# Determining the WIMP mass using the complementarity between direct detection, indirect detection and the ILC

Andreas Goudelis



Laboratoire de Physique Théorique - Orsay, France



INFN - Padova, Italy

Based on N.Bernal, A.G., Y.Mambrini, C.Muñoz, JCAP 0901 (2009) 046 (arXiv:0804.1976 [hep-ph])

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## Outline





- Direct Detection
- Indirect Detection

Oirect vs Indirect Detection

#### Dark Matter at Colliders

- One idea
- Radiative WIMP Production Rate
- ILC: Results

#### 6 Complementarity - Conclusions

Dark matter detection is a quite challenging task...

- Parameter spaces for thermal candidates are seriously starting to be explored now: CDMS-II, Xenon, Fermi, PAMELA...
- The situation can be even worse for non-thermal candidates (gravitinos, right-handed neutrinos etc)
- "Positive detection" reports generate much controversy: HESS, EGRET, DAMA, PAMELA...
- An example: Models can explain  $e^+$  data from PAMELA but fail elsewhere.

 $\implies$  Need to combine as much information as possible.

A slightly different question:

- Suppose we see an excess.
- Suppose this excess can be nicely explained through DM annihilations/scatterings/pair-production etc...

How well shall we be able to constrain the DM properties?

 $\hookrightarrow$  We shall be focusing on WIMPs.

Direct Detection Indirect Detection

## Direct Detection: The event rate

$$\frac{dN}{dE_r} = \frac{\sigma_{\chi-N} \cdot \rho_0}{2 M_r^2 m_{\chi}} F(E_r)^2 \int_{v_{\min}(E_r)}^{v_{esc}} \frac{f(v)}{v} dv$$

Where:

- N: Number of scatterings (s<sup>-1</sup>kg<sup>-1</sup>)
- Er: Nuclear recoil energy (~few keV)
- $m_{\chi}$ : WIMP mass
- $M_r = \frac{m_{\chi} m_N}{m_{\chi} + m_N}$ : WIMP Nucleus Reduced Mass
- $\sigma_{\chi-N}$ : WIMP-Nucleus cross-section (Spin-independent coupling)
- $\rho_0$ : Local WIMP density (0.3 GeV cm<sup>-3</sup>)
- f(v): WIMP local velocity distribution (Maxwell-Boltzmann)
- F: Nuclear form factor (Woods-Saxon)

Direct Detection Indirect Detection

## Direct Detection in a XENON-like experiment

#### Mass and cross-section discrimination

Mass Resolution



Ignoring backgrounds/theoretical uncertainties.

Direct Detection Indirect Detection

Indirect Detection: The  $\gamma$ -ray flux

$$\begin{aligned} p_{\gamma}(E_{\gamma}) &= 0.94 \cdot 10^{-13} \mathrm{cm}^{-2} \mathrm{sec}^{-1} \mathrm{GeV}^{-1} \mathrm{sr}^{-1} \sum_{i} Br_{i} \frac{dN_{\gamma}^{i}}{dE_{\gamma}} \\ &\times \left(\frac{\langle \sigma \mathbf{v} \rangle}{10^{-29} \mathrm{cm}^{3} \mathrm{sec}^{-1}}\right) \left(\frac{100 \mathrm{GeV}}{m_{\chi}}\right)^{2} \overline{J}(\Delta \Omega) \Delta \Omega \end{aligned}$$

Where:

- Bri: Branching Ratio of annihilation into i-th SM particle
- $dN_{\gamma}^{i}/dE_{\gamma}$ : Functions describing SM particles' decays into  $\gamma$ -rays (PYTHIA + fit)
- $\langle \sigma v \rangle$ : Total WIMP self-annihilation cross-section ( $\approx 3 \cdot 10^{-26} \text{cm}^3 \text{sec}^{-1}$ )
- $\overline{J}$ : Astrophysical factor: Depends on DM distribution.

	a (kpc)	$\alpha$	$\beta$	$\gamma$	$\overline{J}(4 \cdot 10^{-3} \mathrm{sr})$
NFW	20	1	3	1	$5.859 \cdot 10^{2}$
NFW <sub>c</sub>	20	0.8	2.7	1.45	$3.254 \cdot 10^4$
Moore et al.	28	1.5	3	1.5	$2.574 \cdot 10^4$
Moore <sub>c</sub>	28	0.8	2.7	1.65	$3.075 \cdot 10^5$

Direct Detection Indirect Detection

## Indirect Detection with FERMI

## Mass and cross-section discrimination for a NFW halo profile





- Again, better resolution for smaller masses (due to strong differences in the spectrum form in the [1, 300] GeV region!).
- Strong dependence on the halo profile.
- Background: EGRET PS + HESS PS + HESS diffuse.

6-year mission, GC in field of view half of the time. Interesting theoretical/experimental uncertainties!

## Direct vs Indirect Detection

Example for  $m_{\chi} = 50 \text{GeV}$ 

Example for  $m_{\chi} = 100 \text{GeV}$ 



• Complementarity more obvious for not too optimistic astrophysical considerations and  $\sigma_{\chi-p}$  cross-sections.

• In any case, results can be improved with a more elaborate statistical treatment.

One idea Radiative WIMP Production Rate ILC: Results

## ILC: The Method

The General Idea: Birkedal, Matchev, Perelstein, arXiv:hep-ph/0403004



One idea Radiative WIMP Production Rate ILC: Results

## ILC: Radiative WIMP Production Rate

$$\frac{d\sigma(e^+e^- \rightarrow 2\chi + \gamma)}{dxd\cos\theta} \approx \frac{\alpha\kappa_e\sigma_{an}}{16\pi} \frac{1 + (1 - x)^2}{x} \frac{1}{\sin^2\theta} 2^{2J_0}$$
$$\times (2S_{\chi} + 1)^2 \left(1 - \frac{4m_{\chi}^2}{(1 - x)s}\right)^{1/2 + J_0}$$

Where:

• 
$$x = 2E_{\gamma}/\sqrt{s}$$

- $\theta$ : Photon emission angle
- $\sigma_{an}$ : Total annihilation cross-section (WMAP) ( $\sim$  7pb, under assupmptions.)
- J<sub>0</sub>: Dominant (s- or t-) annihilation channel
- $S_{\chi}$ : WIMP's spin

• 
$$\kappa_e = \sigma_e^{(J_0)} / \sigma_{an}$$
: Annihilation fraction into  $e^+e^-$  pairs.

One idea Radiative WIMP Production Rate ILC: Results

## **ILC:** Subtleties

Approach valid for soft/collinear photons  $\rightarrow$  Undetectable!

 $\bullet$  Nevertheless, it gives satisfactory results outside the soft/collinear region for nonrelativistic WIMPs.

To ensure the WIMPs' non-relativistic nature we impose the following kinematical cuts:

$$rac{\sqrt{s}}{2}\left(1-rac{8m_\chi^2}{s}
ight)\leq E_\gamma\leq rac{\sqrt{s}}{2}\left(1-rac{4m_\chi^2}{s}
ight).$$

• Main background process: Radiative neutrino production (CalcHEP 2.5).

One idea Radiative WIMP Production Rate ILC: Results

## ILC: Mass Discrimination

#### WIMP mass - annihilation fraction discrimination capacity

#### Relative error in WIMP's mass discrimination



- Discrimination capacity peaks significantly for  $m_{\chi} = 175 \text{GeV}$  (optimal combination of uncut spectrum phase space).
- Significant improvement in mass resolution for polarized beams.

## DM Experiments/ILC complementarity



An example at 95% CL:

- $m_{\chi} = 100 {
  m GeV}$
- 3 years of exposure,

• 
$$\sigma_{\chi-p}=10^{-8}$$
 pb

- NFW profile and a
- 500 GeV unpolarized linear collider with an integrated luminosity of  $500 {\rm fb}^{-1}$

$m_{\chi}$	XENON	GLAST	ILC
50 GeV	$-5/+7 { m GeV}$	$\pm 12$ GeV	—
100 GeV	-19/+75 GeV	$-50/+60 { m GeV}$	$-40/+20  { m GeV}$
175 GeV	-65/ GeV	-125 GeV	-20/+15 GeV
500 GeV	_	_	_

## Summarizing...

- We presented a simple way to exploit simultaneously different kinds of experiments to extract **model-independent** information on WIMP Dark Matter.
- For quite reasonable (i.e. not too optimistic) considerations on the WIMP-nucleus scattering cross-section and the DM halo profile we saw that different kinds of experiments can act highly complementary:
  - $\rightarrow$  The precision is comparable.
  - $\rightarrow$  Possibility to cover different regions in the parameter space.

#### Further questions:

- What about other channels? (p/s, antideuterons, synchrotron etc...)
- More elaborate techniques are starting to develop for colliders:  $M_{T_2}$  approach, EFT techniques.
- Even if we determine the mass of an invisible particle in a collider, another issue is its cosmological relevance!

## A few more details...

## **Discrimination** Method

Analysis based on extended likelihood function:

$$L = \frac{(N_{th}^{scan})^{N_{Exp}}}{N_{Exp}!} \exp\left(-N_{th}^{scan}\right) \prod_{i=1}^{N_{Exp}} f(E; m_{\chi}, \sigma_{\chi-p})$$

- Calculate the theoretical number of events, *N*<sub>th</sub>, for the input mass and cross-section.
- Draw an "experimental" nb of events,  $N_{Exp}$ , from a Poisson distribution.
- Scan the  $(m_{\chi}, \sigma_{\chi-N})$  parameter space and find the experiment's estimation, taking into account the theoretical nb of events of every point in the PS,  $N_{th}^{scan}$ .
- Generate a large number of experiments, repeat the procedure, pick the one that averages all experiments' results.
- From this experiment, plot  $(m_{\chi}, \sigma_{\chi-N})$  non-discrimination regions.

 $\Longrightarrow$  This method allows to account for random deviations from the expected number of events.

## Direct Detection: Uncertainties

#### **Background considerations**

Velocity distribution



Inclusion of an exponential background mimicking the signal.

$$f(v_{\chi}) d^{3}v_{\chi} = \frac{1}{(v_{\chi}^{0})^{3}\pi^{3/2}} e^{-(v_{\chi}/v_{\chi}^{0})^{2}} d^{3}v_{\chi}$$

Indirect Detection: Uncertainties

#### Fun with $\gamma$ !



#### Impact of final states



↔ Recent N-body simulations seem to disfavour highly cusped profiles.
 ↔ What about baryons?
 ↔ And the backgrounds?

 $\hookrightarrow$  We include a  $30\%\tau$  final state.  $\Leftrightarrow$  Subtle for  $\gamma$ 's: Role of other leptonic channels (signal mostly renormalized). Should look into other wavelengths.

## Halo Profiles

The most usual parametrization:

$$\rho(r) = \frac{\rho_0 [1 + (R_0/a)^{\alpha}]^{(\beta - \gamma)/\alpha}}{(r/R_0)^{\gamma} [1 + (r/a)^{\alpha}]^{(\beta - \gamma)/\alpha}}$$

#### Some well motivated profiles







## Spectral Functions

PYTHIA result fit performed through functions of the form:

$$\frac{dN_{\gamma}^{i}}{dx} = \exp[F_{i}(\ln(x))]$$

with  $x = E_{\gamma}/m_{\chi}$  and F being 7th order polynomial functions.



• The  $\tau$  spectrum has a characteristic hard form, other leptons have zero contribution.

Some Points on the ILC Treatment

• The detailed balancing equation:

$$\frac{\sigma(\chi + \chi \to X_i + \bar{X}_i)}{\sigma(X_i + \bar{X}_i \to \chi + \chi)} = 2\frac{v_X^2 (2S_X + 1)^2}{v_\chi^2 (2S_\chi + 1)^2}$$

We can expand the total thermally averaged CS:

$$\sigma_i \mathbf{v} = \sum_{J=0}^{\infty} \sigma_i^{(J)} \mathbf{v}^{2J} \stackrel{\mathbf{v} \ll \mathbf{c}}{\Longrightarrow} \sigma_{an} = \sum_i \sigma_i^{(J_0)}$$

For soft/collinear photons:

$$rac{d\sigma(e^+e^- o 2\chi+\gamma)}{dxd\cos heta}pprox \mathcal{F}(x,\cos heta) ilde{\sigma}(e^+e^- o 2\chi)$$