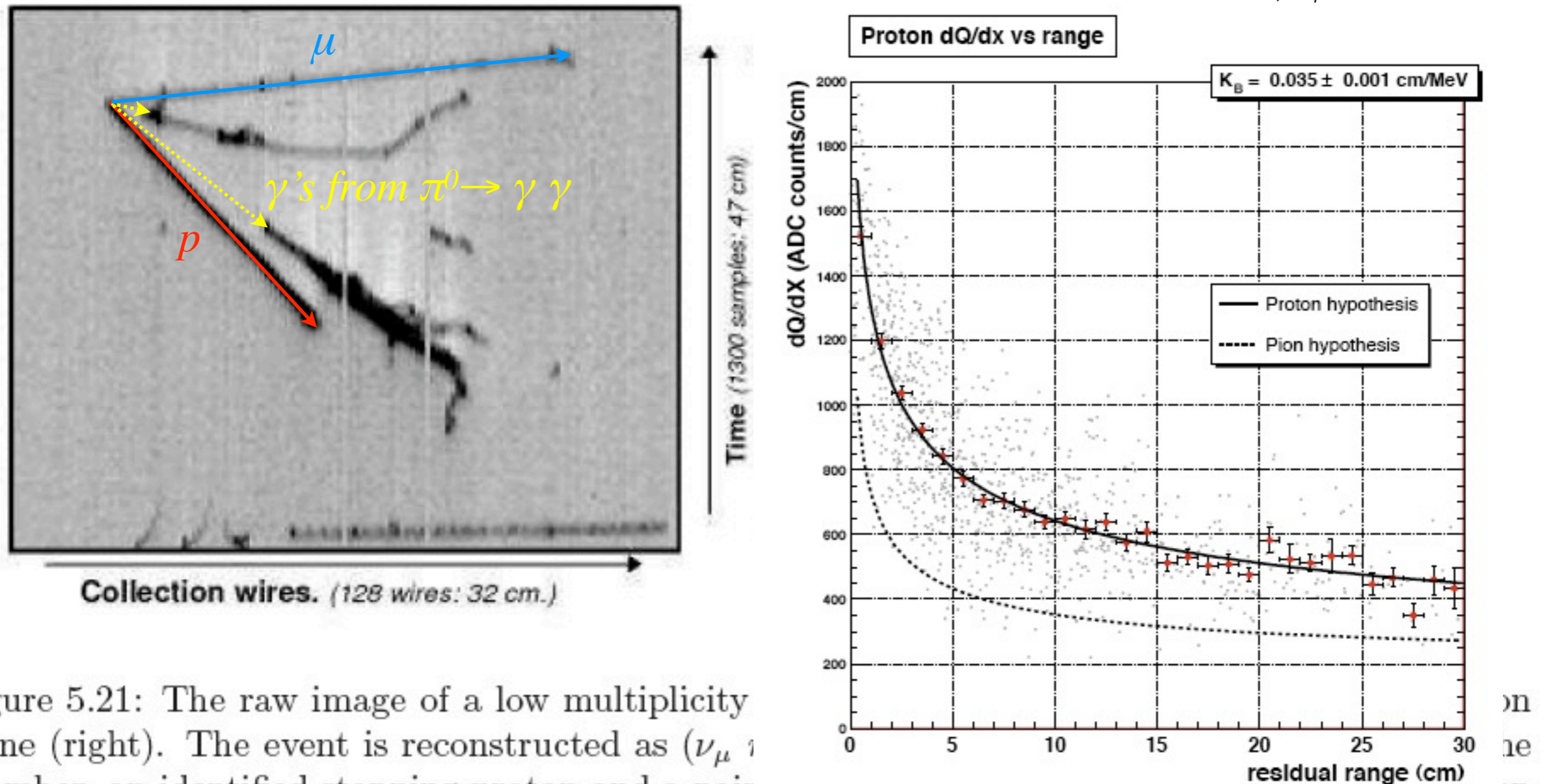
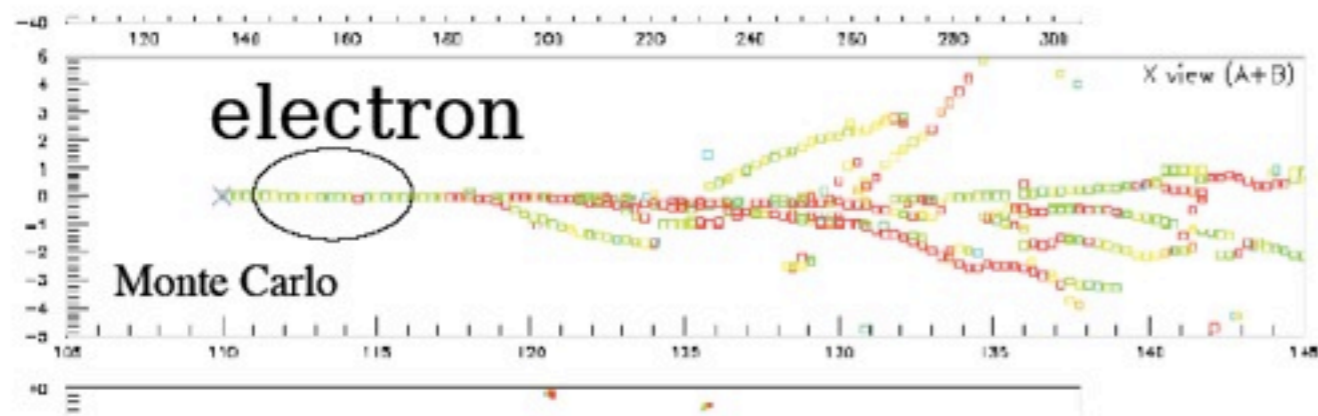


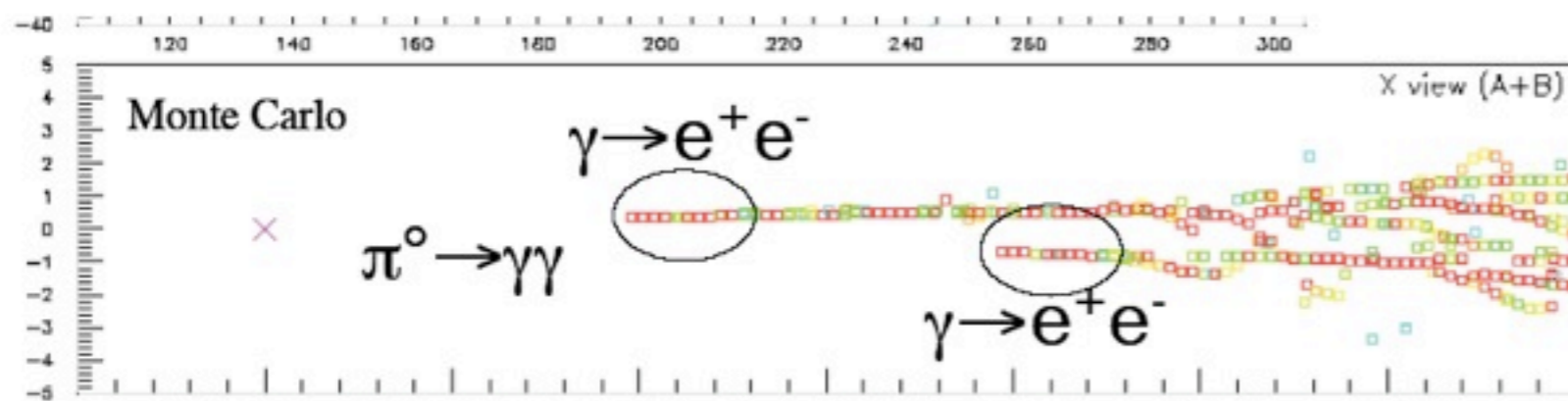
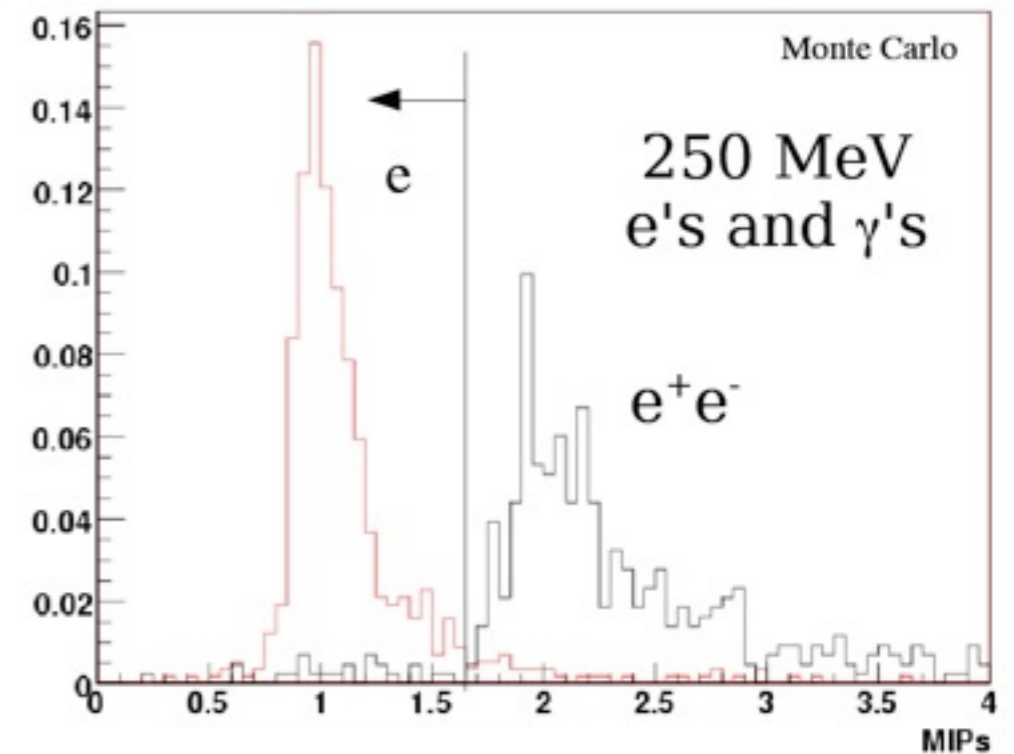
Figure 5.21: The raw image of a low multiplicity plane (right). The event is reconstructed as (ν_μ, μ) in a chamber, an identified stopping proton and a pair of photons. If these photons escape from the chamber, the event is tagged as a *golden event*.

Electron / Photon Separation

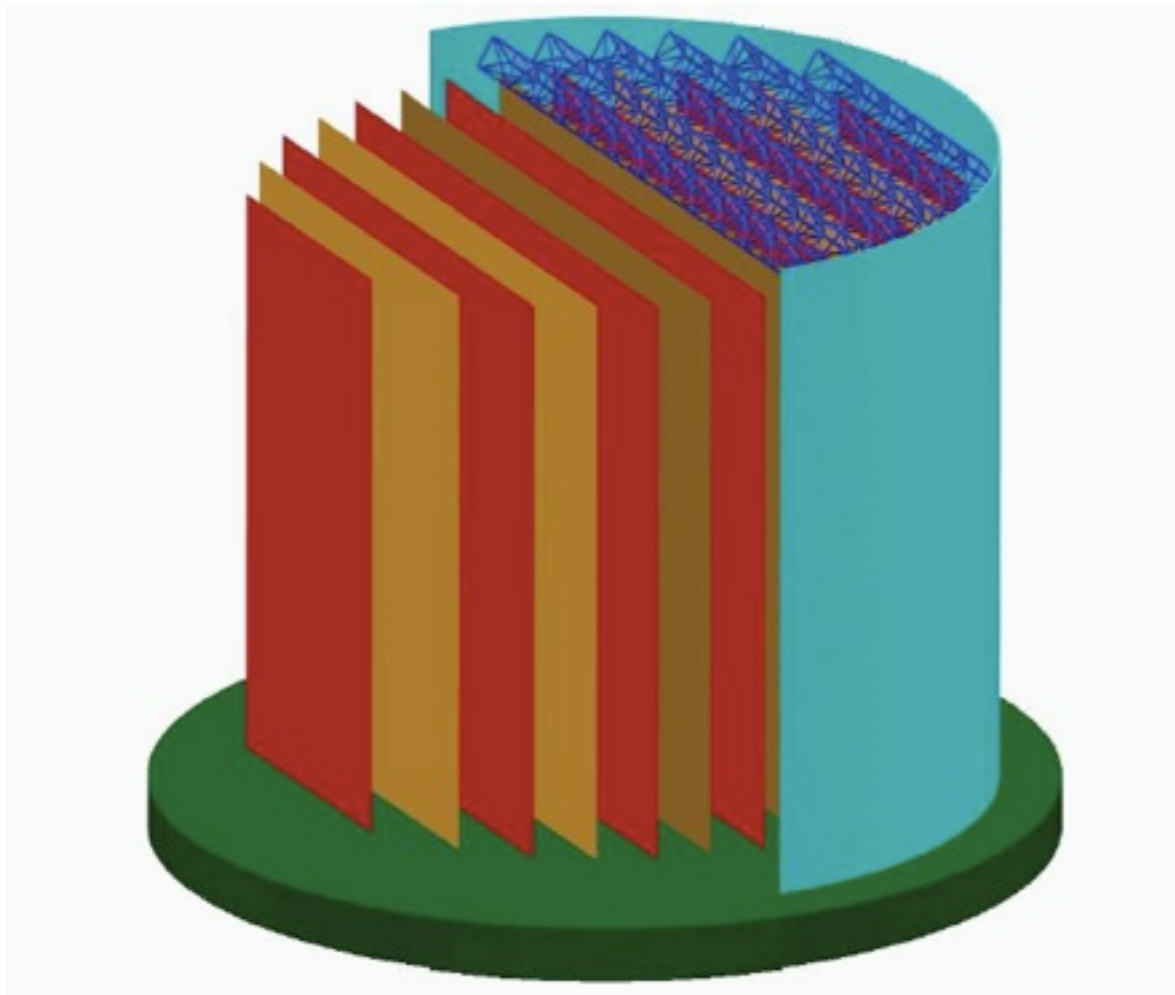
TPC's provide many samples per radiation length. Allows for e/gamma separation by checking dE/dx at start of shower



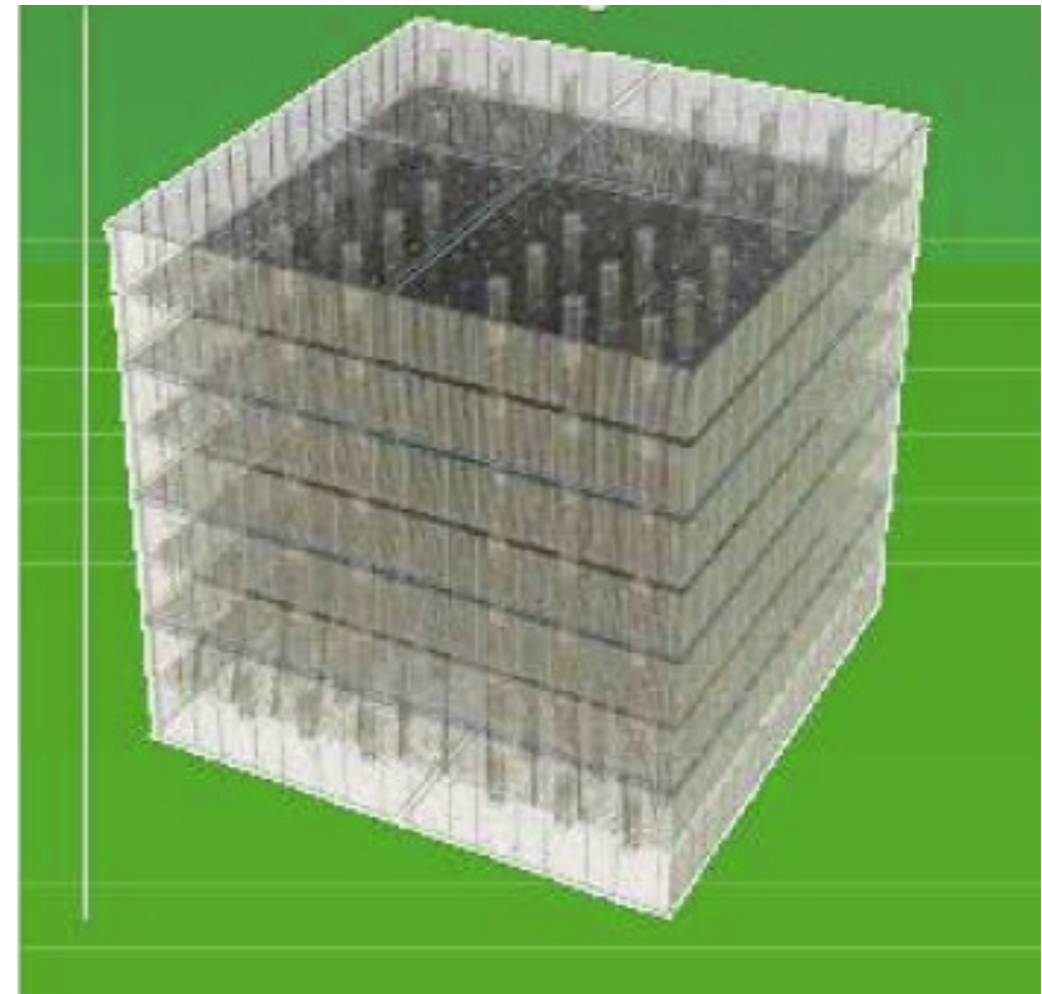
Energy loss in the first 24mm of track: 250 MeV electrons vs. 250 MeV gammas



Some possible designs for big detectors

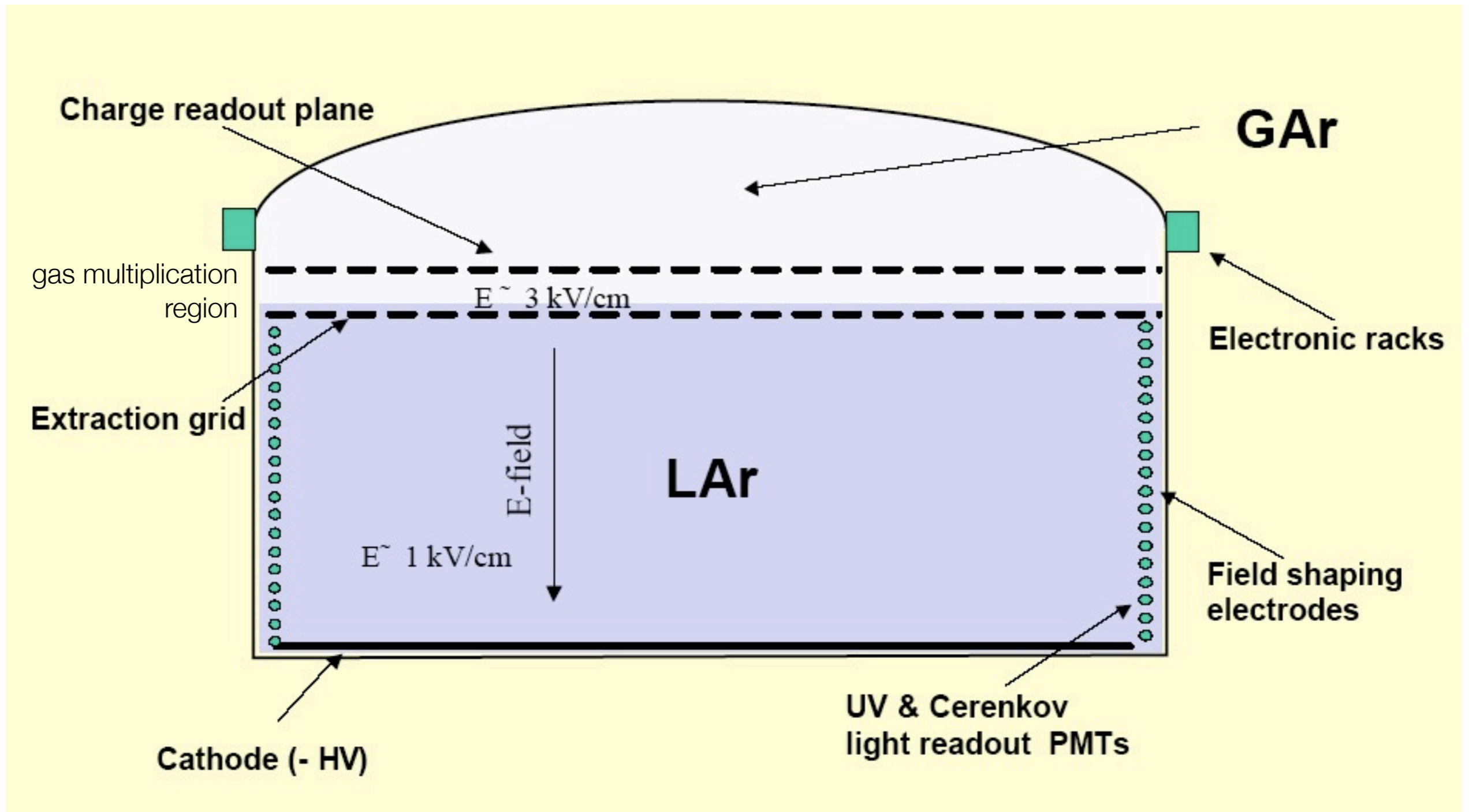


LArTPC: 10-50 kton storage tank. Modular drift regions.



LANDD: Single vessel designed to support vacuum

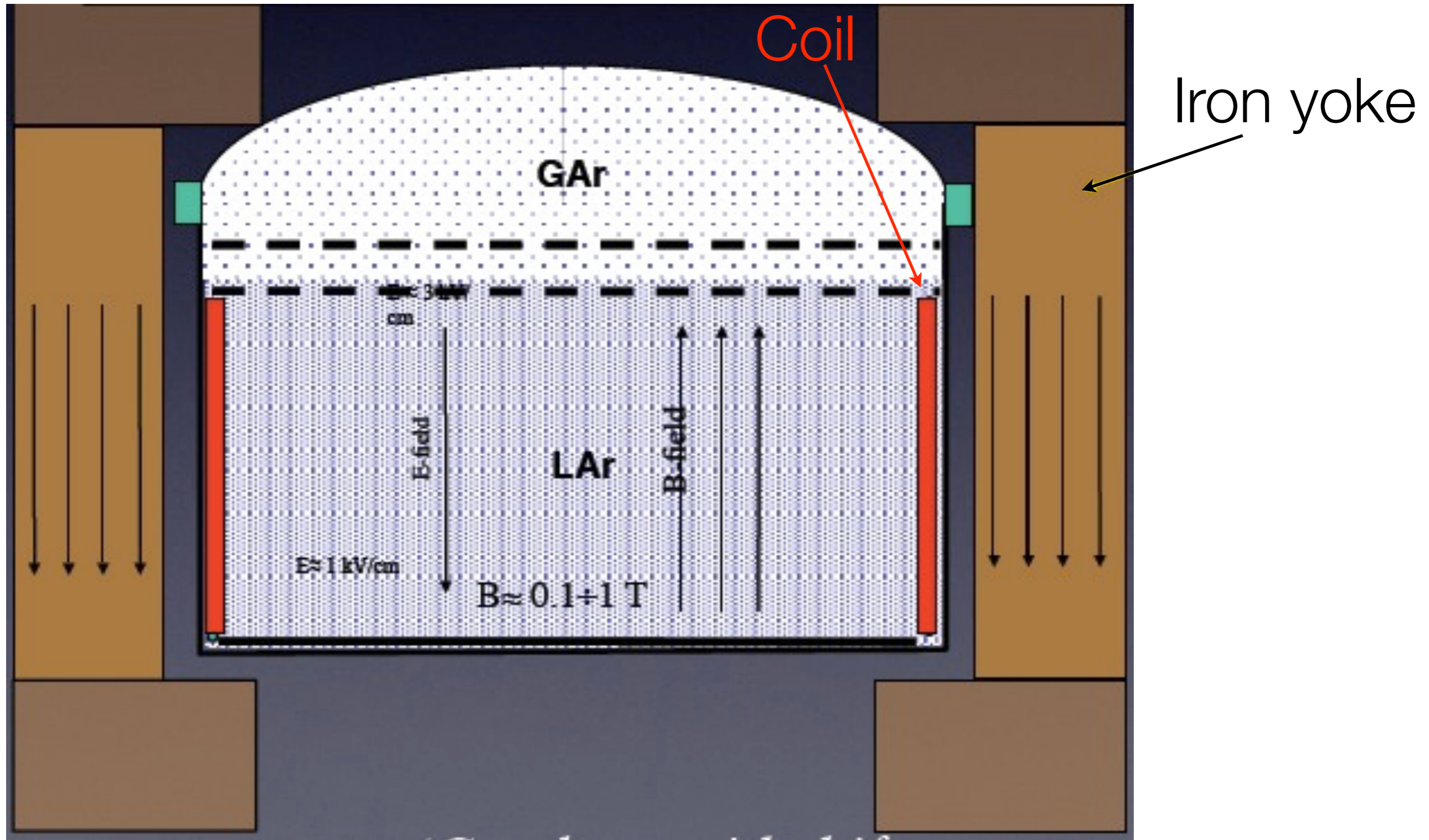
GLACIER Concept



In gas multiplication region, electrons shower in a region of high electric field. Energy/particle goes up as a result of acceleration in the field.

Concepts for large, magnetized, LqAr detectors

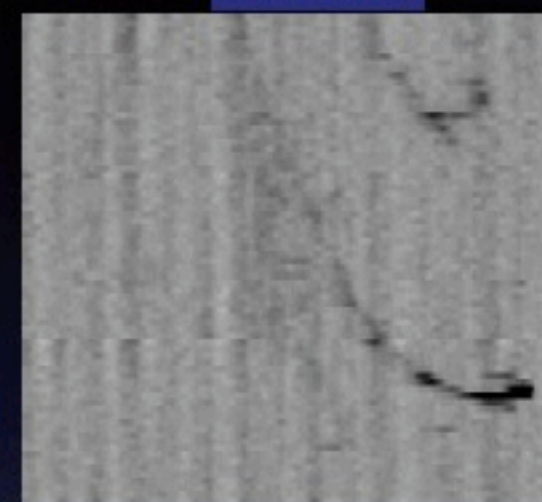
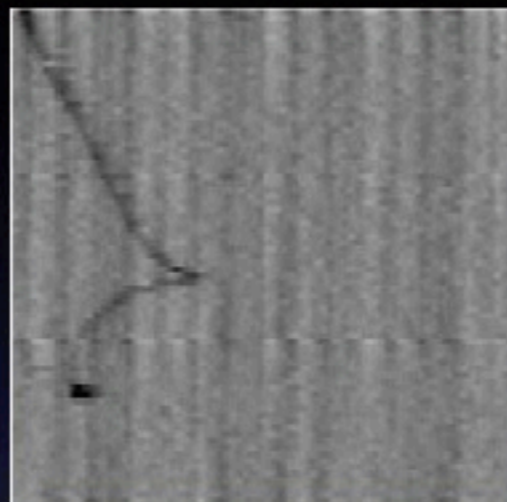
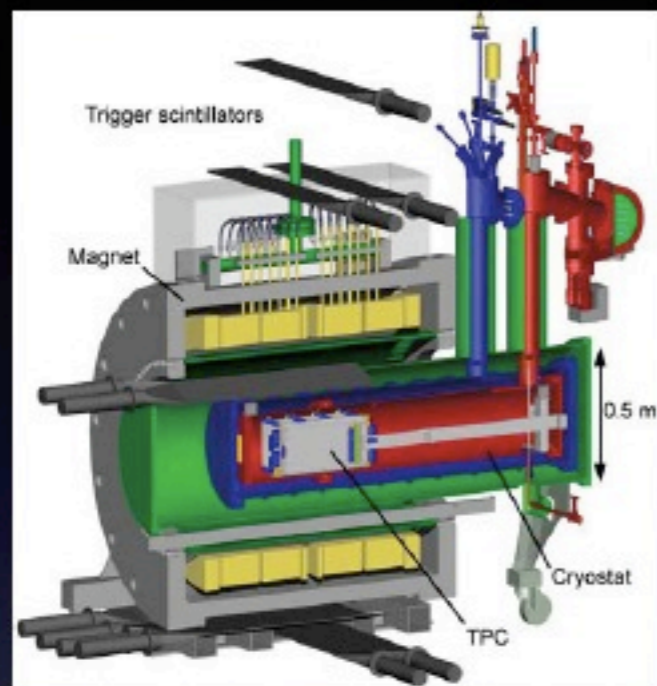
Andre Rubbia, Golden'07



First operation of a LAr TPC embedded in a B-field

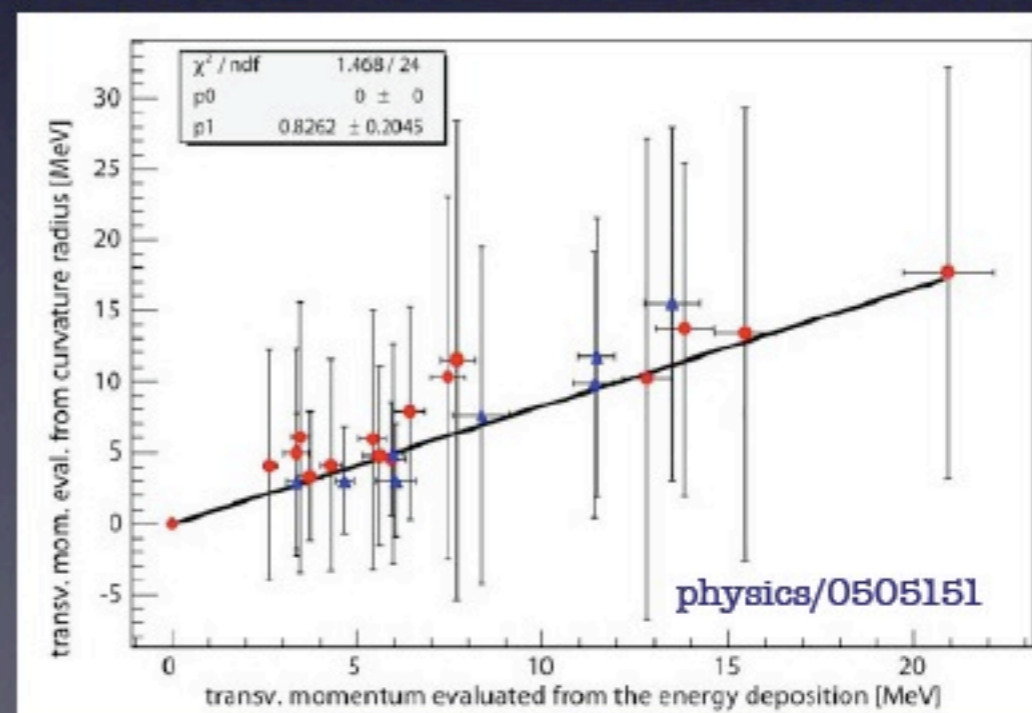
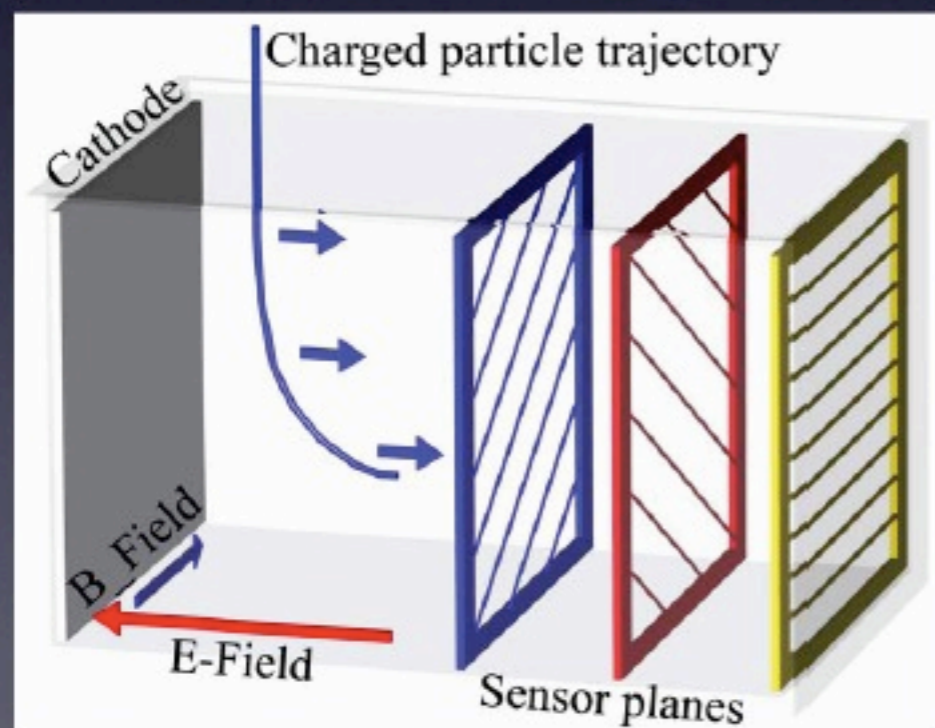
First real events in B-field ($B=0.55T$):

New J. Phys. 7 (2005) 63
NIM A 555 (2005) 294



150 mm
150 mm

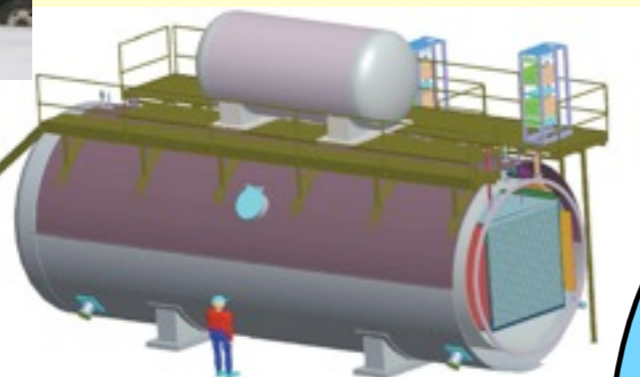
Correlation between calorimetry and magnetic measurement for contained tracks:



Path to large detectors (U.S.)



“ArgoNeut”
running in front of
MINOS now



“MicroBooNE” in 8
GeV beamline at
Fermilab

0.3 ton

x600

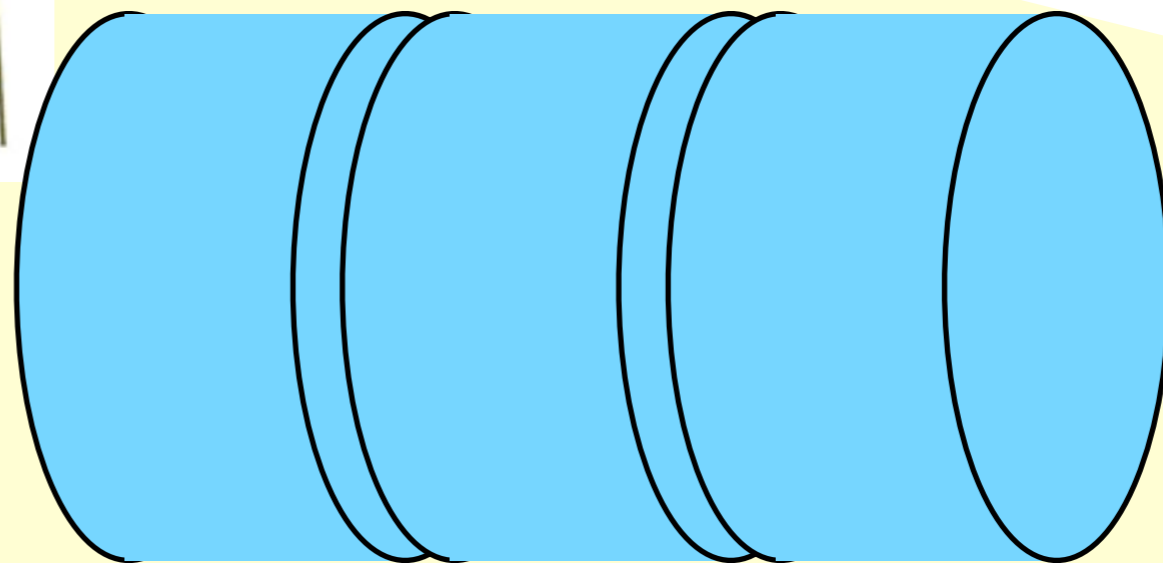
170 ton

x30

5 kton

x20

100 kton



LAr5 in future DUSEL
beamline

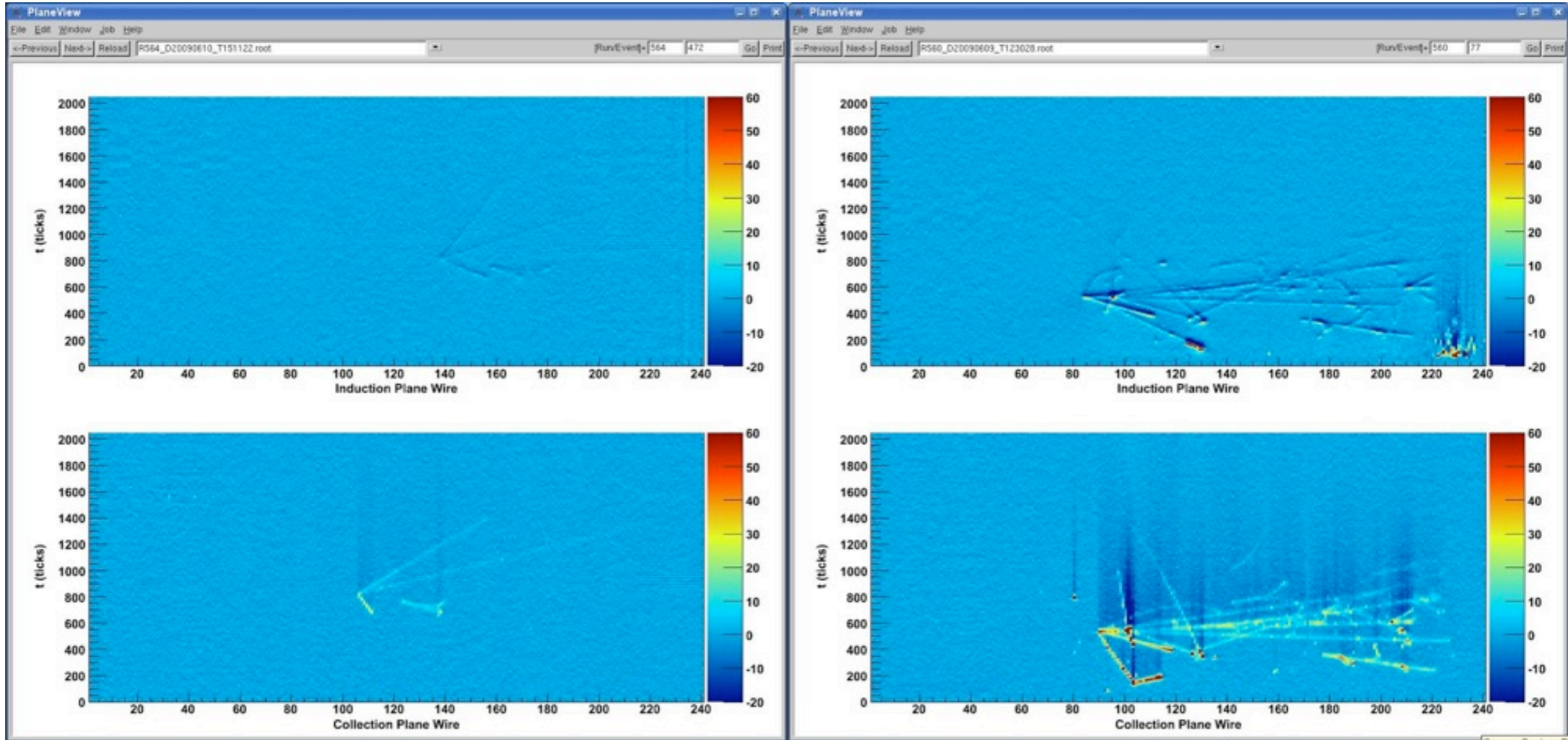


Eventual goal:
100 kt

R&D issues

- Are the drift distances required by large detectors achievable in large cryostats?
- Electronic optimization. Multiplexing? Noise?
- Large wire plane construction

First events from ArgoNEUT



courtesy of B. Fleming

Summary

Neutrino event topologies

- Muons : Long straight, ~constant energy deposit of $2 \text{ MeV cm}^2 / \text{g}$
- Electrons : Create compact showers. Longitudinal size determined by radiation length. Transverse size determined by Moliere radius.
- Photons: Create compact showers after a gap of ~ 1 radiation length.
- Hadrons : Create diffuse showers. Scale determined by interaction length

Specific technologies:

- *Cherenkov*: Best for low rate, low multiplicity, energies below 1 GeV
- *Tracking calorimeters*: Can handle high rate and multiplicities. Best at 1 GeV and above.
- *Unsegmented scintillator calorimeters*: Large light yields at MeV energies. Background considerations dominate design.
- *TPCs*: Great potential for large mass with high granularity. Lots of activity to realize potential