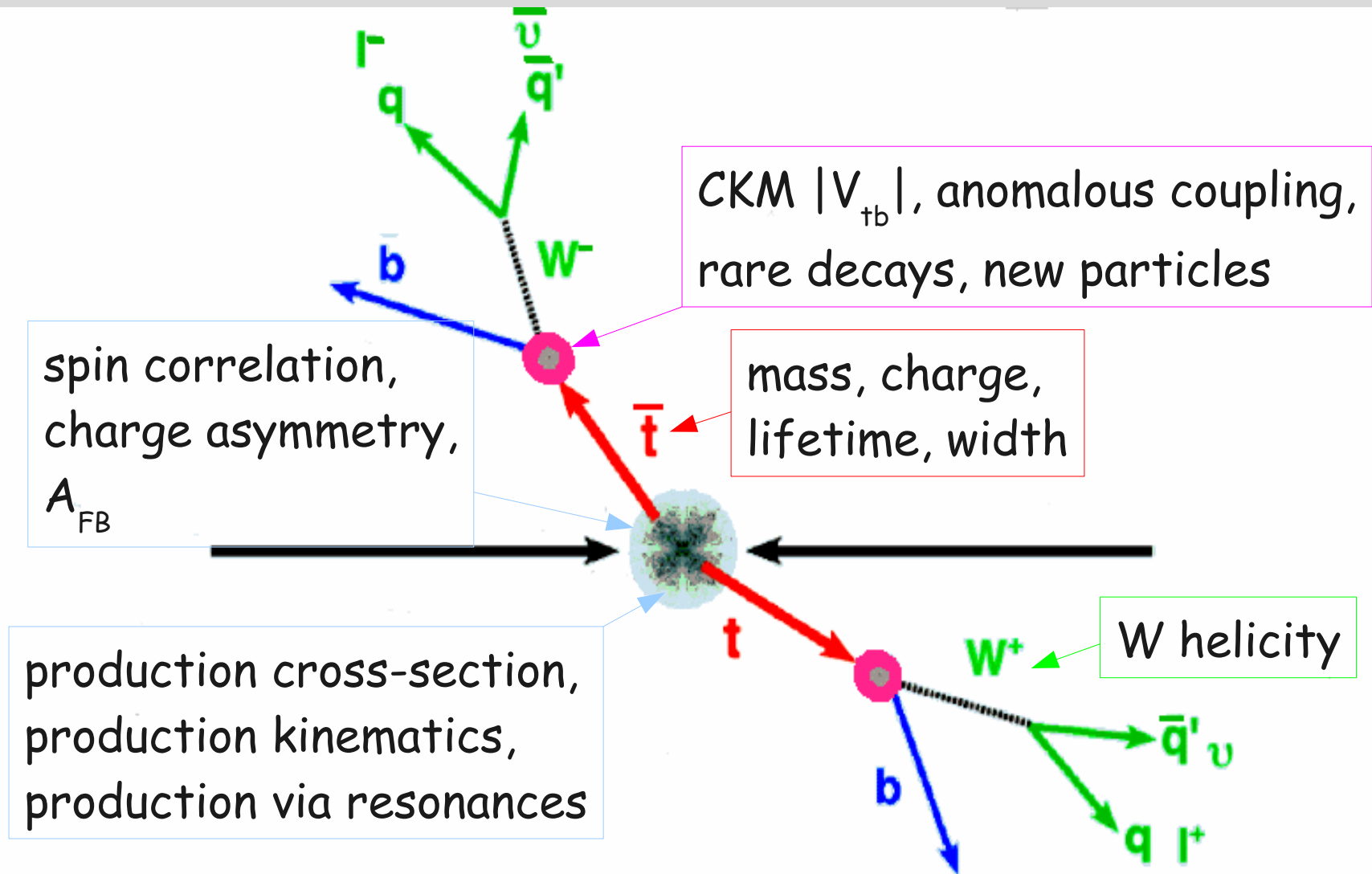


Top Quark Properties



Top Quark Mass Measurement

- Possible to measure the quark mass
- Important ingredient for SM precision tests: $B \rightarrow X_s \gamma$ and $K_L \rightarrow \pi^0 \nu \nu$
- Help predict (verify) the Higgs sector due to the relationship with W and H
- Measure the mass from the reconstructed decay products has low precision due to the presence of jets and neutrino. Use other methods.

Template method

- Choose an observable, x , sensitive to m_T
- x can be: lepton P_t , reconstructed top mass, decay length
- Predict the x distribution as a function of m_T using Monte Carlo
- For each event evaluate the likelihood for each m_T value
- Maximize the likelihood for the entire sample

Matrix Element

- Use all information from the event integration over the least known variables

Top Quark Mass Measurement: Template Method

Method: build top mass and JES template for signal and background
Use the templates as pdf in the Likelihood. Extract top mass and JES

Hadronic decay channel

Reconstruct the event kinematic by minimizing:

$$\chi^2 = \frac{(m_{jj}^{(1)} - M_W)^2}{\Gamma_W^2} + \frac{(m_{jj}^{(2)} - M_W)^2}{\Gamma_W^2} + \frac{(m_{jjb}^{(1)} - m_t^{rec})^2}{\Gamma_t^2} + \frac{(m_{jjb}^{(2)} - m_t^{rec})^2}{\Gamma_t^2} + \sum_{t=1}^n \frac{(P_{T,t}^{fit} - P_{T,t}^{meas})^2}{\sigma_t^2}$$

m_{jj} = invariant mass of two light jets

m_{jjb} = invariant mass of three jets

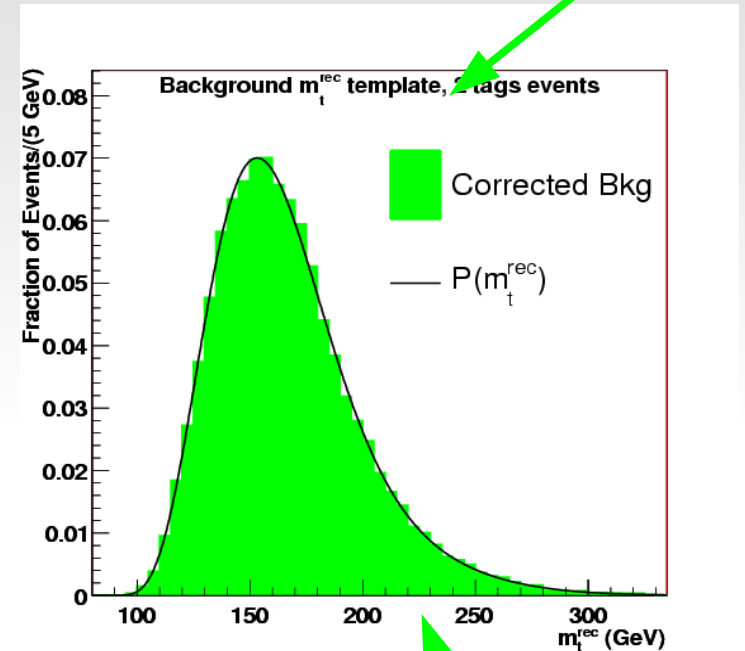
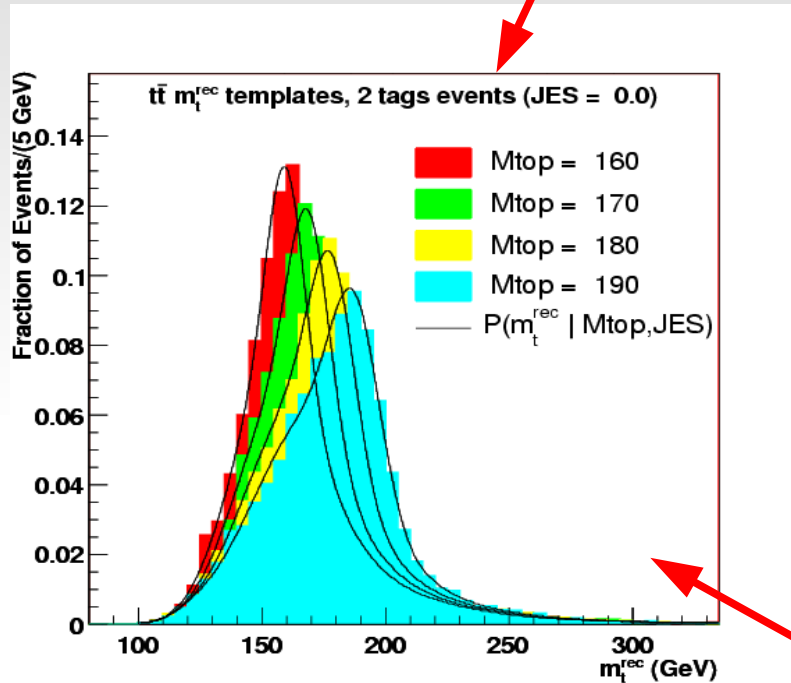
P_T^{fit} = top transv. momentum

For each permutation we obtain m_t^{rec} this forms the template for signal (MC) and background (data)

Top Quark Mass Measurement: Template Method

Signal template: Monte Carlo data

Background template: data



$$L_{tot} = L_{M_{top}} + L_{JES}$$

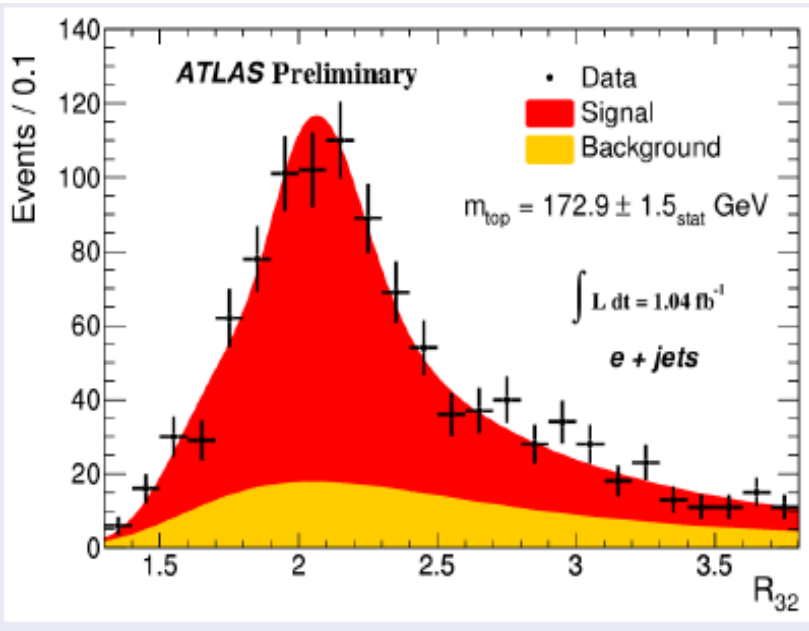
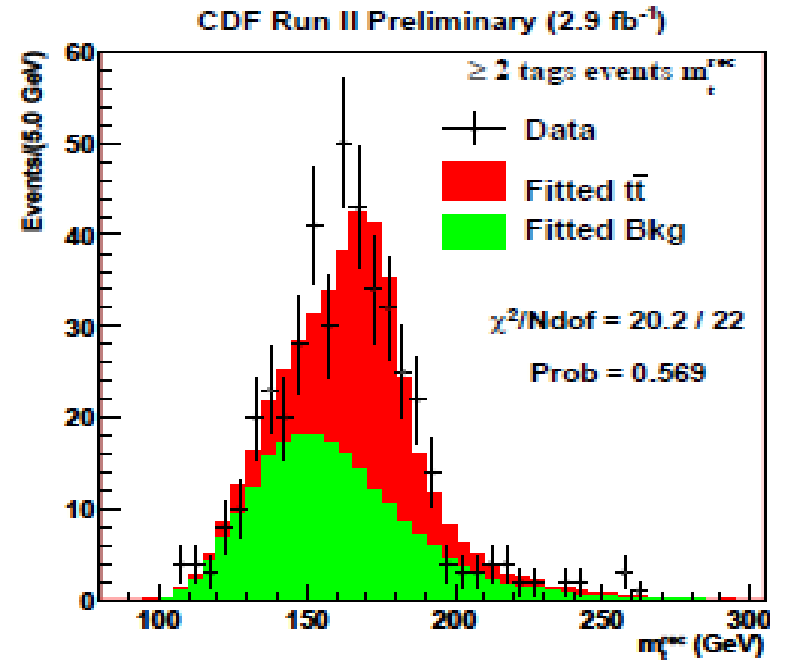
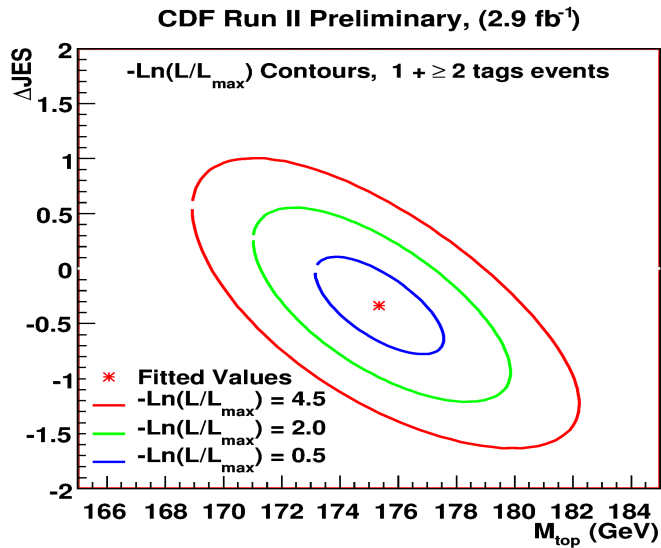
The 2 free parameters are: top mass and JES

$$\mathcal{L}_{M_{top}} = \prod_{i=1}^{N_{obs}} \frac{\nu_s \cdot P_{sig}^{m_t^{rec}}(m_{t,i} | M_{top}, JES) + \nu_b \cdot P_{bkg}^{m_t^{rec}}(m_{t,i})}{\nu_s + \nu_b}$$

$$\mathcal{L}_{JES} = \prod_{i=1}^{N_{obs}} \frac{\nu_s \cdot P_{sig}^{m_W^{rec}}(m_{W,i} | M_{top}, JES) + \nu_b \cdot P_{bkg}^{m_W^{rec}}(m_{W,i})}{\nu_s + \nu_b}$$

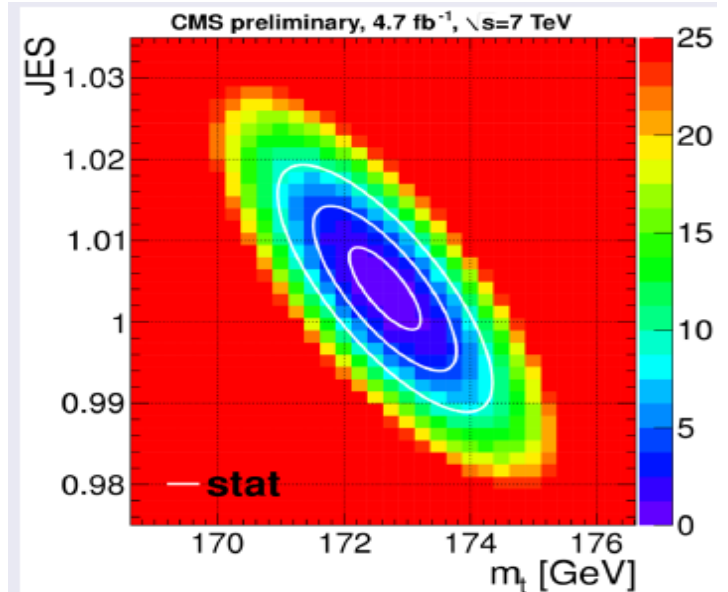
Top Quark Mass Measurement: Template Method

Tevatron
all hadronic
decay
channel



LHC
μ+jets
decay
channel

$$R_{32} = \frac{m_t^{reco}}{m_W^{reco}}$$



Top Quark Mass Measurement: Matrix Element

Observables: measured momenta of jets and leptons

Question: for an observed set of kinematic variables x what is the most probable top mass

Method: start with an observed set of events of given kinematics and find maximum of the likelihood, which provides the best measurement of top quark mass

Our sample is a mixture of signal and background

$$P_{evt}(x, m_t) = f_{top} \cdot P_{sgn}(x, m_t) + (1 - f_{top}) \cdot P_{bkg}(x)$$

$P_{bkg}(x)$ depends on the decay channel

$$P_{sgn}(x; m_t) = \frac{1}{\sigma(m_t)} \int d^n \sigma(y; m_t) dq_1 dq_2 f(q_1) f(q_2) W(x, y)$$

Top Quark Mass Measurement: Matrix Element

probability to observe a set of kinematic variables x for a given top mass

$$P_{\text{sgn}}(x; m_t) = \frac{1}{\sigma(m_t)} \int d^n \sigma(y; m_t) dq_1 dq_2 f(q_1) f(q_2) W(x, y)$$

Normalization depends on m_t
Includes acceptance effects

$d^n \sigma$ is the differential cross section
Contains **matrix element** squared

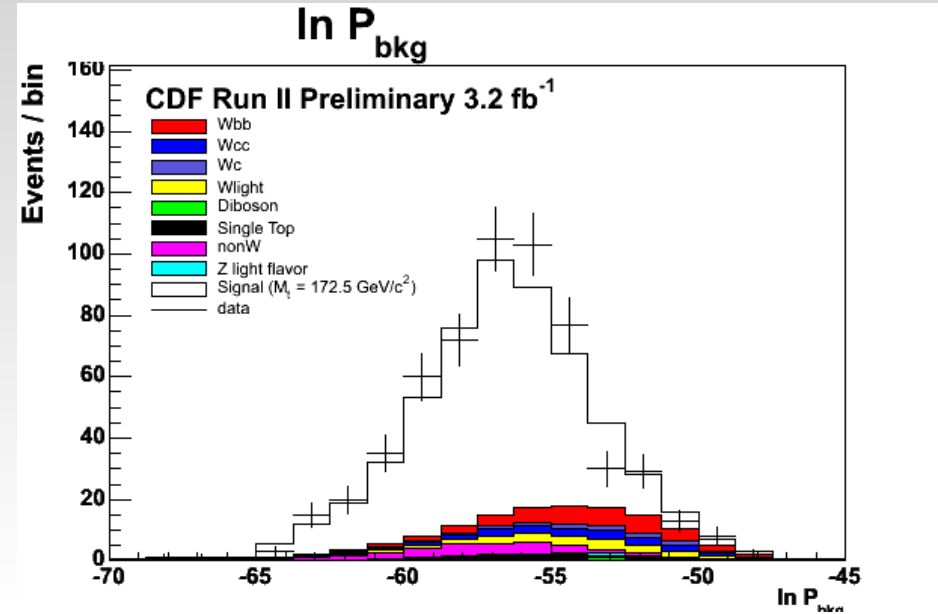
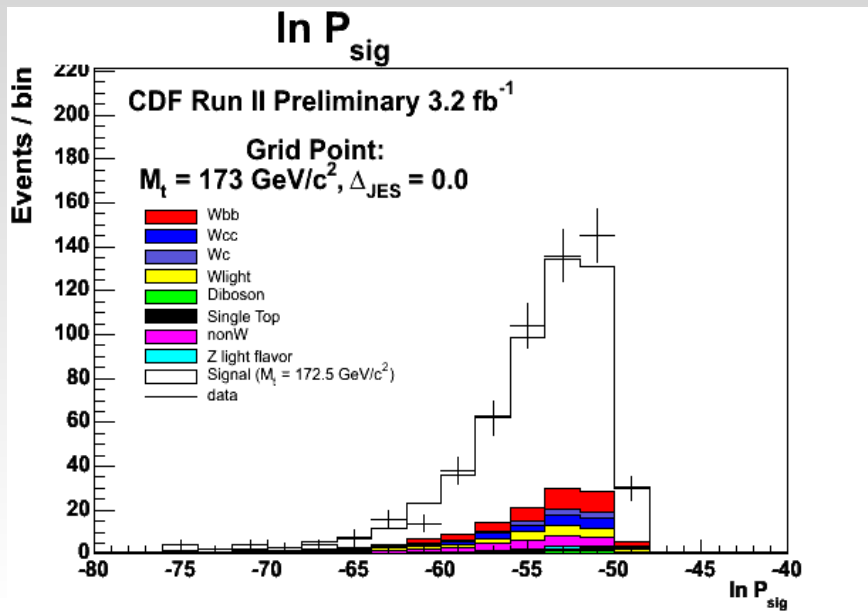
$W(x, y)$ is the probability that a parton level set of variables y will be measured as a set of variables x .
Parton Energy \leftrightarrow Jet Energy

$f(q)$ is the probability distribution that a parton will have a momentum q

Integrate over unknown:
kinematical variable q_1, q_2 of initial states parton and final states parton y
Approximations: LO matrix element and $qq \rightarrow t\bar{t}$ process only (no gluon fusion - 15%)

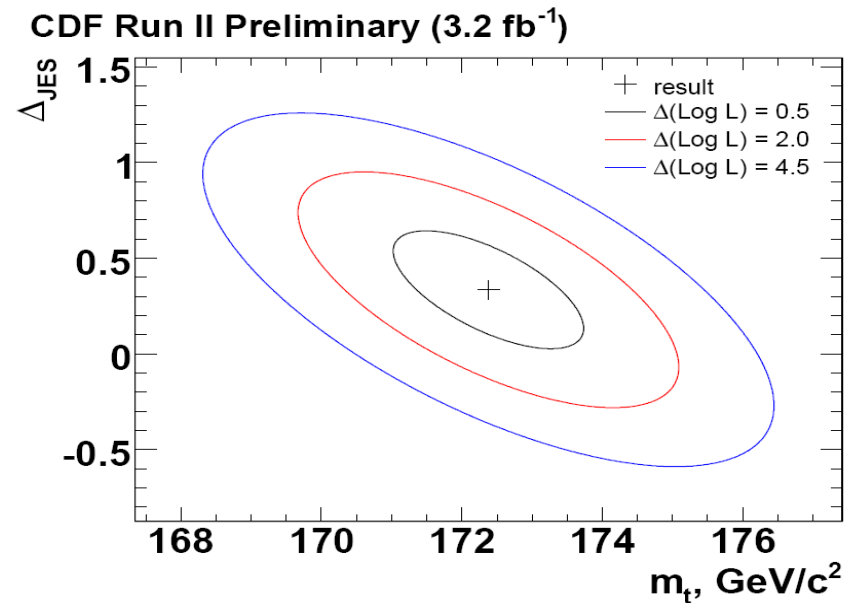
In a similar way is constructed $P_{\text{bkg}}(x)$

Top Quark Mass Measurement with Matrix Element

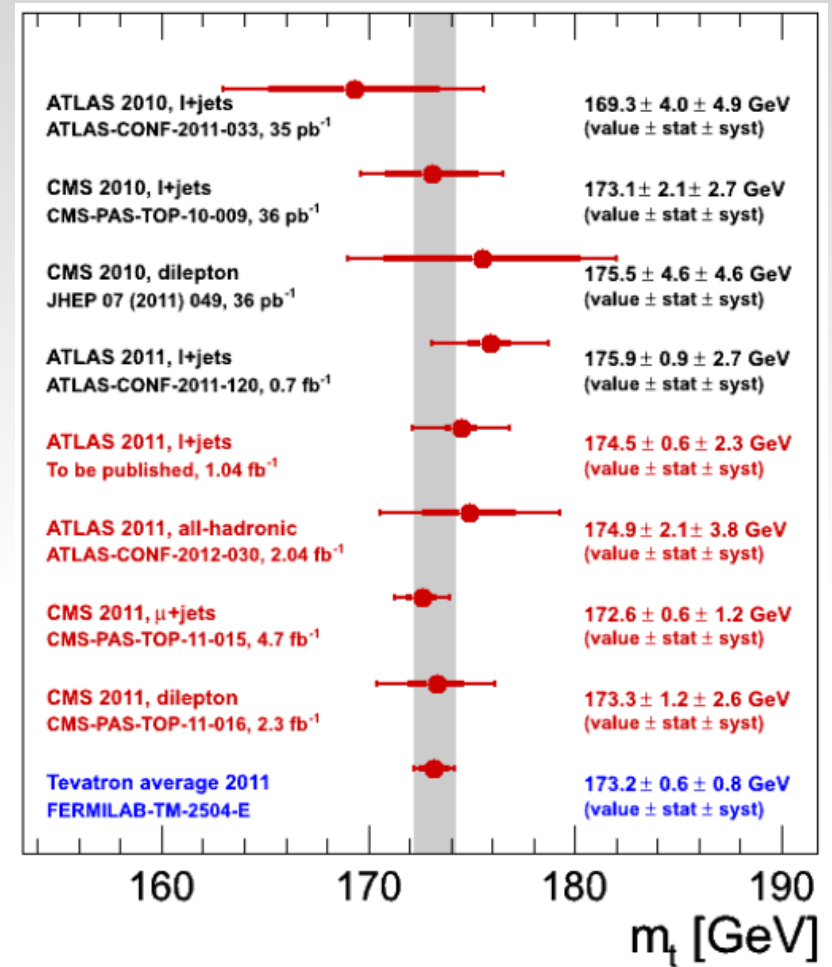
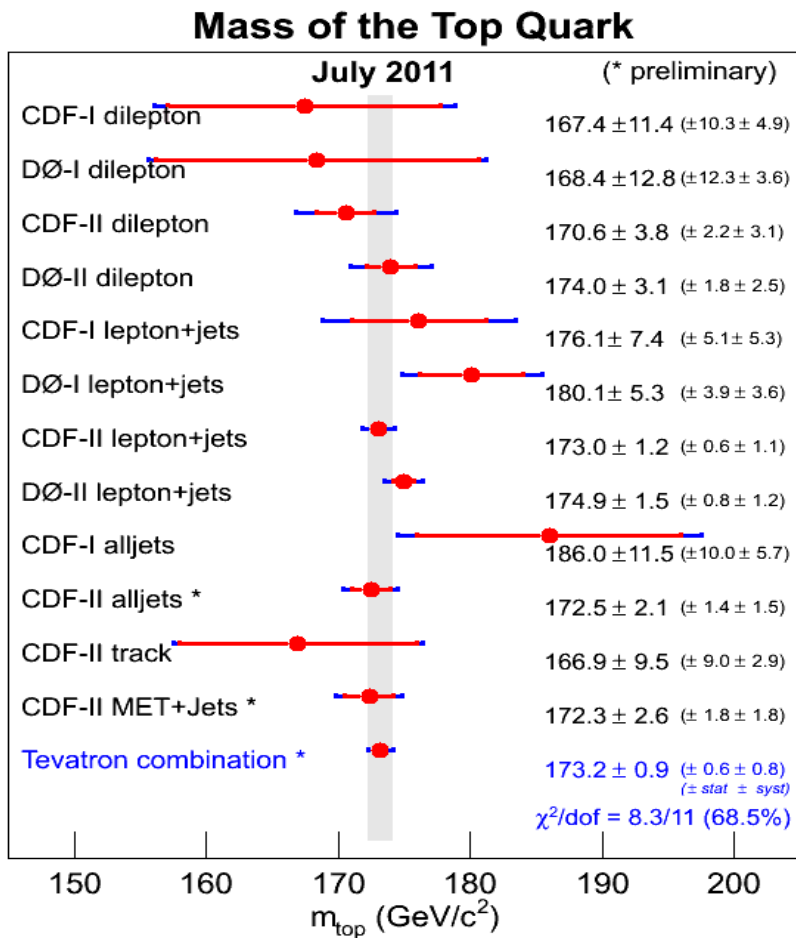


Then the likelihood which uses
 is minimized to obtain

$$P_{evt}(x, m_t) = f_{top} \cdot P_{sgn}(x, m_t) + (1 - f_{top}) \cdot P_{bkg}(x)$$

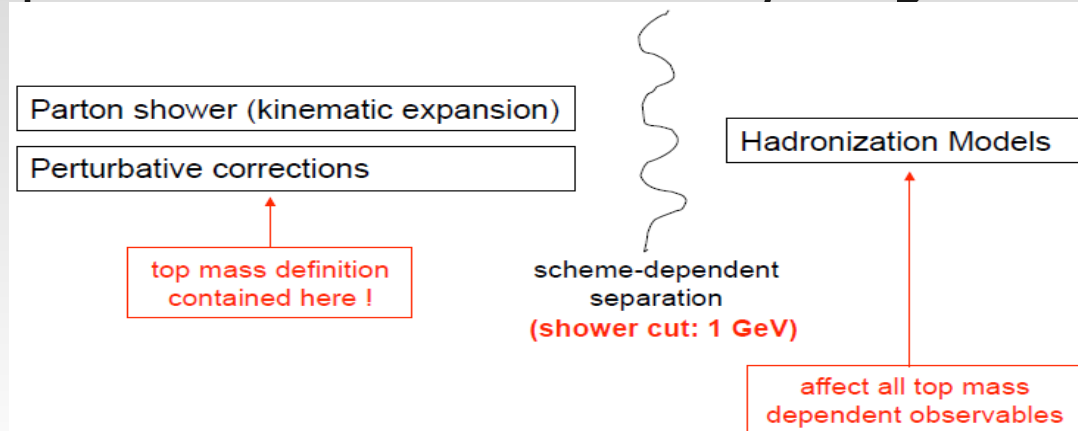


Top Quark Mass Results



Top Quark Mass: What are we measuring?

- Direct M_{top} measurements make heavy usage of Monte Carlo



- So we measure the Monte Carlo top mass! What is in the Monte Carlo?
- Masses in Quantum Field theory:

Pole mass: based on the concept of free particle, usable only in perturbation theory ($i \frac{p+m}{p^2 - m^2 + i\epsilon}$), does not apply to quark

MS (Mass Scheme): $m_{top}^{pole} +$ corrections due to the interaction

Conclusion
$$m_{top}^{MC}(R_{sc}) = m_{top}^{pole} - R_{sc} c \left[\frac{\alpha_s}{\pi} \right] \quad R_{sc} \approx 1 \text{ GeV} \text{ Shower cut-off}$$

Detailed discussion <http://arxiv.org/abs/0808.0222v2>

Top Mas Measurements from jets and the Tevatron Top-Quark Mass

Andre' H. Hoang, Iain W. Stewart +talks pf A.H. Hoang

Top Quark Mass Measurement from Cross Section

- The top quark mass can be extracted from the cross section measurement using final states that have weak dependence on the top mass. The measured cross section is compared to the NNLO theory prediction where the top mass is a parameter and can be defined in a not ambiguous way
- This measurement is a important QCD test where the $\sigma(m_{top}^{pole})$ is verified.
- Method used:
 - The theoretical cross section as function of m_{top}^{pole} is calculated using different NNLO approximation.
 - Cross section parameterization is extracted from data:

$$\sigma_{t\bar{t}}(m_t^{MC}) = \frac{1}{(m_t^{MC})^4} [a + b(m_t^{MC} - m_0) + c(m_t^{MC} - m_0)^2 + d(m_t^{MC} - m_0)^3]$$

Parameters a, b determined from data
 $m_0 = 170 \text{ GeV}$

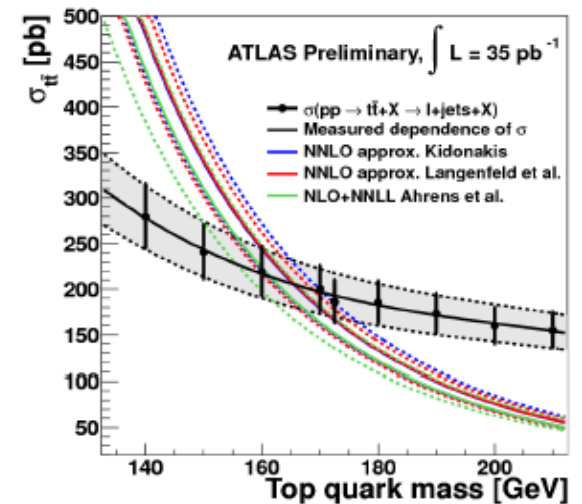
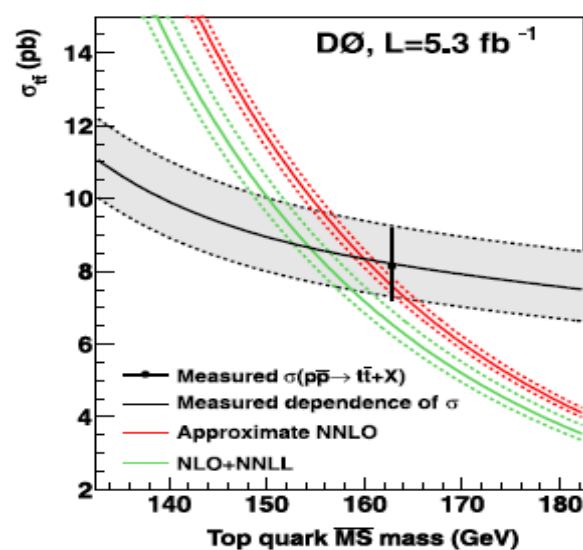
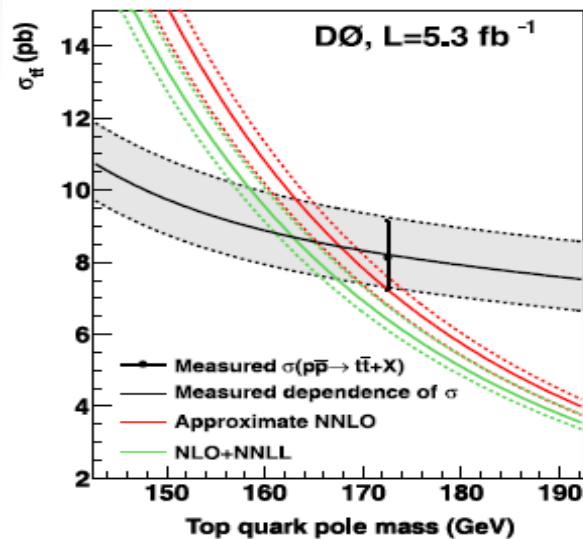
Top Quark Mass Measurement from Cross Section

m_{top} is determined from the joint likelihood

$$L(m_t) = \int f_{\text{exp}}(\sigma | m_t) [f_{\text{scale}}(\sigma | m_t) \otimes f_{\text{PDF}}(\sigma | m_t)] d\sigma$$

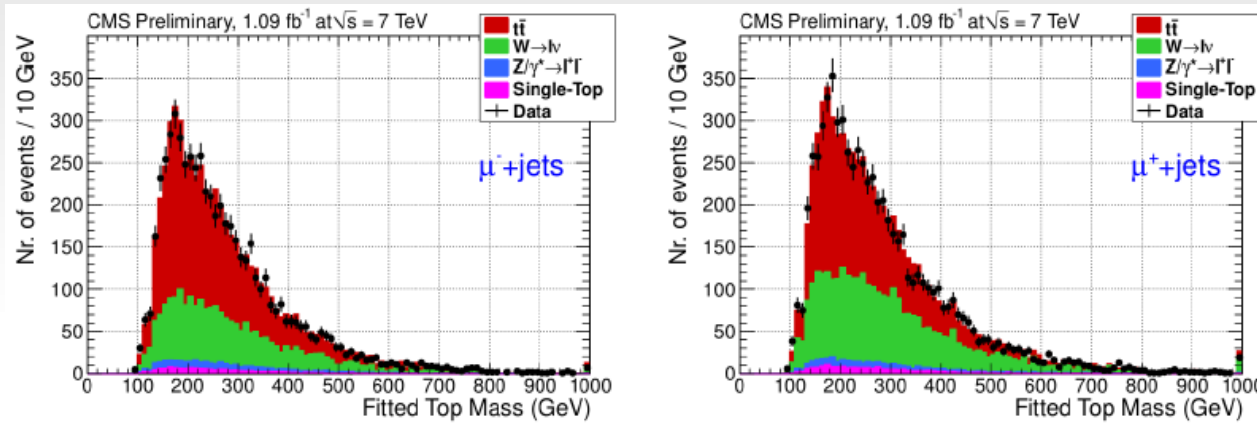
Experimental function

Theoretical function

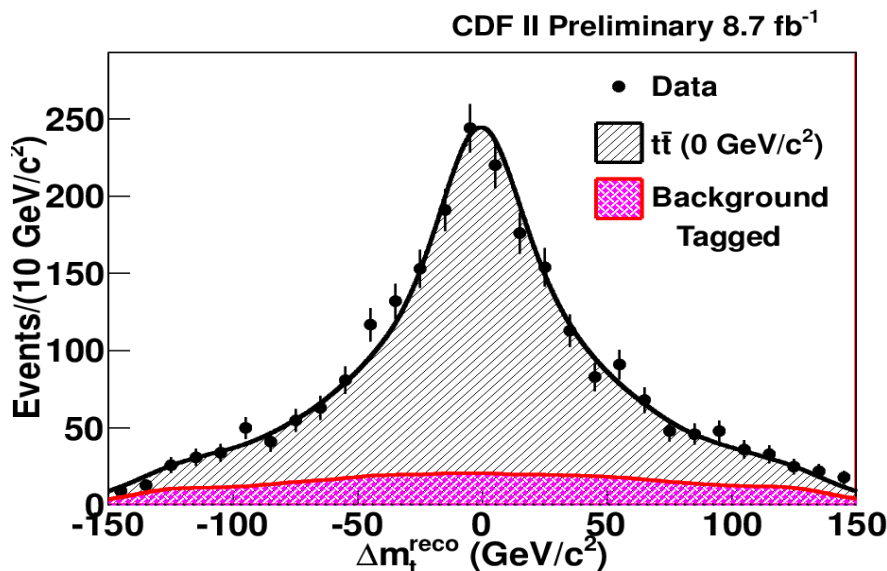


Top-anti-Top Quark Mass Difference

With the top quark is possible to test the CPT invariance in the quark system. The data used to measure the mass is also fitted for the mass difference Δm .



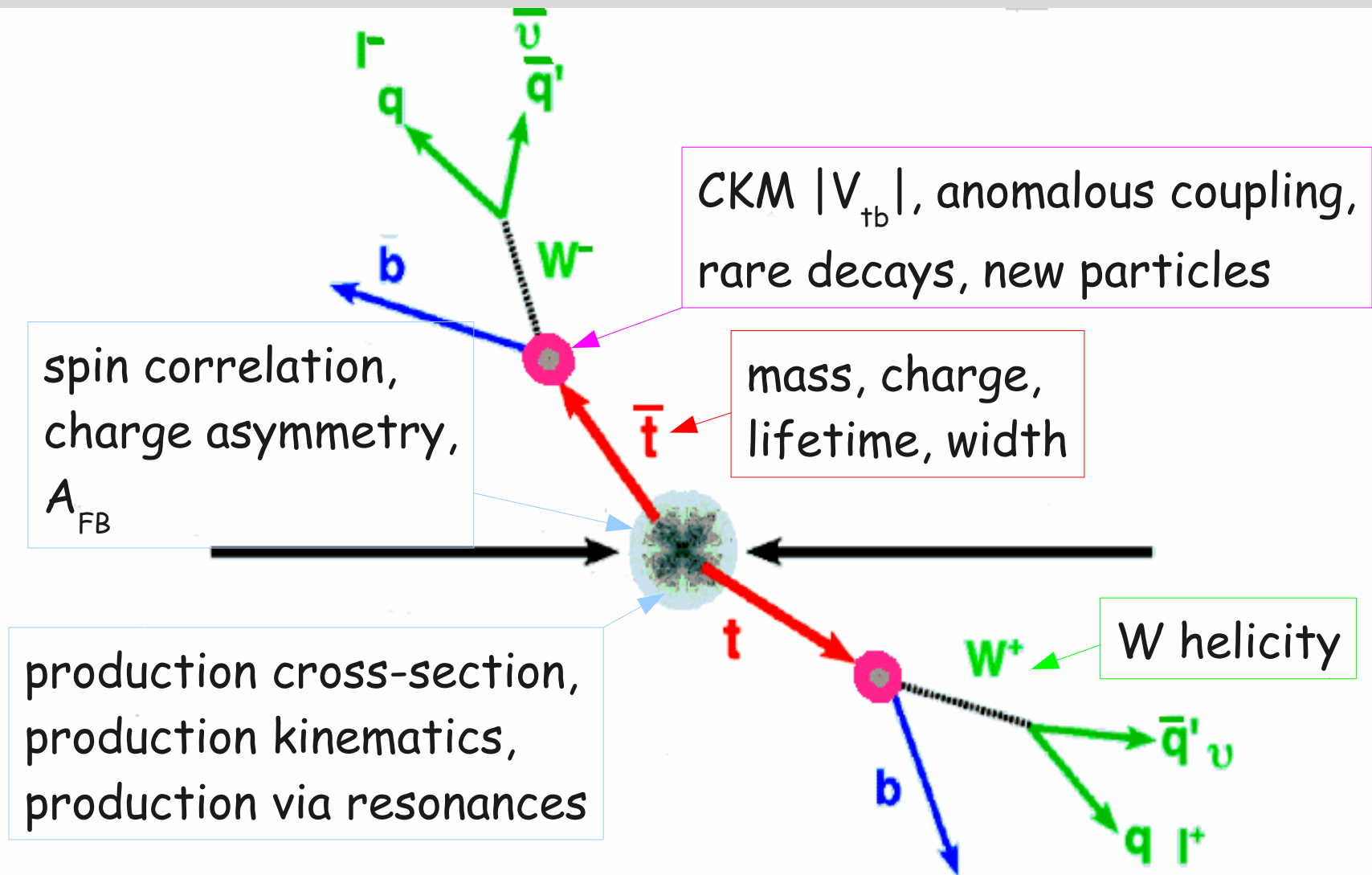
LHC measure the m_{top} and $m_{\text{anti-top}}$



Tevatron measure Δm

Δm consistent with 0

Top Quark Properties



CKM V_{tb} measurement: Introduction

- In the SM $SU(2) \times U(1)$ quarks and leptons are assigned to be left-handed doublets and right-handed singlet
- Quark mass eigenstates are not the same as the weak eigenstates, the matrix relating these bases defined for 6 quarks and parameterized by Kobayashi and Maskawa by generalization of 4 quark case described by the Cabibbo angle
- By convention, the matrix is often expressed in terms of a 3x3 unitary matrix, V , operating on the charge $-1/3$ quark eigenstates (d, s, b):

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \underbrace{\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}}_{V_{CKM}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Elements depend on 4 real parameters (3 angles and 1 CPV phase)
 V_{CKM} is the only source of CPV in the SM

CKM V_{tb} measurement

V_{tb} can be measured using top events in two different way:

1. indirect, by using $t\bar{t}$ events
2. direct with single-top events

1. the ratio $R = \frac{BR(t \rightarrow Wb)}{BR(t \rightarrow Wq)} = \frac{|V_{tb}|^2}{|V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2}$ is obtained from

events with 0, 1 and 2 tags (two different taggers are used)

$$N_{00} = n_0 + (1 - \epsilon_l)(1 - \epsilon_s)n_1 + (1 - \epsilon_l)^2(1 - \epsilon_s)^2n_2 + F_{00} \quad n_0 = N_{top}[a_0 + (1 - R)a_1 + (1 - R)^2a_2]$$

$$N_{01} = \epsilon_l(1 - \epsilon_s)n_1 + \epsilon_l(2 - \epsilon_l)(1 - \epsilon_s)^2n_2 + F_{01} \quad n_1 = N_{top}[Ra_1 + 2R(1 - R)a_2]$$

$$N_1 = \epsilon_s n_1 + 2\epsilon_s(1 - \epsilon_s)n_2 + F_1 \quad n_2 = N_{top}R^2a_2$$

$$N_2 = \epsilon_s^2 n_2 + F_2$$

Number of background events

Number of $t\bar{t}$ events

$\epsilon_{s,l}$ a_i Tagging efficiency and acceptance

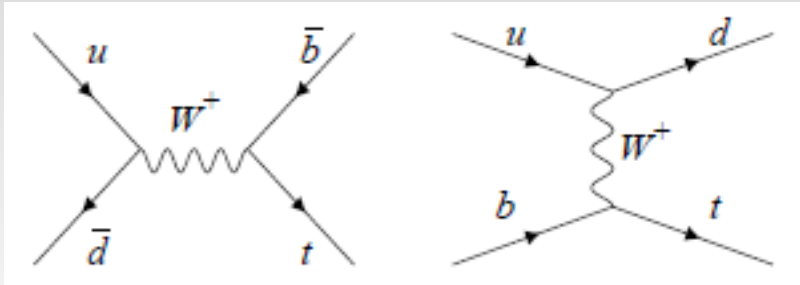
$$R = 0.94^{+0.31}_{-0.24}$$

Three generation unitarity \rightarrow denominator=1

$$V_{tb} = 0.97^{+0.16}_{-0.12}$$

CKM V_{tb} measurement

2. Single top production is dominated by s and t channel



The cross section $\propto |V_{tb}|^2$ from which it can be extracted:

$$|V_{tb}|_{\text{measured}}^2 = \sigma_{s+t}^{\text{measured}} |V_{tb}|_{\text{SM}}^2 / \sigma_{s+t}^{\text{SM}}$$

With the assumption that $t \rightarrow Wb$, $|V_{tb}|^2 \gg |V_{td}|^2 + |V_{ts}|^2$

With the measured cross section

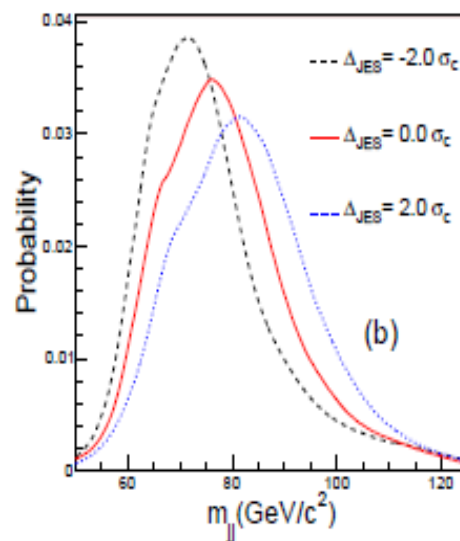
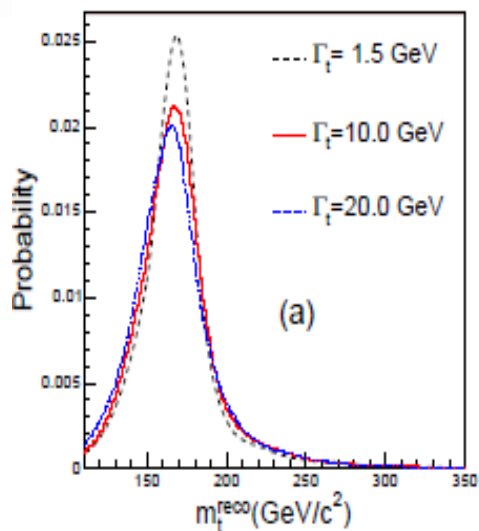
$$|V_{tb}| = 0.91 \pm 0.11(\text{stat.} + \text{syst.}) \pm 0.07(\text{theory})$$

Top Quark Width and Lifetime

At LO the total top width $\Gamma_t^0 = |V_{tb}|^2 G_F m_t^3 / 8\pi\sqrt{2}$

If $|V_{tb}| \approx 1$ $\Gamma_t^0 = 1.3 \text{ GeV}$ assuming $m_{\text{top}} = 172.5 \text{ GeV}$ which correspond to a lifetime of $5 \times 10^{-25} \text{ s}$. Deviation from the expected value could indicate new physics.

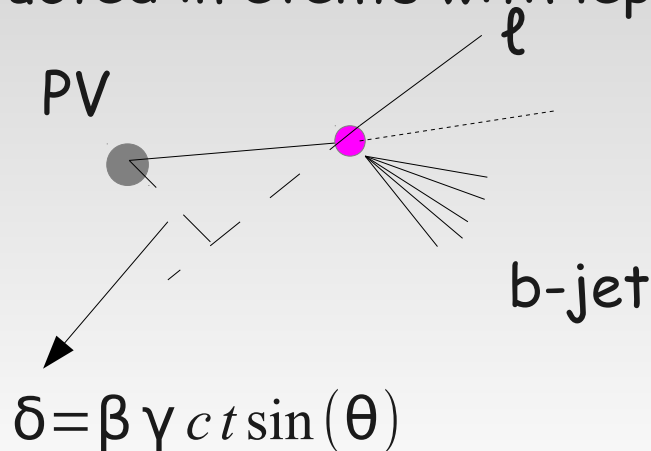
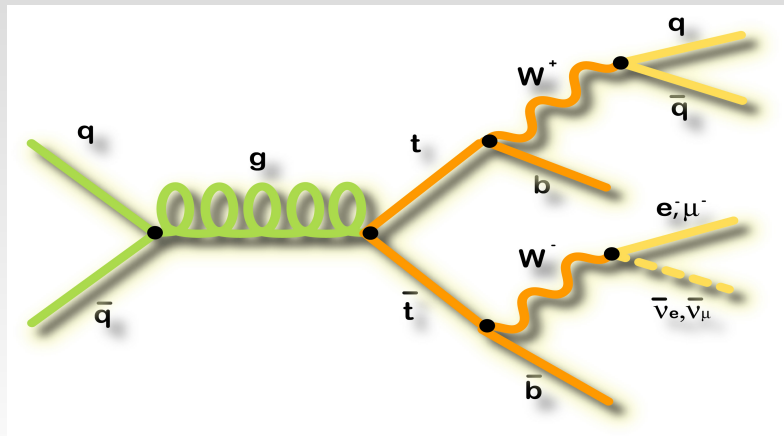
In events lepton+jets two observable are used : $m_{\text{+}}^{\text{reco}}$ and m_{jj} and reconstructed for each event as function of Γ and ΔE



A likelihood, built using these template is minimized from which is extracted $\Gamma_{\text{+}} < 7.6 \text{ GeV @95 CL}$

Top Quark Width and Lifetime

Top quark candidate are reconstructed in events with lepton + jets



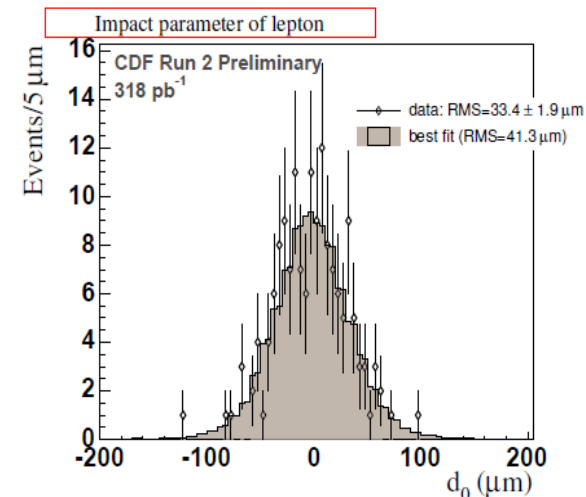
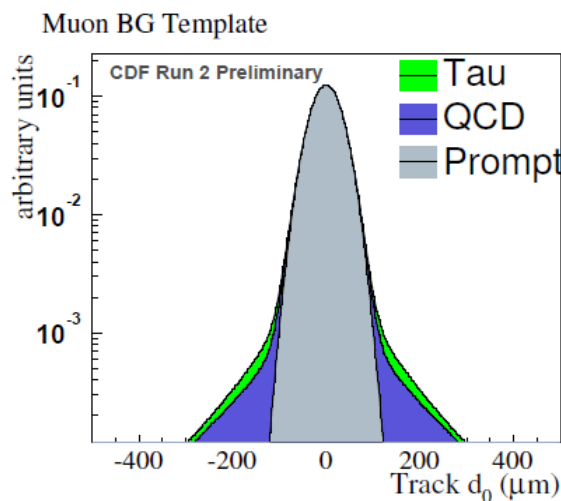
The impact parameter distribution is proportional to the lifetime.

The measured impact parameter distribution has several components:

- detector resolution
- background
- top quark

$$c\tau < 52.5 \mu\text{m}$$

$$\rightarrow 17.5 \times 10^{-14} \text{s}$$



Top Quark Charge

In the SM top quark is supposed to have charge $+2/3$, $t \rightarrow W^+b$ but for long time its charge was not measured.

In hep-ph/9810531 it is proposed an exotic 4th generation model: what has been observed it is not the SM top but a particle of this family decaying $t \rightarrow W^-b$ with charge $-4/3$. In this scenario the top quark has mass ~ 230 GeV.

Top quark charge measurement: $t \rightarrow W^+b$ lepton+jets

1. infer the charge of the W by using the charge of the lepton

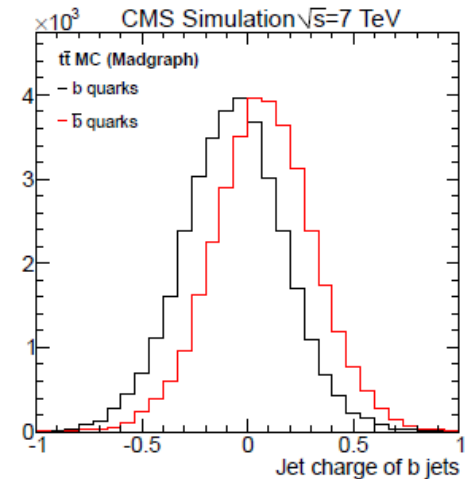
2. determine the charge of the b jet:

a. jet charge method: for each b-tagged jet

loop over
all tracks,
evaluated Q

$$Q_{b\text{-jet}} = \frac{\sum_i q_i \cdot (\vec{p}_i \cdot \hat{a})^x}{\sum_i (\vec{p}_i \cdot \hat{a})^x}$$

x = weighting factor
 \hat{a} = jet axis
 \vec{p}_i = track momentum



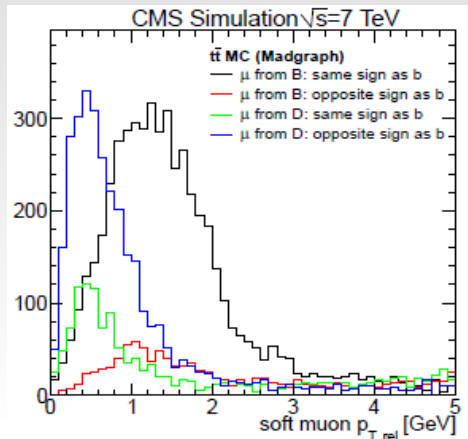
Jet charge method is calibrated on MC and data, it has high efficiency but low purity: efficiency $\sim 98\%$ purity $\sim 60\%$

Top Quark Charge - 2

b. soft lepton method: $b \rightarrow \ell^-$

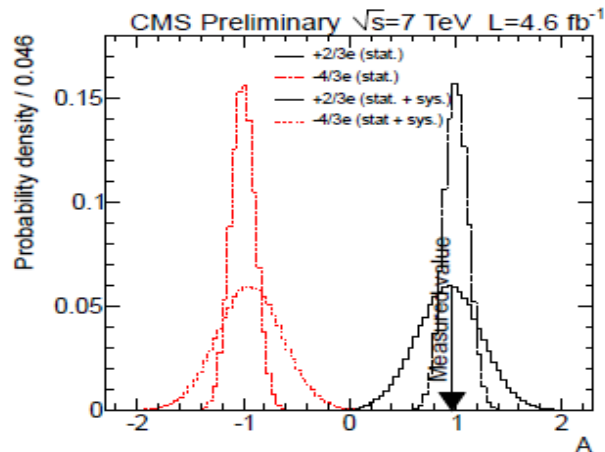
identify a lepton with low momentum inside the b-jet

use $p_{T,rel}$ to distinguish lepton from b and charm



3. pair W and b to form "right" top: fit three jets invariant mass, one b-tagged, the best combination \rightarrow the candidate top.

The other b-jet is paired to the W.



Exotic top-quark model is excluded at 99%CL

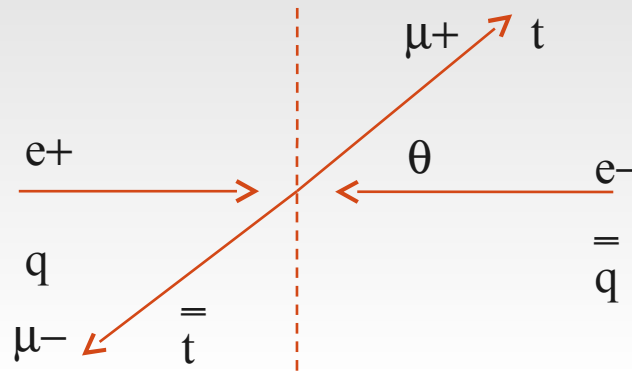
Top anti-Top Quark Charge Asymmetry

At LO the number of top quark produced at a given angle is expected to be almost equal to the number of anti-top quark produced at the same angle.

In analogy to the muon production asymmetry for the $t\bar{t}$ system it is possible to define the

asymmetry using the rapidity, in lepton+jets events

$\Delta y = y_t - y_{\bar{t}} = q_l (y_{leptonic} - y_{hadronic})$ rapidity difference invariant to z-boost

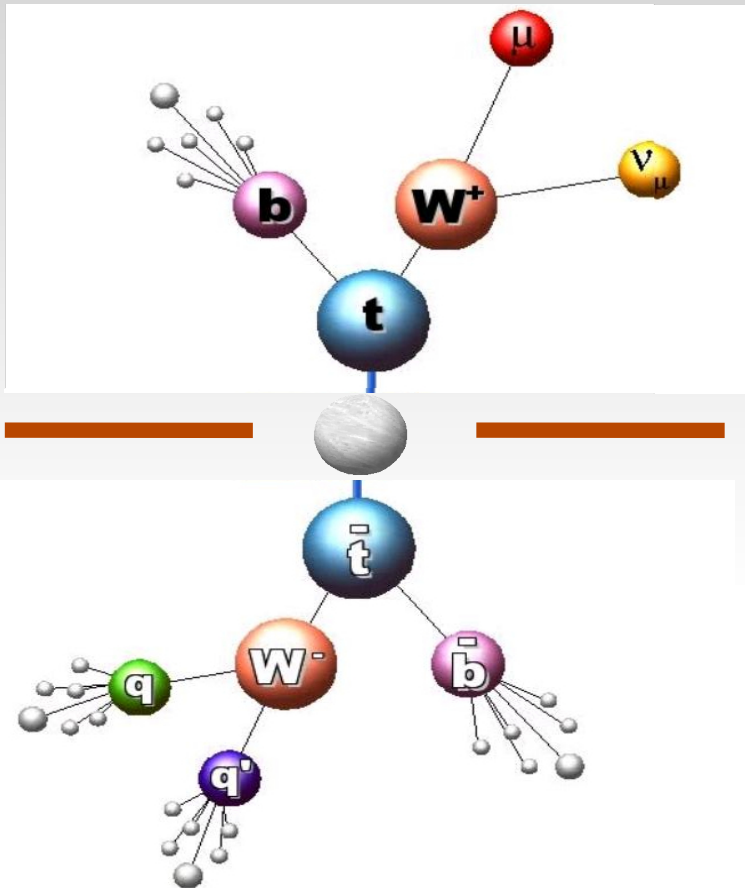


$$A = \frac{N(\cos\theta > 0) - N(\cos\theta < 0)}{N(\cos\theta > 0) + N(\cos\theta < 0)}$$

$$A = \frac{N(\Delta y > 0) - N(\Delta y < 0)}{N(\Delta y > 0) + N(\Delta y < 0)}$$

Asymmetry as function of Δy is the same in the lab and $t\bar{t}$ frame

Top anti-Top Quark Charge Asymmetry



Usual requirements for the lepton+jets reconstruction.

The charge of the lepton determine which reconstructed quark is top.

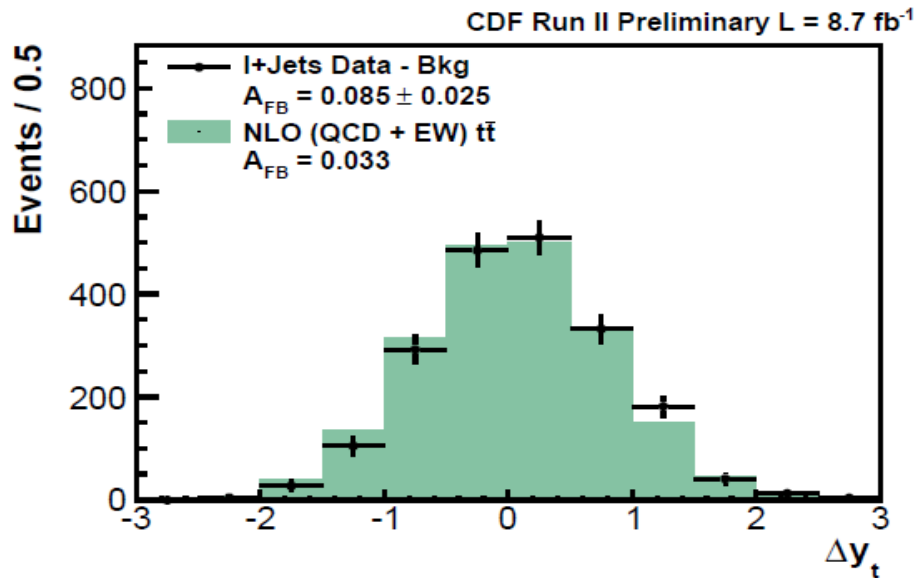
Lepton angles are very well measured.

The SM predictions are calculated by using different Monte Carlo

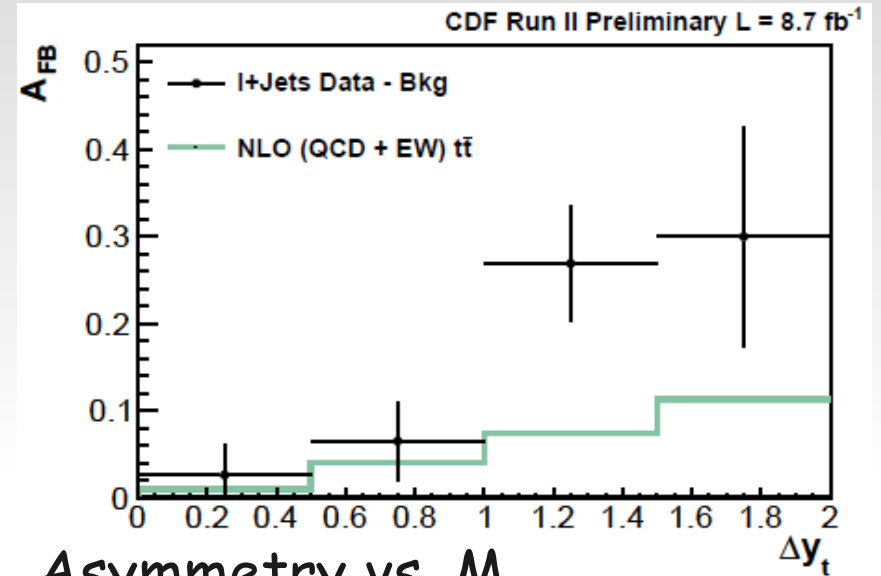
	MC@NLO	POWHEG	MCFM
Inclusive	0.067	0.066	0.073
$ \Delta y < 1$	0.047	0.043	0.049
$ \Delta y > 1$	0.130	0.139	0.150
$M_{t\bar{t}} < 450 \text{ GeV}/c^2$	0.054	0.047	0.050
$M_{t\bar{t}} > 450 \text{ GeV}/c^2$	0.089	0.100	0.110

Top anti-Top Quark Charge Asymmetry

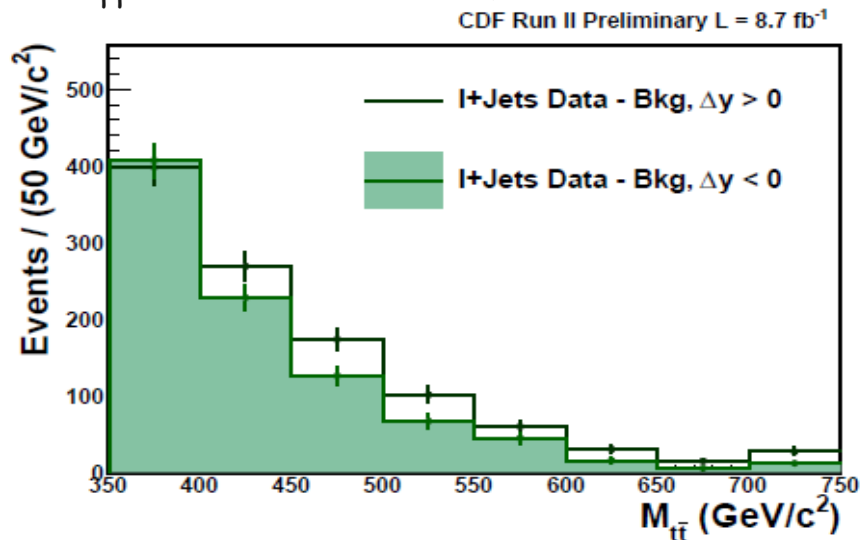
Asymmetry



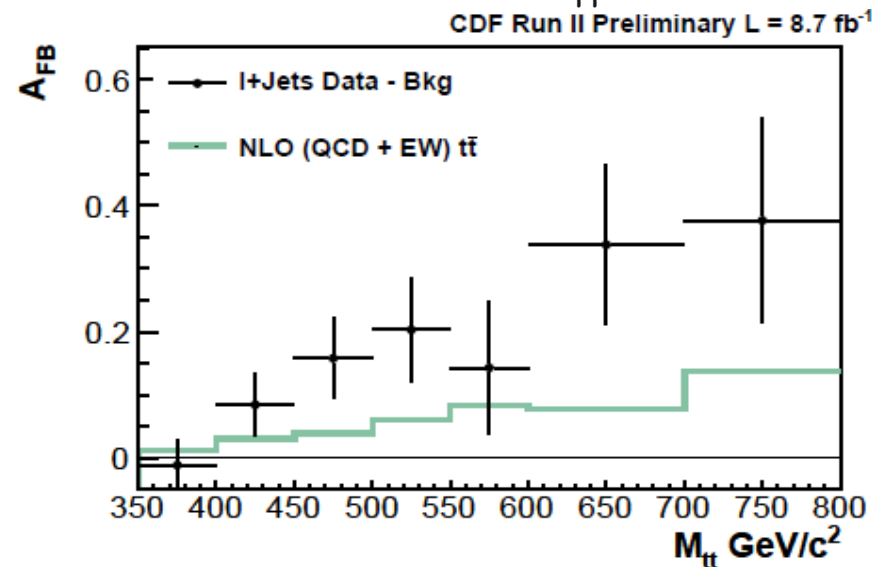
Asymmetry vs. Δy



$M_{t\bar{t}}$ for events $\Delta y > 0$

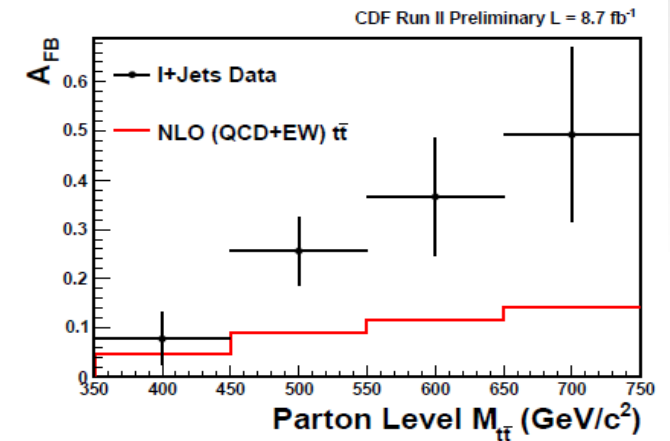
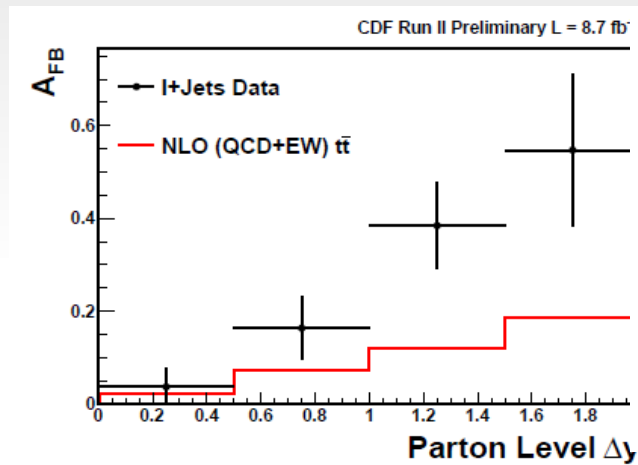
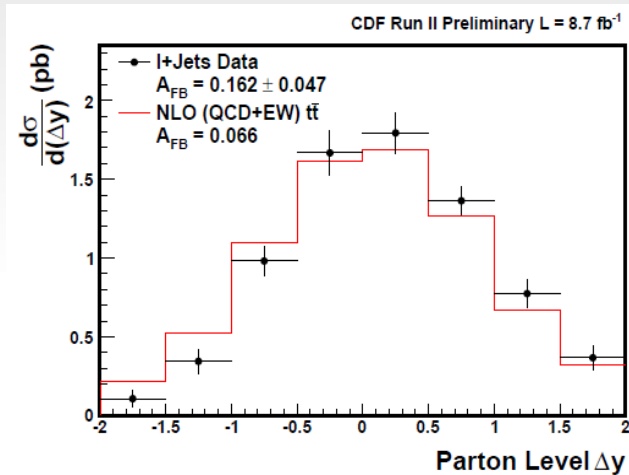


Asymmetry vs. $M_{t\bar{t}}$



Top anti-Top Quark Charge Asymmetry

In order to compare the measured asymmetry directly to theory Prediction we have to go back to the parton level asymmetry. This is done with the "unfolding" procedure.



Disagreement with SM

Top anti-Top Quark Charge Asymmetry @ LHC

Differences respect to Tevatron:

1. large fraction of $t\bar{t}$ is produced by gluon fusion and the asymmetry is present only in $q\bar{q}$ initial states
2. at LHC quarks are mainly valence quarks while anti-quarks are sea quarks and the larger average momentum of the valence quarks produce an excess of top quark produced in the forward region

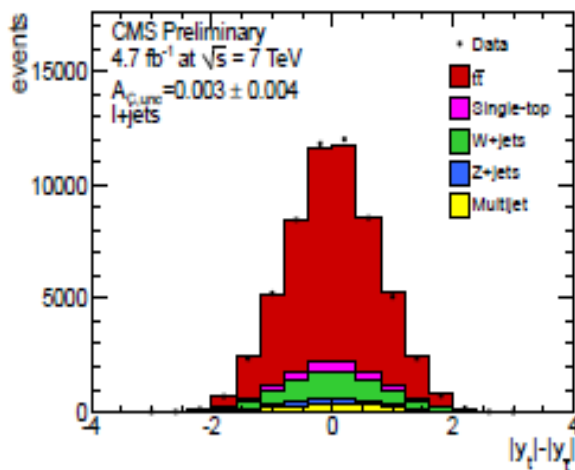


The expected asymmetry $A(\text{theory}) = 0.0115 \pm 0.0006$

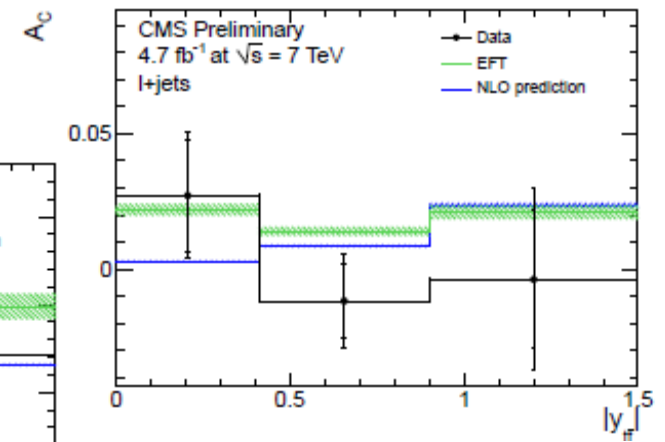
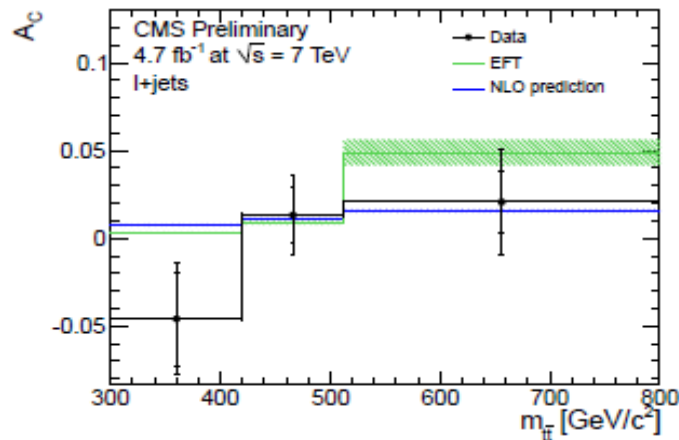
Top anti-Top Quark Charge Asymmetry @ LHC

The asymmetry is based on the fully reconstructed four-momenta of t and \bar{t} in each event

The reconstructed four-vectors are used to obtain the inclusive and differential distributions of Δy and the charge asymmetry is calculated by counting the entries with $\Delta y > 0$ and the entries with $\Delta y < 0$

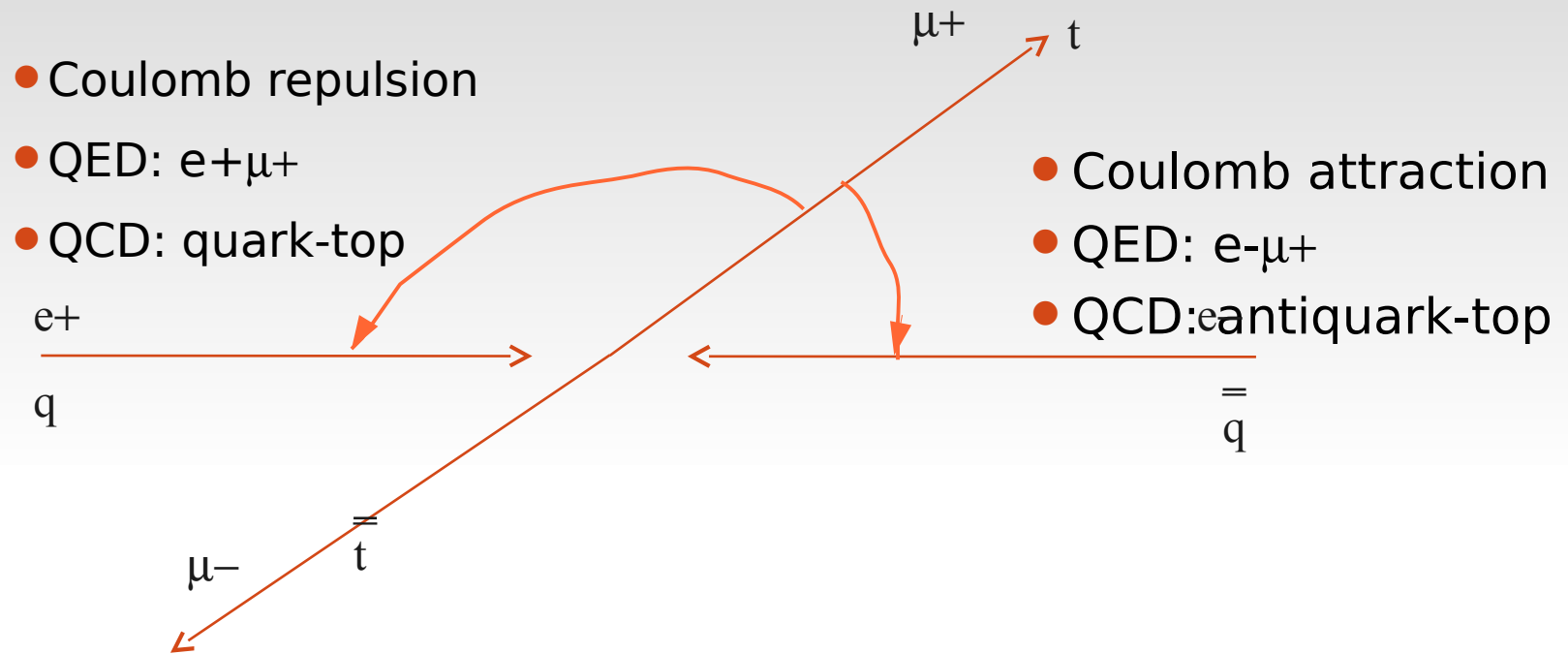


After the unfolding



No disagreement with SM

Top anti-Top Quark Charge Asymmetry Meaning



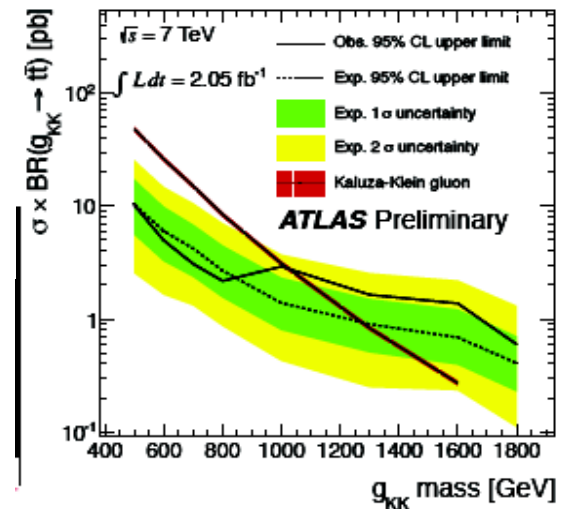
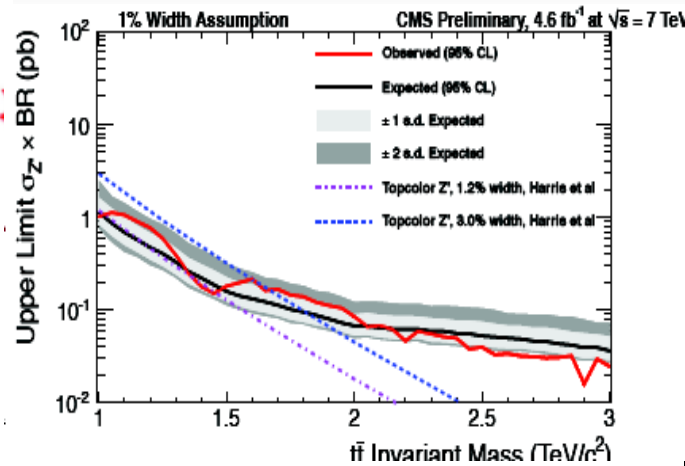
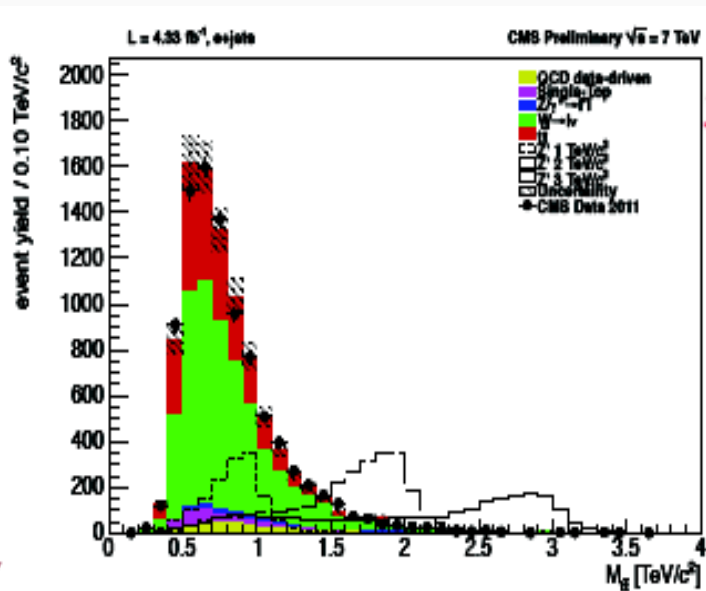
New kind of interactions to explain the $t\bar{t}$ asymmetry

- gluon interferes with an axial object arising from an extended strong gauge group or extra-dimensions
- objects with flavor violating couplings create an asymmetry via a $u/d \rightarrow t$ flavor change into the forward Rutherford peak.

Mtt resonances

Several model of Physics beyond SM assume particle decaying into $t\bar{t}$.

1. $Z' \rightarrow t\bar{t}$
 2. Heavy states expressed as Kaluza-Klein gluons couple to $t\bar{t}$
 3. axigluon models, proposed to solve the discrepancy in the top pair production forward-backward asymmetry from the Tevatron
- Analysis performed in any top decay channel



No evidence up \sim TeV